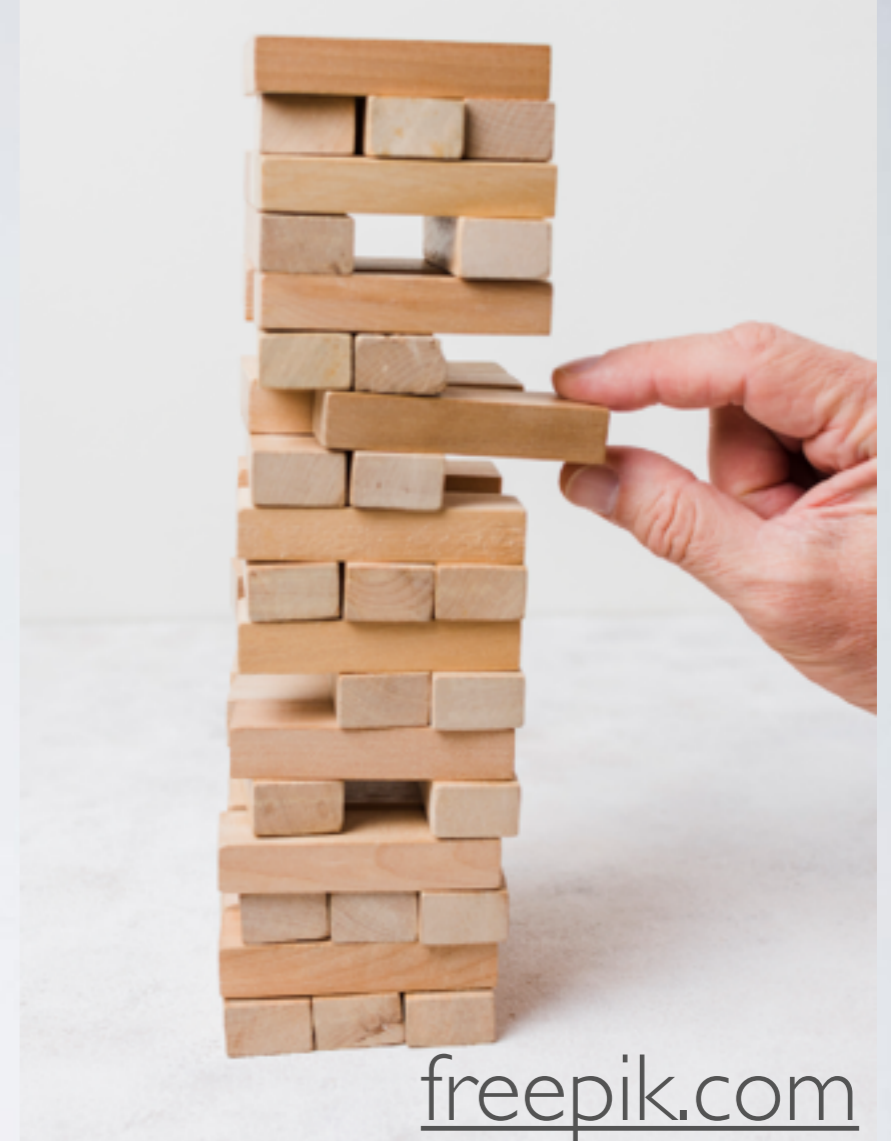


HIGGS COUPLINGS MEASUREMENTS AND THE SCALE OF NEW PHYSICS



Spencer Chang (U. Oregon)
w/ F. Abu-Ajamieh, M. Chen, M. Luty
JHEP 2020, 140 (2020) and arXiv:2009.11293

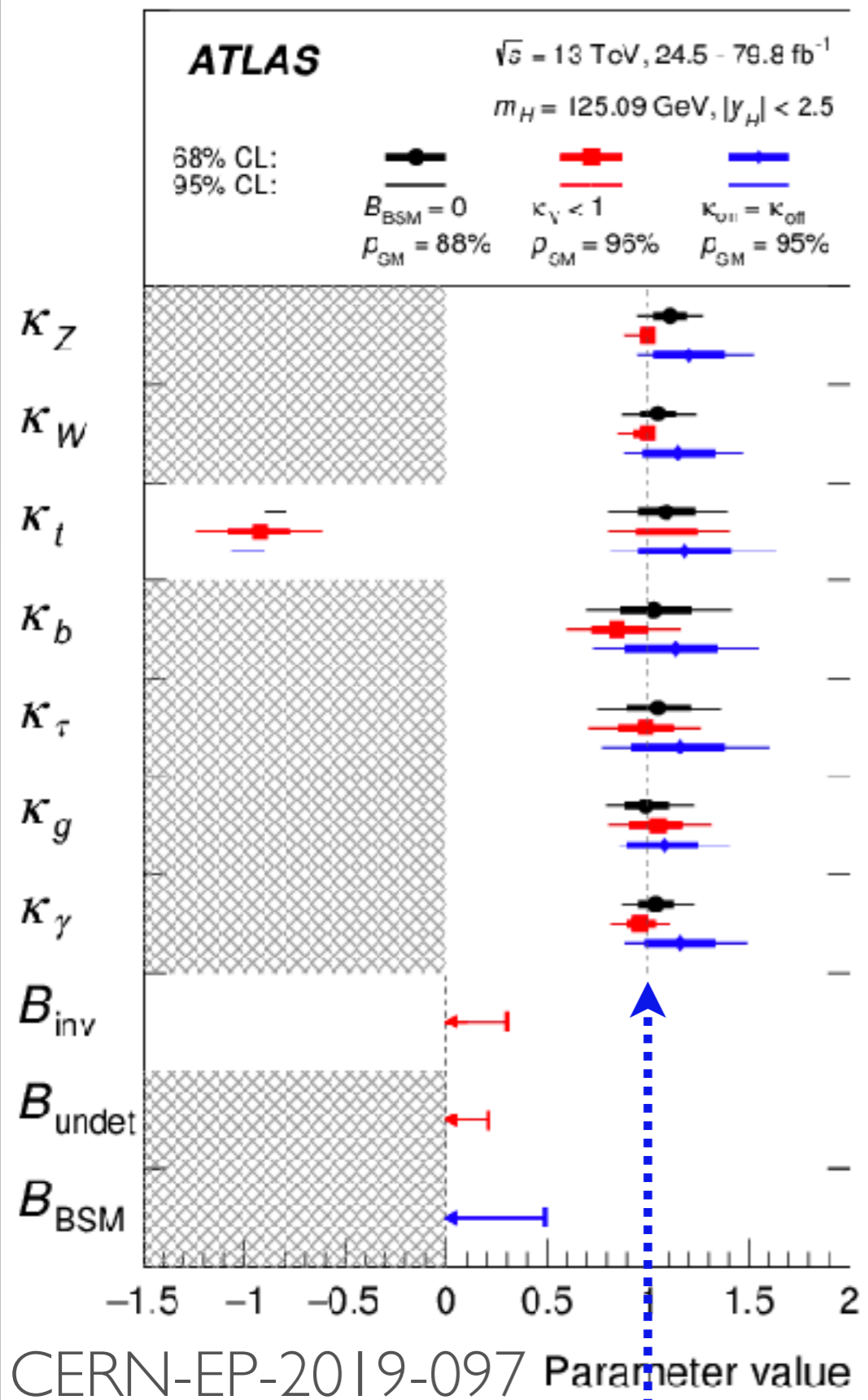
Oklahoma State 2/18/21 Seminar

PINNING DOWN HIGGS PROPERTIES



Post-discovery, goal of LHC and future colliders to measure Higgs properties to test EWSB mechanism, mass generation, and look for new physics beyond the Standard Model

HIGGS COUPLINGS MEASUREMENTS



Standard
Model values

Fits to $\sigma \times$ Branching Ratios, for Higgs couplings have **10-25%** errors and currently agree with SM value

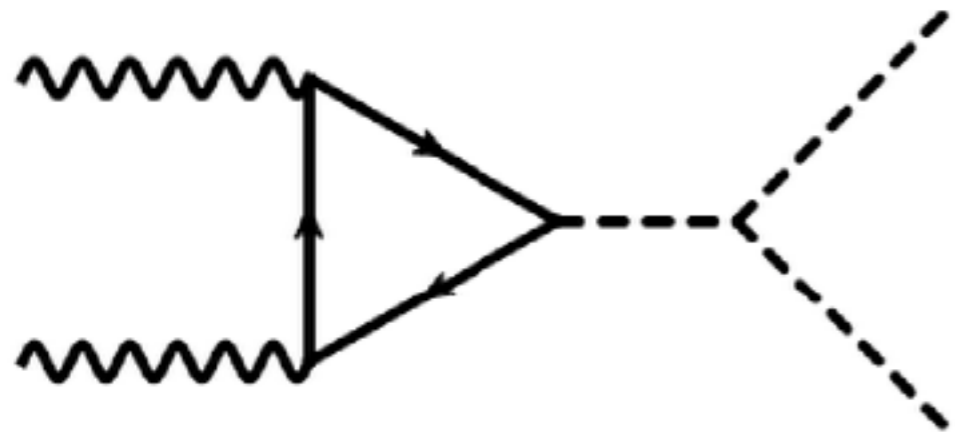
HIGGS COUPLINGS IN FUTURE

kappa-0	HL-LHC	LHeC	HE-LHC		ILC			CLIC			CEPC	FCC-ee		FCC-ee/eh/hh
			S2	S2'	250	500	1000	380	15000	3000		240	365	
κ_W [%]	1.7	0.75	1.4	0.98	1.8	0.29	0.24	0.86	0.16	0.11	1.3	1.3	0.43	0.14
κ_Z [%]	1.5	1.2	1.3	0.9	0.29	0.23	0.22	0.5	0.26	0.23	0.14	0.20	0.17	0.12
κ_g [%]	2.3	3.6	1.9	1.2	2.3	0.97	0.66	2.5	1.3	0.9	1.5	1.7	1.0	0.49
κ_γ [%]	1.9	7.6	1.6	1.2	6.7	3.4	1.9	98*	5.0	2.2	3.7	4.7	3.9	0.29
$\kappa_{Z\gamma}$ [%]	10.	—	5.7	3.8	99*	86*	85*	120*	15	6.9	8.2	81*	75*	0.69
κ_c [%]	—	4.1	—	—	2.5	1.3	0.9	4.3	1.8	1.4	2.2	1.8	1.3	0.95
κ_t [%]	3.3	—	2.8	1.7	—	6.9	1.6	—	—	2.7	—	—	—	1.0
κ_b [%]	3.6	2.1	3.2	2.3	1.8	0.58	0.48	1.9	0.46	0.37	1.2	1.3	0.67	0.43
κ_μ [%]	4.6	—	2.5	1.7	15	9.4	6.2	320*	13	5.8	8.9	10	8.9	0.41
κ_τ [%]	1.9	3.3	1.5	1.1	1.9	0.70	0.57	3.0	1.3	0.88	1.3	1.4	0.73	0.44

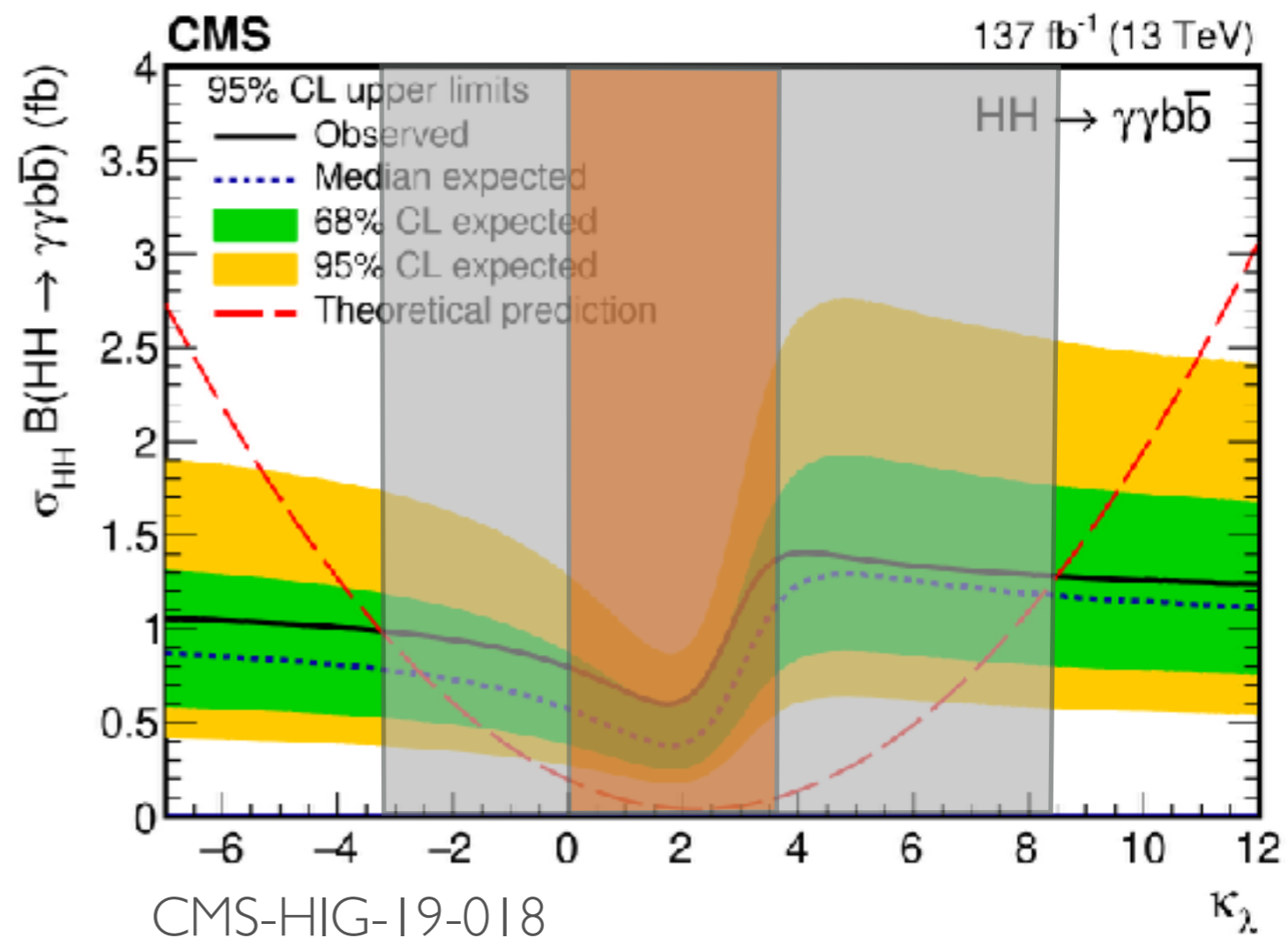
Higgs@FutureColliders report (1905.03764)

Coupling sensitivities playing a role in next collider discussion

TRILINEAR SEARCH



Trilinear probed by
search for Double Higgs
production



Currently only sensitive to $O(10)$ variations, but
projections estimate trilinear sensitivity
to $\sim [-0.2, 3.6]$ at HL-LHC w/ 3 ab^{-1} and
20-30% at future colliders

TRIPLE HIGGS PROCESS

Papaefstathiou and Sakurai
See also Chien et.al.

hh and hhh at one loop
e.g. Bizon et.al.

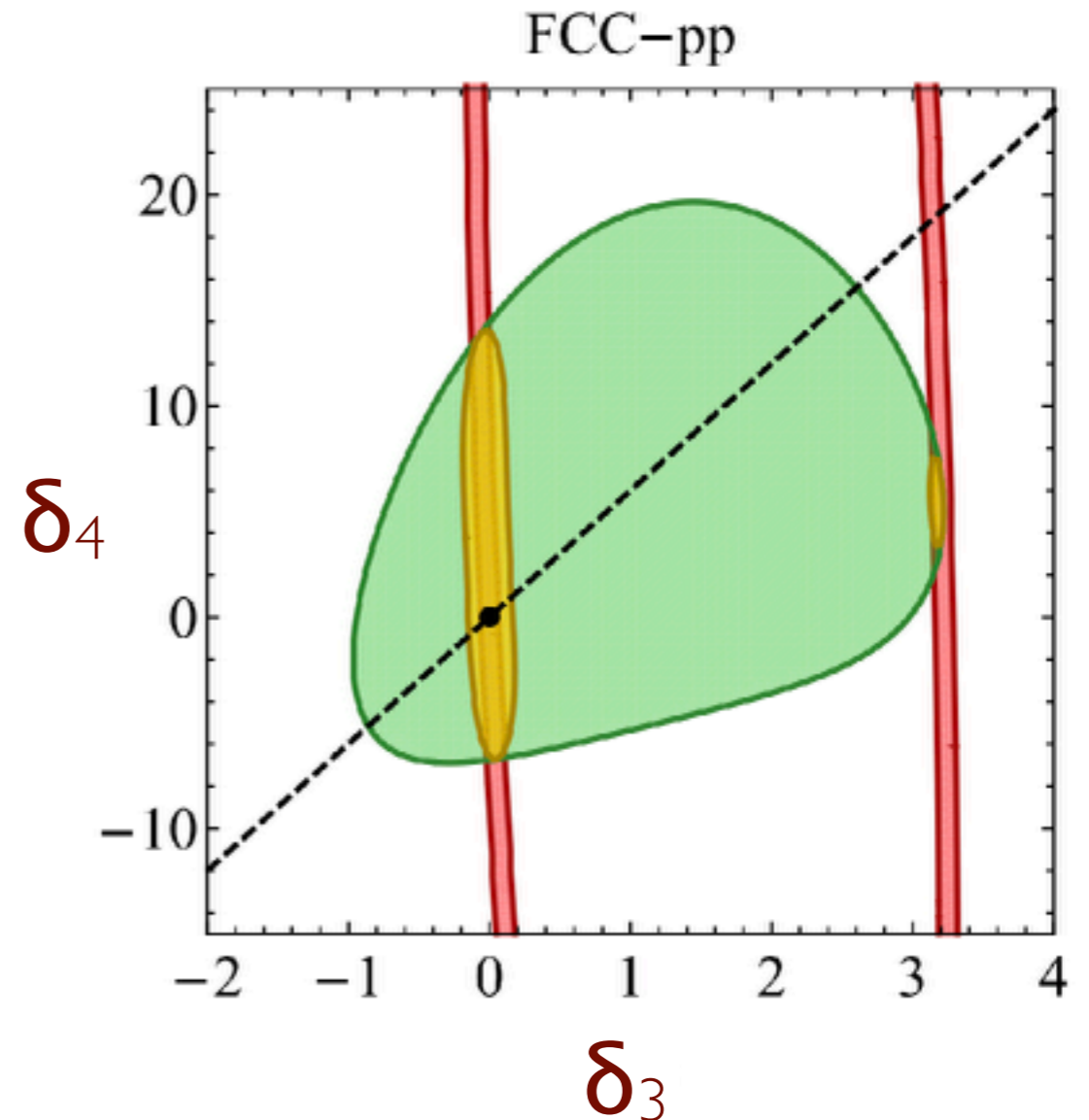
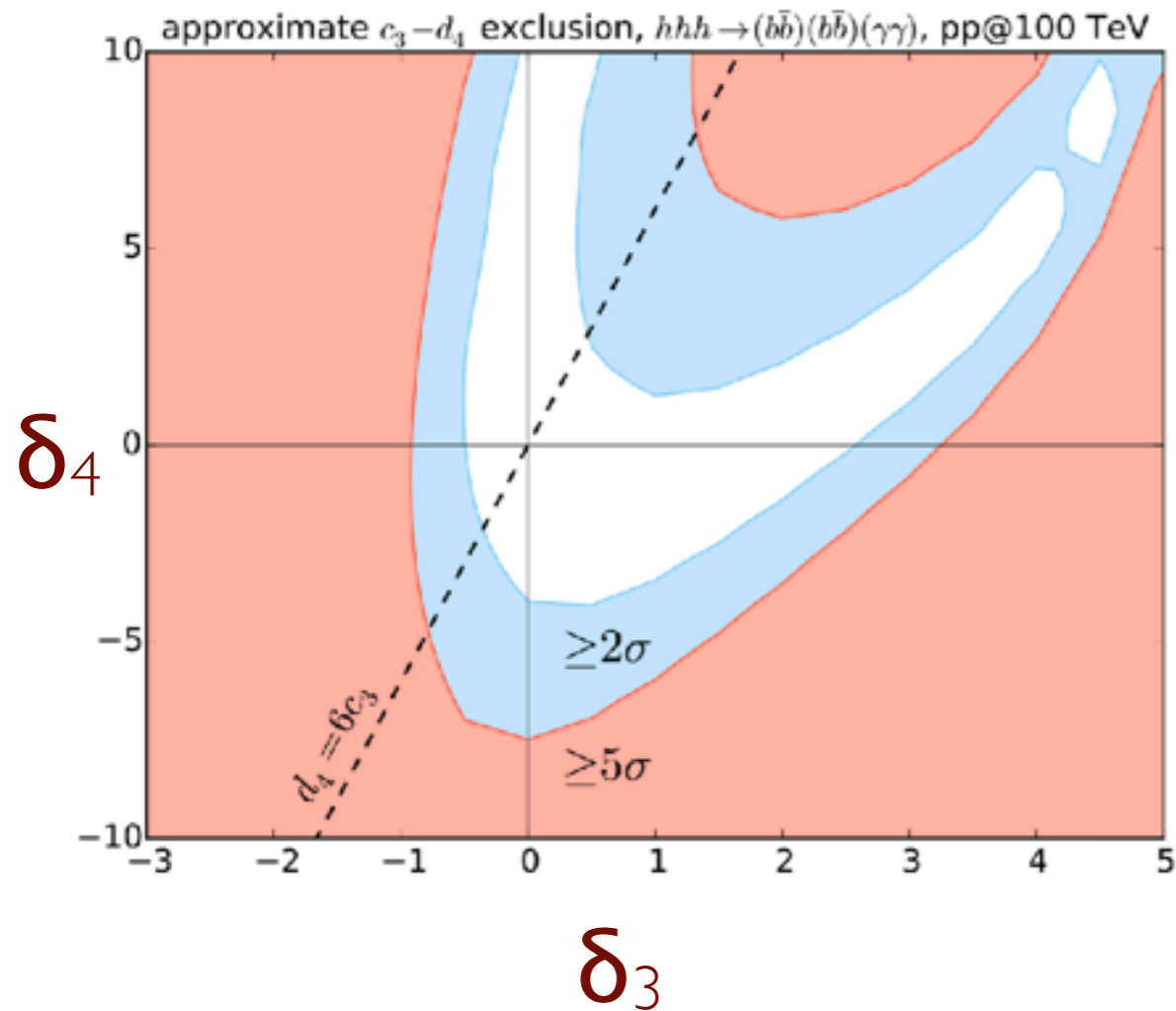


FIG. 6: The approximate expected 2σ (blue) and 5σ (red) exclusion regions on the $c_3 - d_4$ plane after 30 ab^{-1} of integrated luminosity, derived assuming a constant signal efficiency, calculated along the $d_4 = 6c_3$ line in $c_3 \in [-3.0, 4.0]$.

Sensitivity to Higgs quartic is poor even in optimistic cases

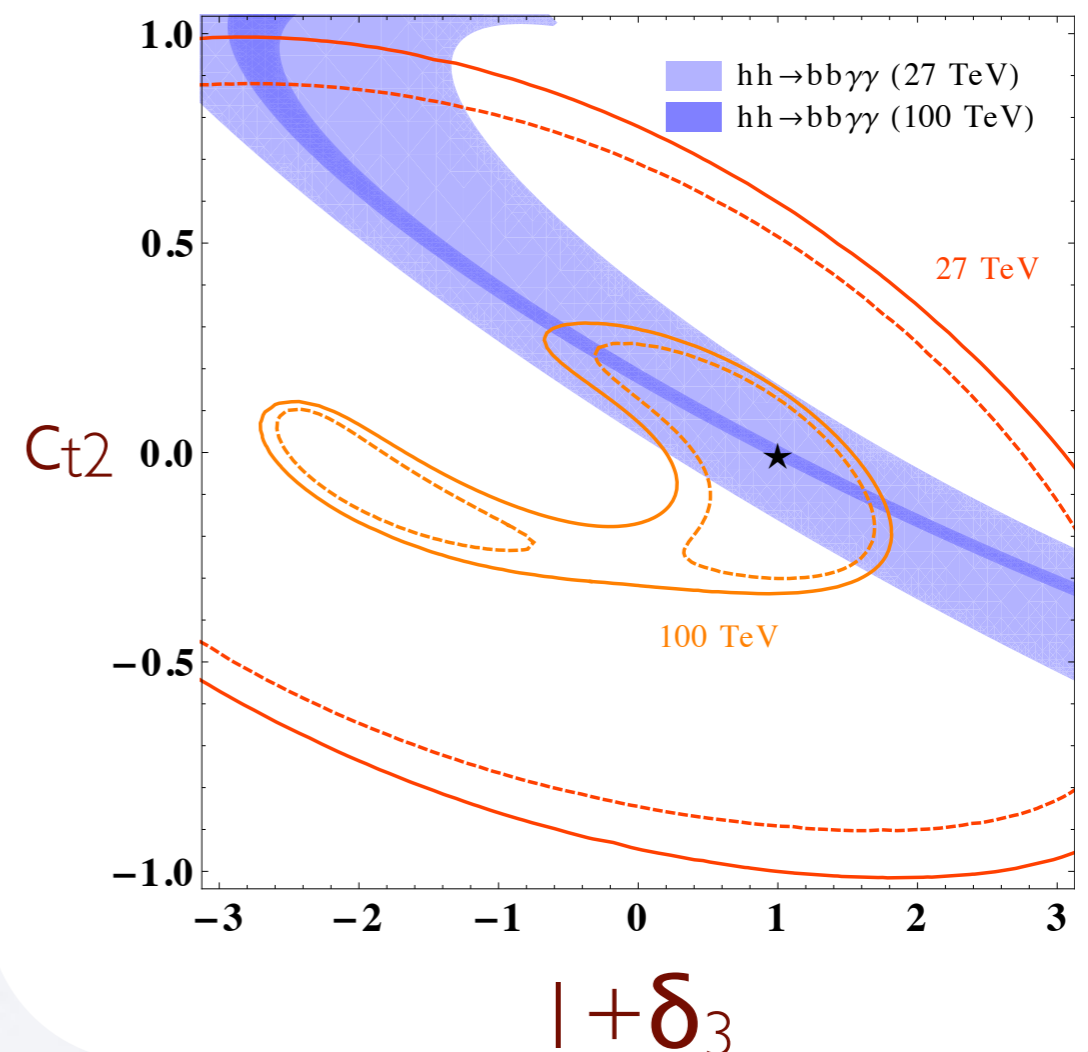
W/Z AND TOP COUPLINGS TO HH

VVhh measured in
VBF DiHiggs to 4b's
Bishara et.al. (1611.03860)

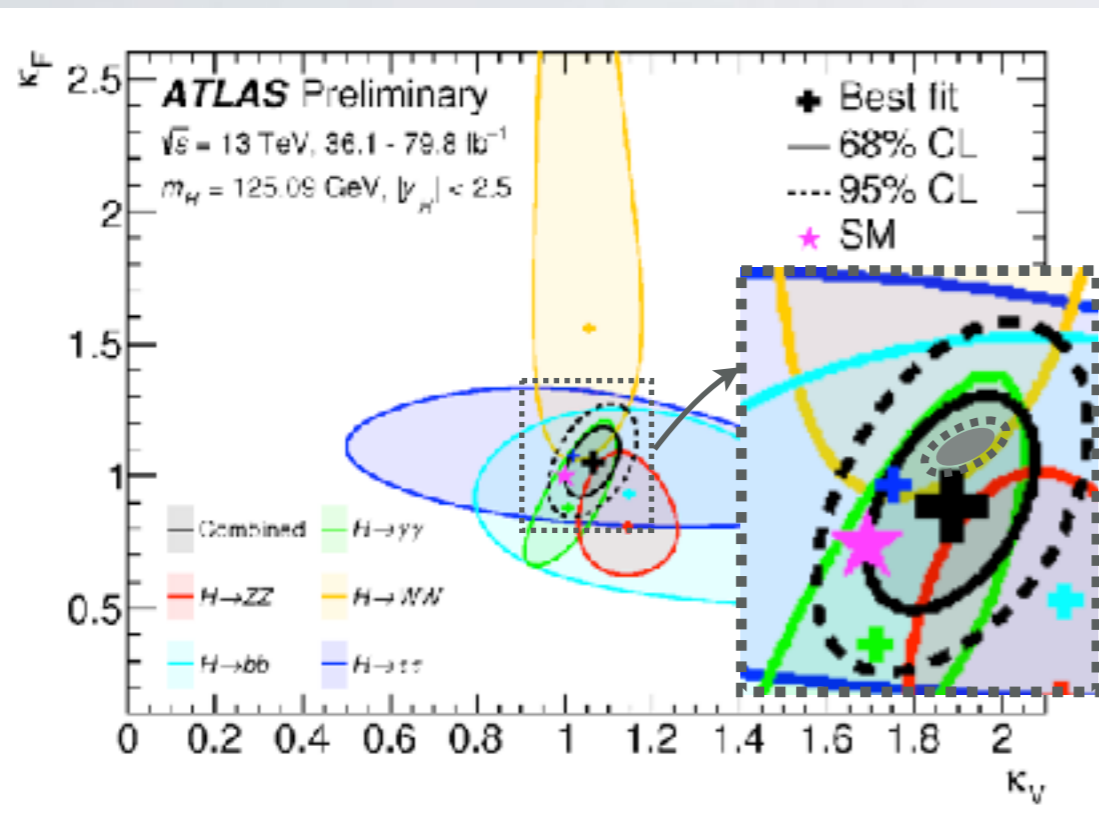
	68% probability interval on δ_{hhVV}	
	$1 \times \sigma_{\text{bkg}}$	$3 \times \sigma_{\text{bkg}}$
LHC ₁₄	[-0.37, 0.45]	[-0.43, 0.48]
HL-LHC	[-0.15, 0.19]	[-0.18, 0.20]
FCC ₁₀₀	[0, 0.01]	[-0.01, 0.01]

Sensitivity to $O(.1-1)$ for
quadratic Higgs couplings

tthh coupling
probed by tthh
production
Li et.al. (1905.03772)



NEW PHYSICS SCALE BOUND FROM UNITARITY VIOLATION

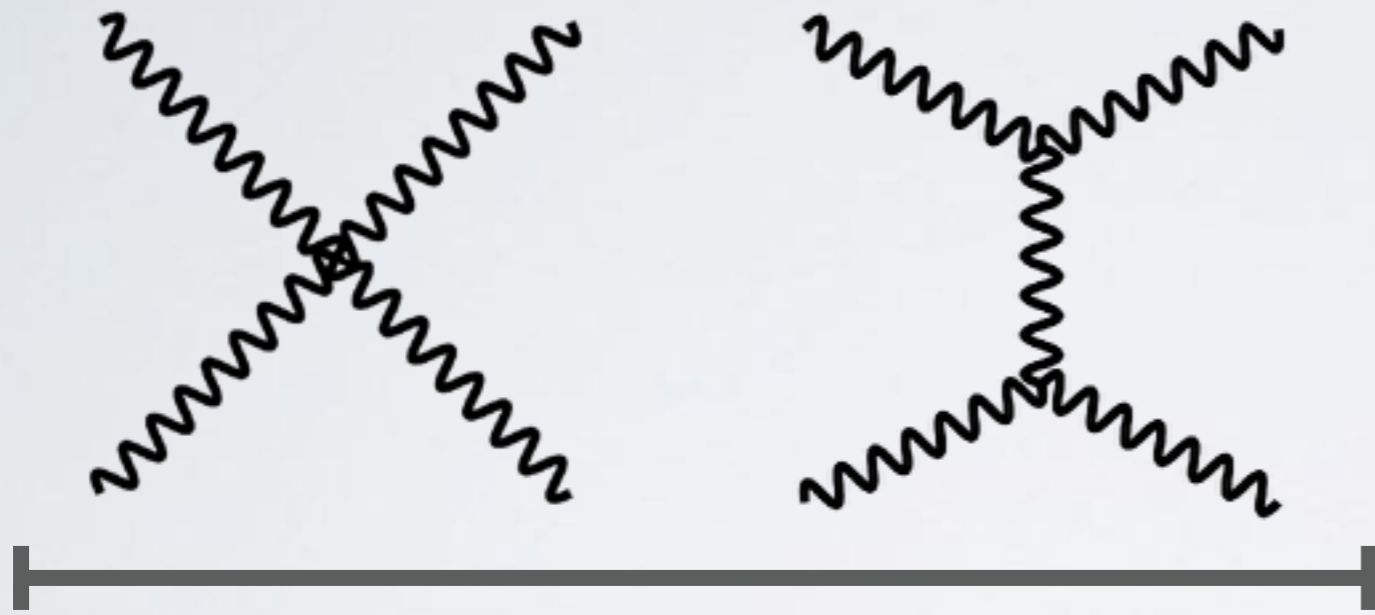


What are the new physics implications of a Higgs coupling deviation?

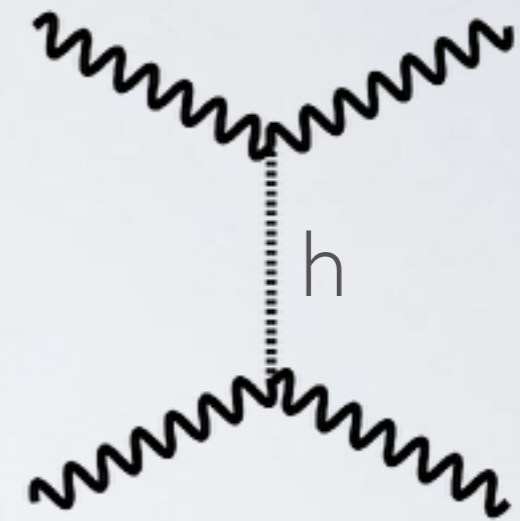
Any Higgs coupling deviation from SM prediction leads to unitarity violation at high energies, placing an upper bound on new physics. Also, leads to interesting processes to measure (see Kilian et.al. 1808.05534, Henning et.al. 1812.09299 & Stolarski, Wu 2006.09374)

CLASSIC EXAMPLE

SCATTERING $Z_L Z_L \Leftrightarrow W^+_L W^-_L$



$$M = c \text{ Energy}^2 + \dots$$



$$M = -c \text{ Energy}^2 + \dots$$

Higgs exchange cancels high energy growth if its couplings are SM-like, matrix element is unitary if $m_H \approx 1 \text{ TeV}$ (Lee, Quigg, Thacker), motivating LHC design

GENERAL HIGGS COUPLINGS

$$\mathcal{L} = \mathcal{L}_{\text{SM}} - \delta_3 \frac{m_h^2}{2v} h^3 - \delta_4 \frac{m_h^2}{8v^2} h^4 - \sum_{n=5}^{\infty} \frac{c_n}{n!} \frac{m_h^2}{v^{n-2}} h^n + \dots \text{ Higgs Potential Couplings}$$

$$+ \delta_{Z1} \frac{m_Z^2}{v} h Z^\mu Z_\mu + \delta_{W1} \frac{2m_W^2}{v} h W^{\mu+} W_\mu^- + \delta_{Z2} \frac{m_Z^2}{2v^2} h^2 Z^\mu Z_\mu + \delta_{W2} \frac{m_W^2}{v^2} h^2 W^{\mu+} W_\mu^-$$

$$+ \sum_{n=3}^{\infty} \left[\frac{c_{Zn}}{n!} \frac{m_Z^2}{v^n} h^n Z^\mu Z_\mu + \frac{c_{Wn}}{n!} \frac{2m_W^2}{v^n} h^n W^{\mu+} W_\mu^- \right] + \dots \text{ W/Z Couplings}$$

$$- \delta_{t1} \frac{m_t}{v} h \bar{t} t - \sum_{n=2}^{\infty} \frac{c_{tn}}{n!} \frac{m_t}{v^n} h^n \bar{t} t + \dots \quad \text{top Couplings}$$

Any nonzero δ or \mathbf{c} coupling is a sign of new physics, which leads to unitarity violation at high energies (higher dim. operators), placing an upper bound on this new physics

ASIDE: TECHNIQUE DETAILS



OUR GENERAL UNITARITY VIOLATION APPROACH

$|P, \alpha\rangle$ Define states of total momentum P
w/ other properties α (e.g. # Higgses)

Properly normalized $\langle P', \alpha' | P, \alpha \rangle = (2\pi)^4 \delta(P - P') \delta_{\alpha\alpha'}$

Leads to unitarity bounds $|T_{\alpha\alpha'}| \leq 1$

$$\langle P', \alpha' | T | P, \alpha \rangle = (2\pi)^4 \delta(P - P') T_{\alpha\alpha'}$$

Allows us to go beyond 2 to 2 processes and set better bounds

EXAMPLE STATES

Only
Scalars

$$|P, k_1, \dots, k_r\rangle = C_{k_1 \dots k_r} \int d^4x e^{-iP \cdot x} \prod_{i=1}^r \left[\phi_i^{(-)}(x) \right]^{k_i} |0\rangle$$

$$\frac{1}{|C_{k_1 \dots k_r}|^2} = \frac{\prod_i k_i!}{8\pi (\sum_i k_i - 1)! (\sum_i k_i - 2)!} \left(\frac{E}{4\pi} \right)^{2 \sum_i k_i - 4}$$

Two
Fermions

$$|P; k_1, \dots, k_r, L/R\rangle \equiv C_k'' \int d^4x e^{-iP \cdot x} \phi_1^{(-)}(x)^{k_1} \dots \phi_r^{(-)}(x)^{k_r} \bar{\psi}_{R/L}^{a(-)}(x) \psi_{L/R}^{a(-)}(x) |0\rangle$$

$$\frac{1}{|C_k''|^2} = \frac{2N_c E^2 \prod_i k_i!}{8\pi (\sum_i k_i + 1)! (\sum_i k_i)!} \left(\frac{E}{4\pi} \right)^{2 \sum_i k_i}$$

EQUIVALENCE THEOREM

For general h couplings, restore $SU(2) \times U(1)$ invariance, by introducing Goldstone bosons to use equivalence theorem for W_L, Z_L amplitudes

Higgs
self-interactions

$$\begin{aligned} X &\equiv \sqrt{2|H|^2} - v = \sqrt{(v+h)^2 + \vec{G}^2} - v \\ &= h + \frac{1}{2v} \vec{G}^2 - \frac{1}{2v^2} h \vec{G}^2 + \dots \end{aligned}$$

For W/Z and
top interactions

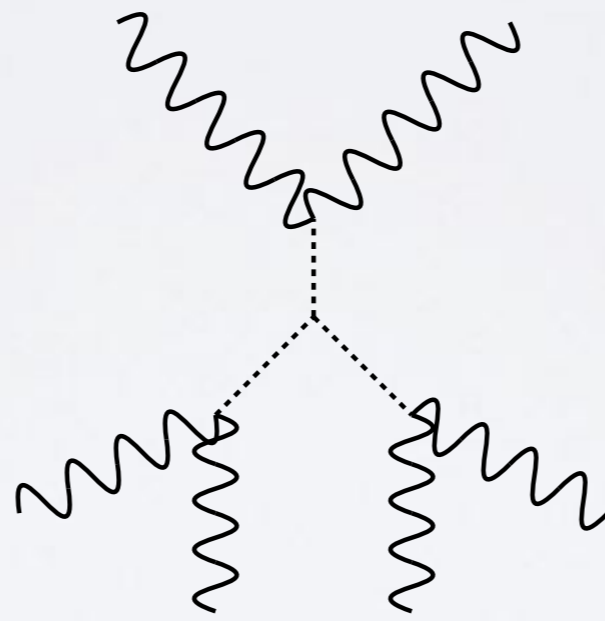
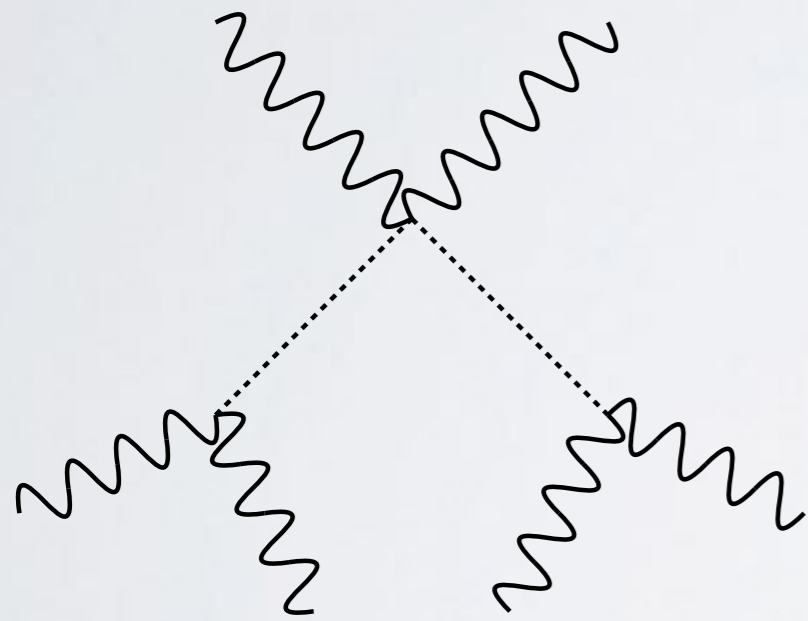
$$P = \frac{H}{\sqrt{|H|^2}} = \begin{pmatrix} 0 \\ 1 \end{pmatrix} + \begin{pmatrix} \sqrt{2}G^+/v \\ iG^0/v \end{pmatrix} + \dots$$

END ASIDE



EXAMPLE: TRILINEAR UNITARITY VIOLATION

Modifying trilinear from SM value automatically leads to Unitarity violation at high energies



Example:

$$Z_L Z_L Z_L \Leftrightarrow Z_L Z_L Z_L$$

Cancellation to get
 $M \sim 1/\text{Energy}^2$
requires SM
trilinear value!

MODEL DEPENDENCE OF INTERACTIONS

$$X^3 \sim h^3 + \vec{G}^2(\boxed{h^2} + h^3 + \dots) + \vec{G}^4(\boxed{h} + h^2 + \dots) + \boxed{\vec{G}^6}(1 + h + \dots) + \vec{G}^8(1 + h + \dots) + \vec{G}^{10}(1 + h + \dots) + \dots,$$

$$X^4 \sim h^4 + \vec{G}^2(h^3 + h^4 + \dots) + \vec{G}^4(h^2 + h^3 + \dots) + \vec{G}^6(h + h^2 + \dots) + \vec{G}^8(1 + h + \dots) + \vec{G}^{10}(1 + h + \dots) + \dots,$$

$$X^5 \sim h^5 + \vec{G}^2(h^4 + h^5 + \dots) + \vec{G}^4(h^3 + h^4 + \dots) + \vec{G}^6(h^2 + h + \dots) + \vec{G}^8(h + h^2 + \dots) + \vec{G}^{10}(1 + h + \dots) + \dots,$$

(Schematic without coefficients, but we know cancellations can occur due to SMEFT description)

Terms circled can only come from trilinear!

BEST CHANNELS FOR HIGGS TRILINEAR

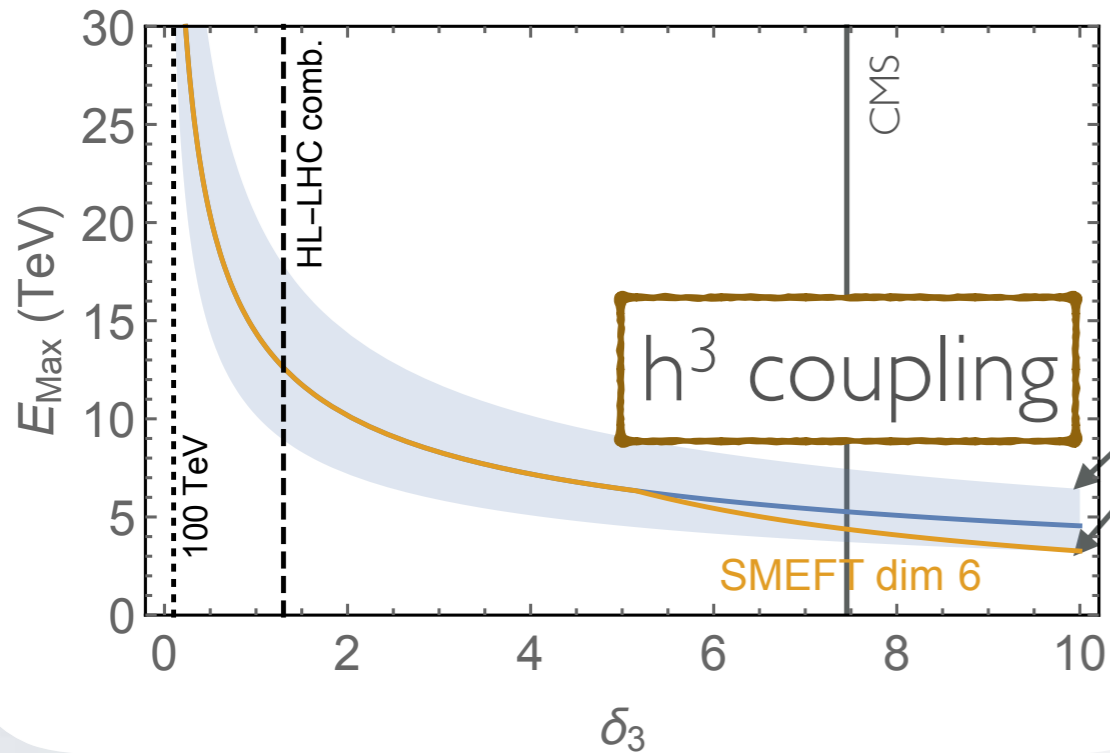
$$hW_L^+W_L^- \rightarrow W_L^+W_L^- : E_{max} = \frac{9.4 \text{ TeV}}{\left|\frac{\delta_3}{7.5}\right|}$$

$$W_L^+W_L^+W_L^- \rightarrow W_L^+W_L^+W_L^- : E_{max} = \frac{5.2 \text{ TeV}}{\sqrt{\left|\frac{\delta_3}{7.5}\right|}}$$

(Normalized to largest deviation consistent with ATLAS and CMS di-Higgs 95%CL constraints)

Takeaway: Current constraints still allow low unitarity bound w/ nearby new physics, a measured coupling deviation from SM places an upper bound on new physics

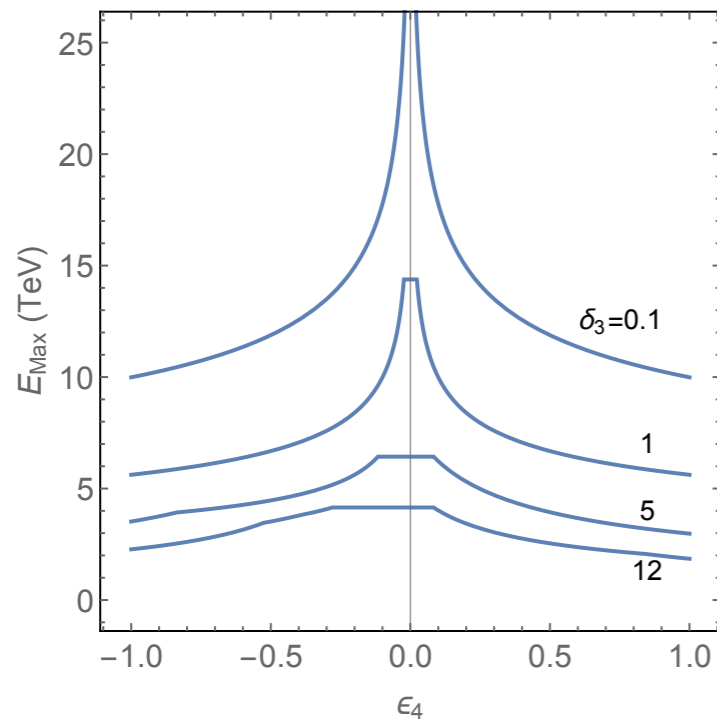
HIGGS TRILINEAR COUPLING DEVIATION



Estimated

Theoretical uncertainty of
Unitarity violating scale

**Current bound allows new
physics below ~ 5 TeV**



**Unitarity requires higher order
couplings to be correlated**

Quartic deviation must satisfy SMEFT-like
relation, $\delta_4 = 6\delta_3(1 + \epsilon_4)$ as predicted by $|H|^6$,
to keep new physics above 10 TeV

W/Z, TOP COUPLINGS

Processes that
only depend on
 hWW, hZZ
couplings

$$W_L^+ W_L^+ \rightarrow W_L^+ W_L^+ : E_{\max} \simeq \frac{1.2 \text{ TeV}}{|\delta_{W1}|^{1/2}},$$

$$Z_L Z_L \rightarrow W_L^+ W_L^- : E_{\max} \simeq \frac{1.5 \text{ TeV}}{|\delta_{Z1} + \delta_{W1}|^{1/2}},$$

$$W_L^+ h \rightarrow W_L^+ Z_L : E_{\max} \simeq \frac{1.0 \text{ TeV}}{|\delta_{Z1} - \delta_{W1}|^{1/2}},$$

$$W_L^+ W_L^+ W_L^- \rightarrow W_L^+ Z_L : E_{\max} \simeq \frac{1.5 \text{ TeV}}{|\delta_{Z1} - \delta_{W1}|^{1/3}}.$$

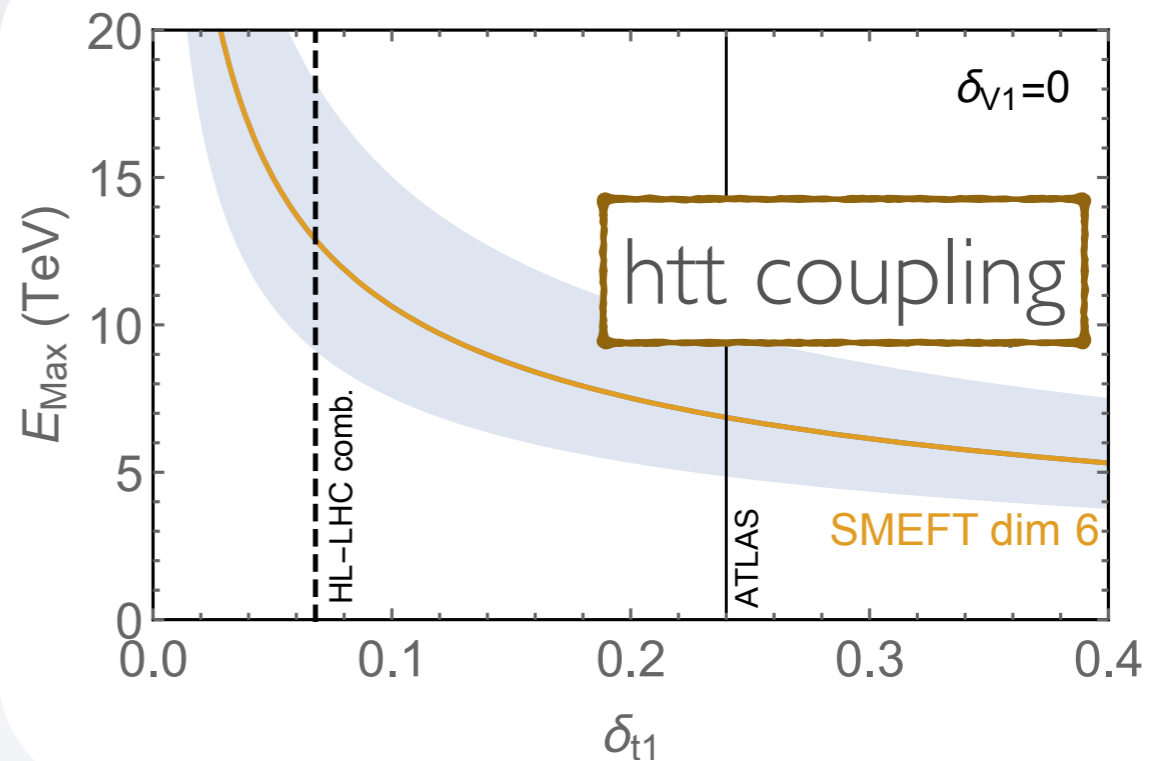
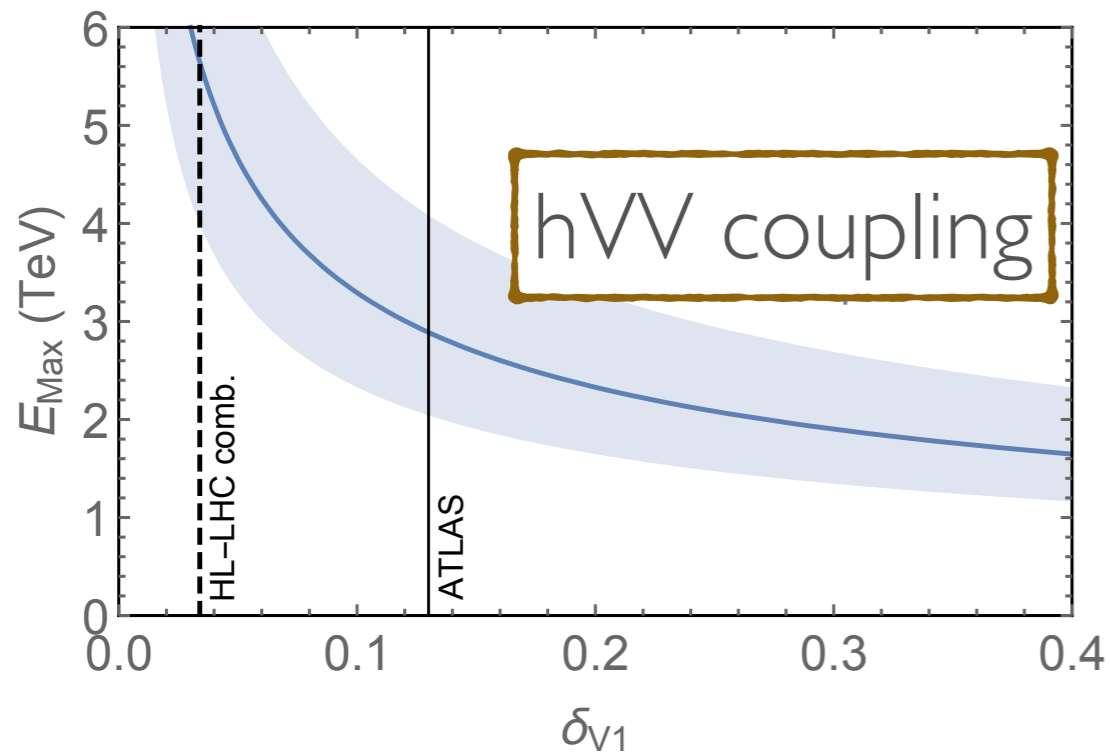
Processes that
only depend on
 htt (and hVV)
coupling

$$t_R \bar{t}_R \rightarrow W_L^+ W_L^- : E_{\max} \simeq \frac{5.1 \text{ TeV}}{|\delta_{t1} + \delta_{V1}|},$$

$$t_R \bar{b}_R \rightarrow W_L^+ h : E_{\max} \simeq \frac{3.6 \text{ TeV}}{|\delta_{t1} - \delta_{V1}|}$$

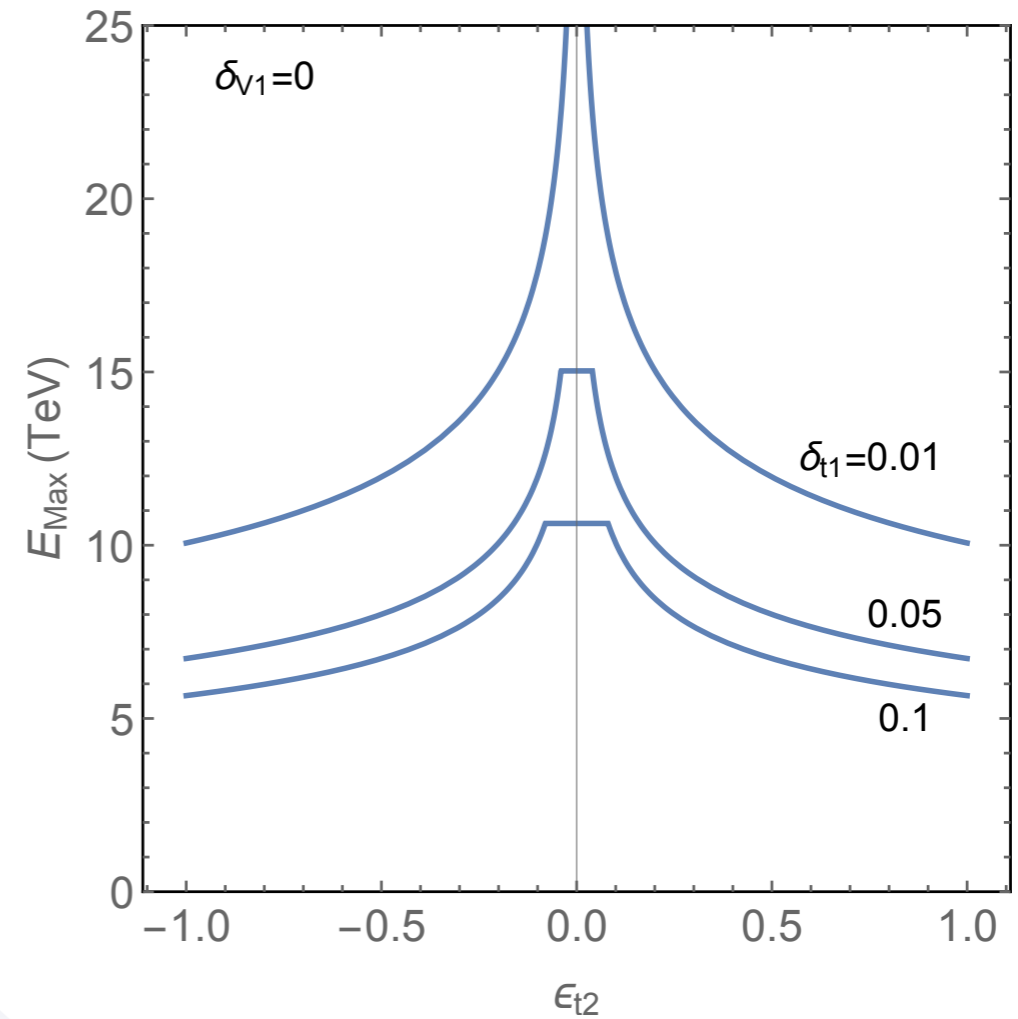
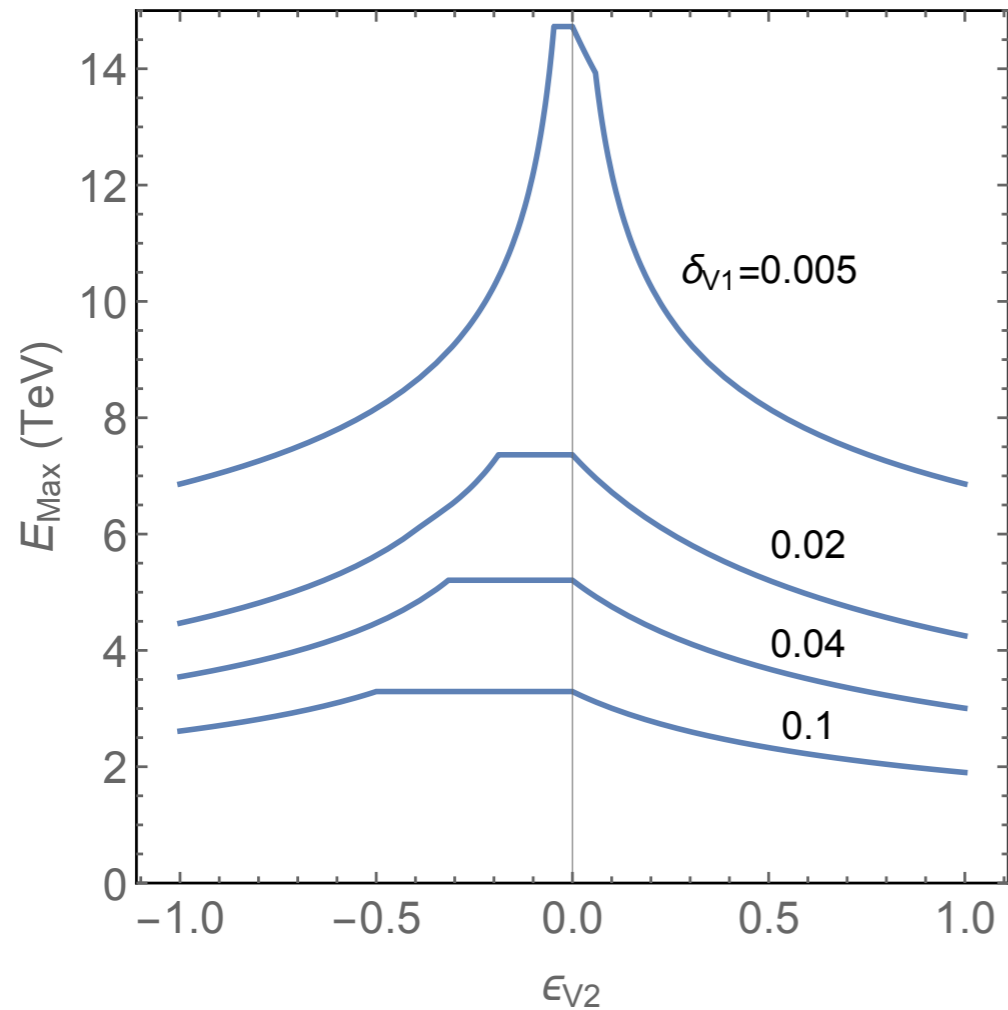
$$t_R \bar{b}_R \rightarrow W_L^+ W_L^+ W_L^- : E_{\max} \simeq \frac{3.3 \text{ TeV}}{\sqrt{|\delta_{t1} - \frac{1}{3}\delta_{V1}|}},$$

W/Z AND TOP BOUNDS



Existing strong bounds on these couplings still allow future deviations where new physics has to appear below $\sim 3\text{-}8$ TeV. In fact, hVV is more powerful than h^3 !

SMEFT CONSTRAINTS



$$|H|^2 |D_\mu H|^2$$

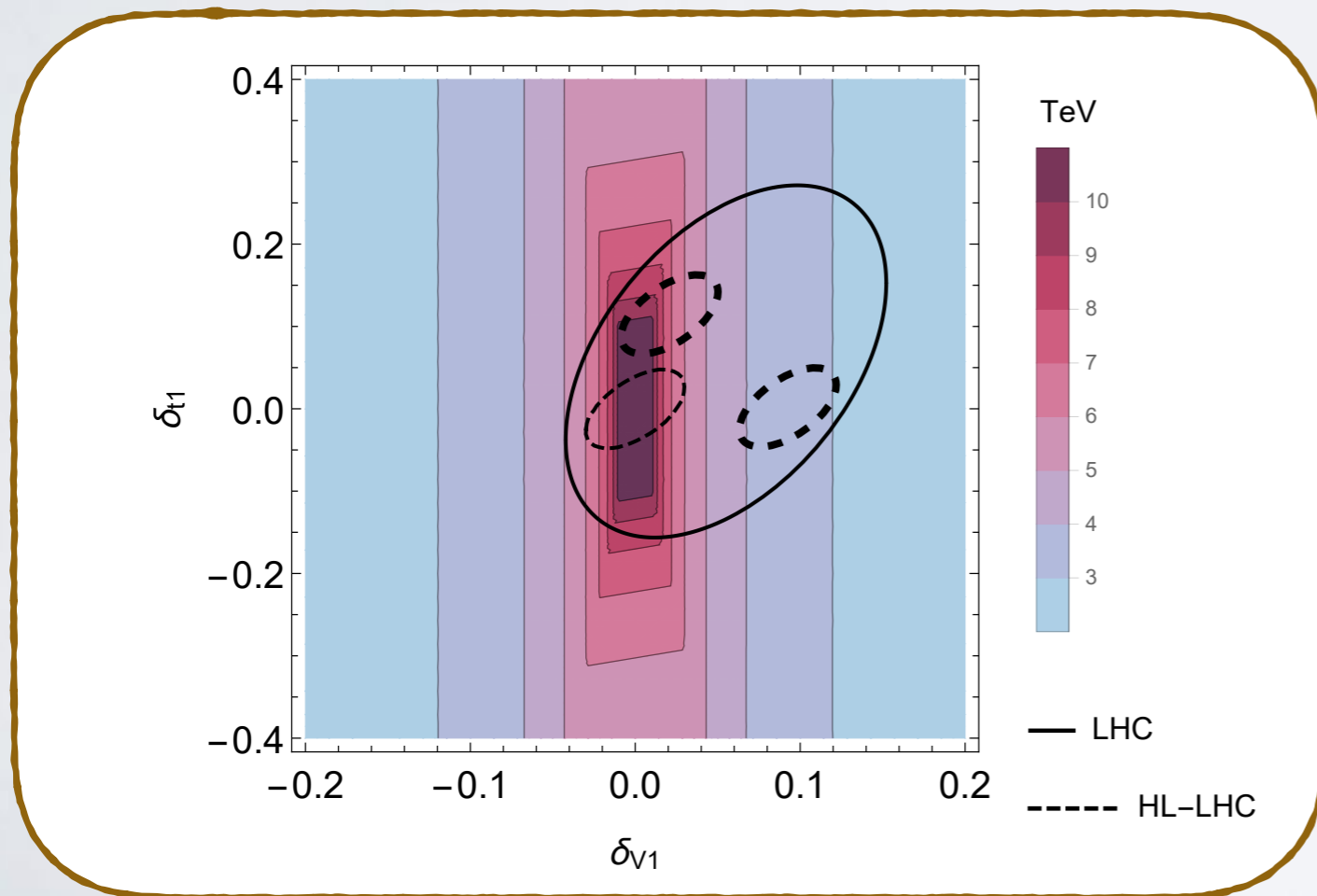
$$|H|^2 \bar{Q} H P_L Q + h.c.$$

Again, for new physics to be much higher than a TeV, need a SMEFT-like structure

WHAT DOES PRECISION BUY?

kappa-0	HL-LHC	LHeC	HE-LHC		ILC			CLIC			CEPC	FCC-ee		FCC-ee/eh/hh
			S2	S2'	250	500	1000	380	15000	3000		240	365	
κ_W [%]	1.7	0.75	1.4	0.98	1.8	0.29	0.24	0.86	0.16	0.11	1.3	1.3	0.43	0.14
κ_Z [%]	1.5	1.2	1.3	0.9	0.29	0.23	0.22	0.5	0.26	0.23	0.14	0.20	0.17	0.12
κ_g [%]	2.3	3.6	1.9	1.2	2.3	0.97	0.66	2.5	1.3	0.9	1.5	1.7	1.0	0.49
κ_γ [%]	1.9	7.6	1.6	1.2	6.7	3.4	1.9	98*	5.0	2.2	3.7	4.7	3.9	0.29
$\kappa_{Z\gamma}$ [%]	10.	—	5.7	3.8	99*	86*	85*	120*	15	6.9	8.2	81*	75*	0.69
κ_c [%]	—	4.1	—	—	2.5	1.3	0.9	4.3	1.8	1.4	2.2	1.8	1.3	0.95
κ_t [%]	3.3	—	2.8	1.7	—	6.9	1.6	—	—	2.7	—	—	—	1.0
κ_b [%]	3.6	2.1	3.2	2.3	1.8	0.58	0.48	1.9	0.46	0.37	1.2	1.3	0.67	0.43
κ_μ [%]	4.6	—	2.5	1.7	15	9.4	6.2	320*	13	5.8	8.9	10	8.9	0.41
κ_τ [%]	1.9	3.3	1.5	1.1	1.9	0.70	0.57	3.0	1.3	0.88	1.3	1.4	0.73	0.44

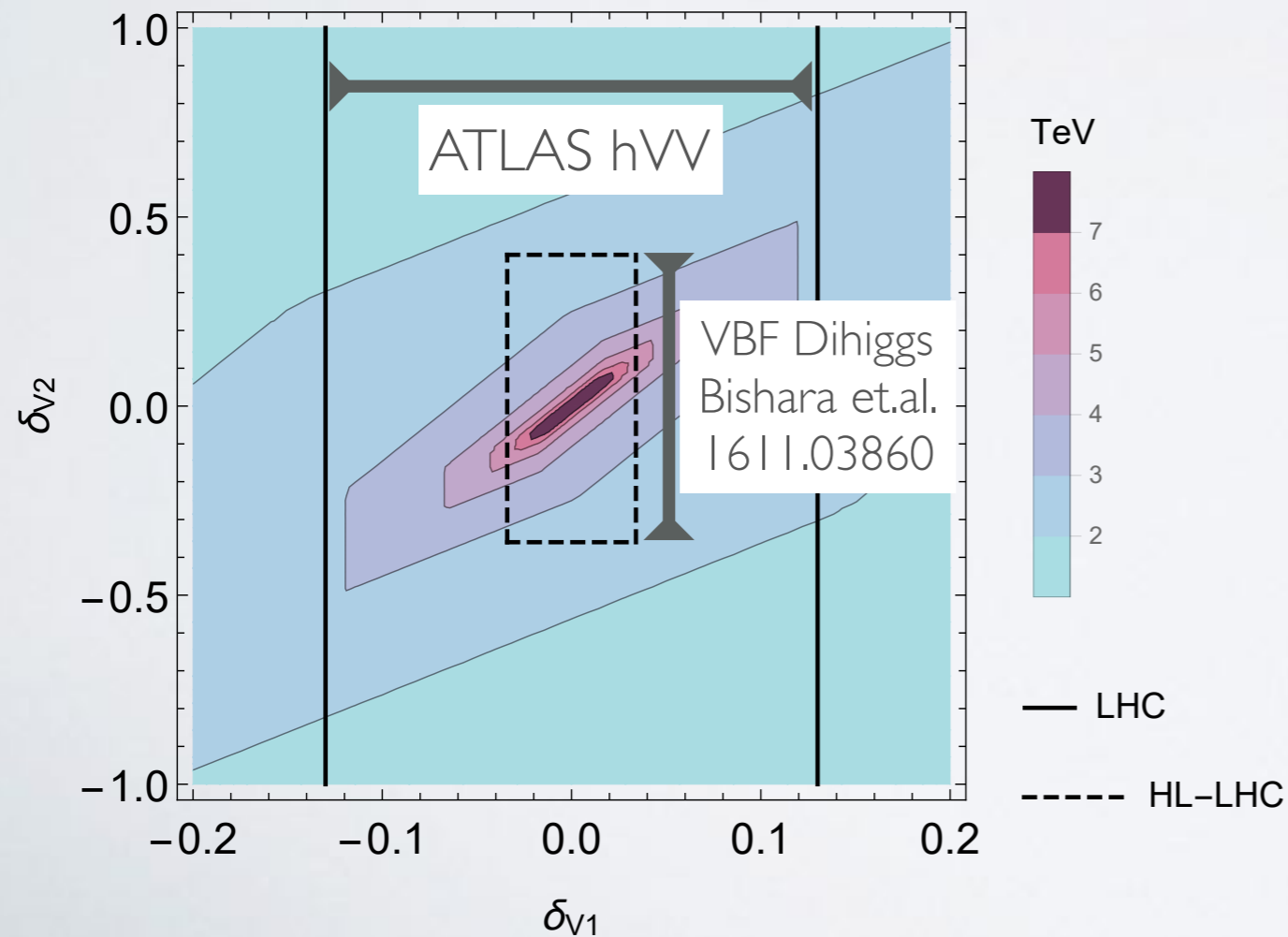
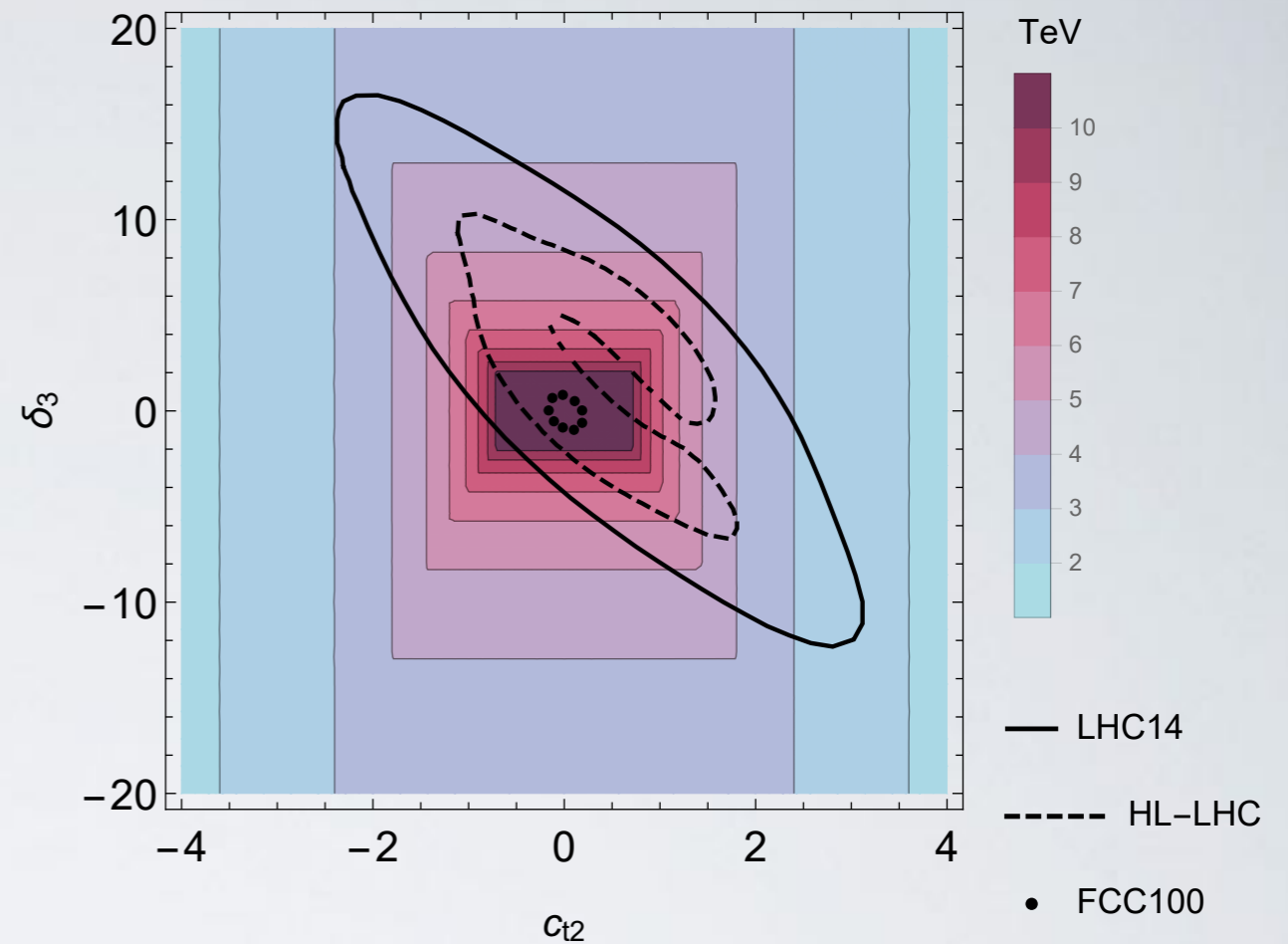
Higgs@FutureColliders
report
(1905.03764)



Unitarity bounds
gives a quantitative reason
to improve precision
outside of better
precision

DI-HIGGS INTERPLAY

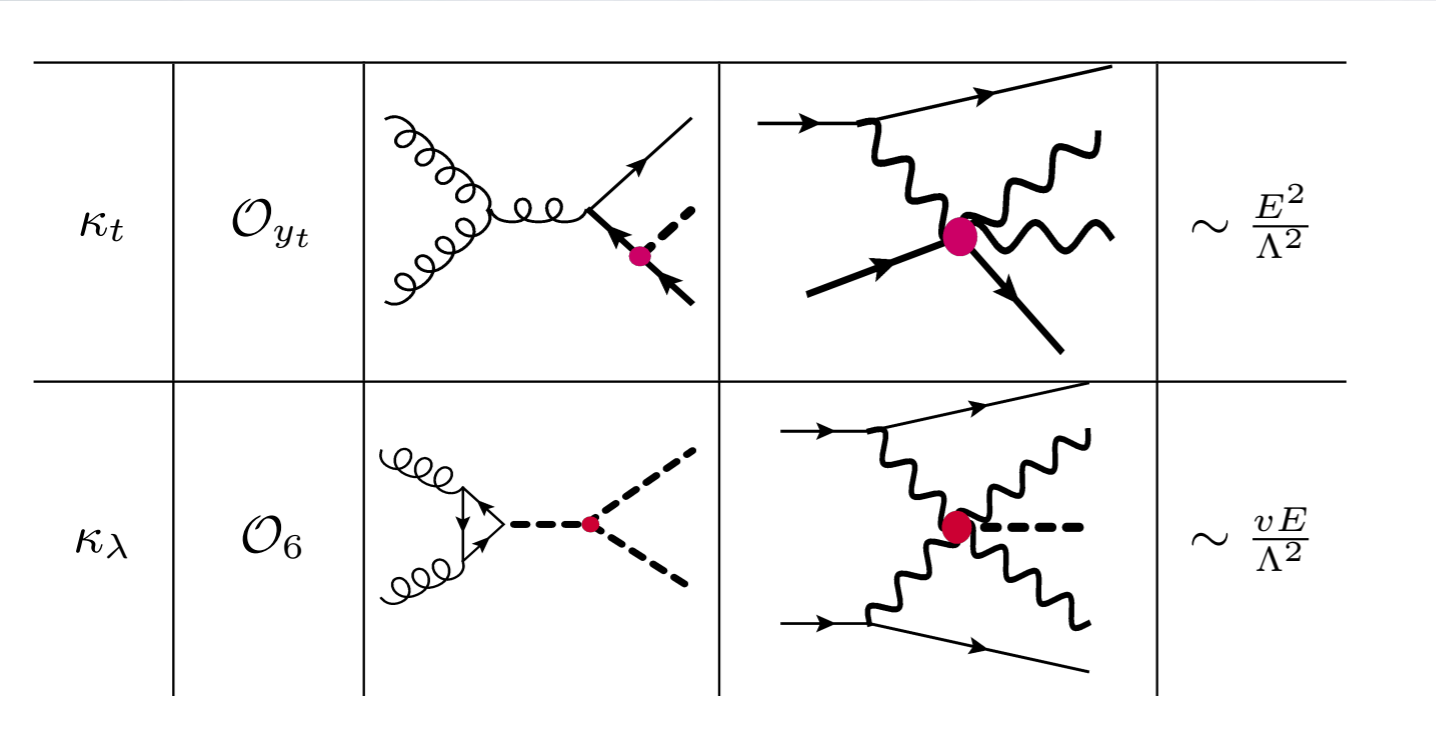
Projected contours from
Azatov et.al. 1502.00539
on di-Higgs from gluon fusion



Vector boson di-Higgs
production constrains
hhVV

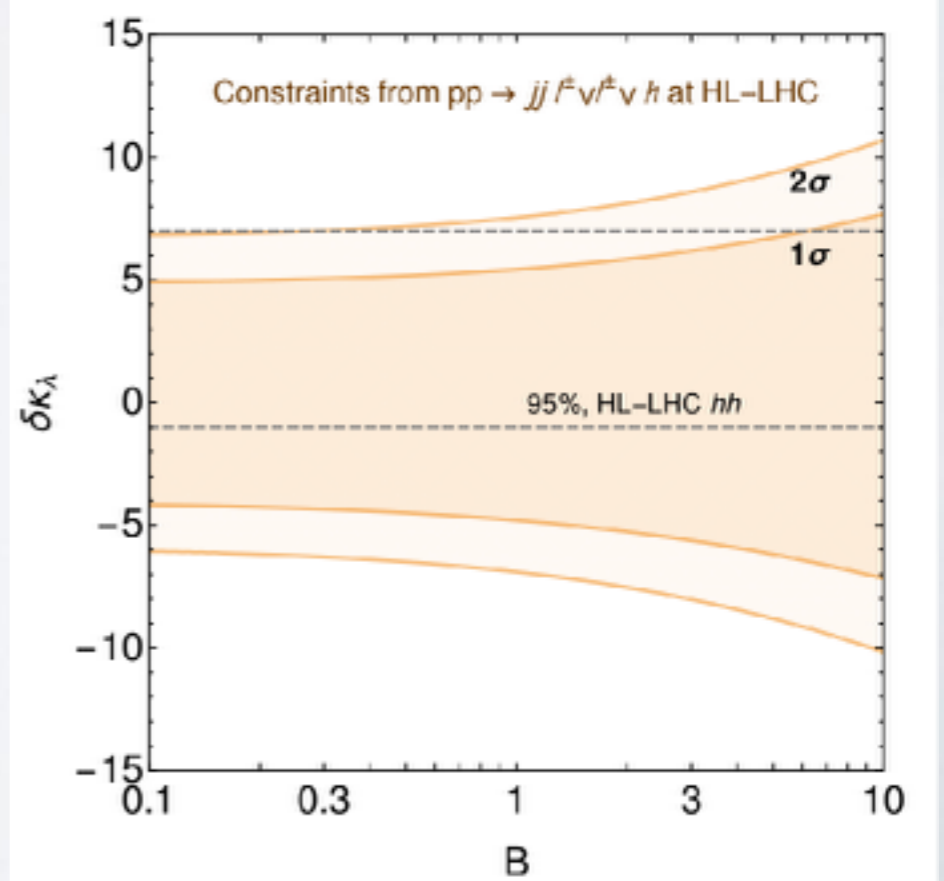
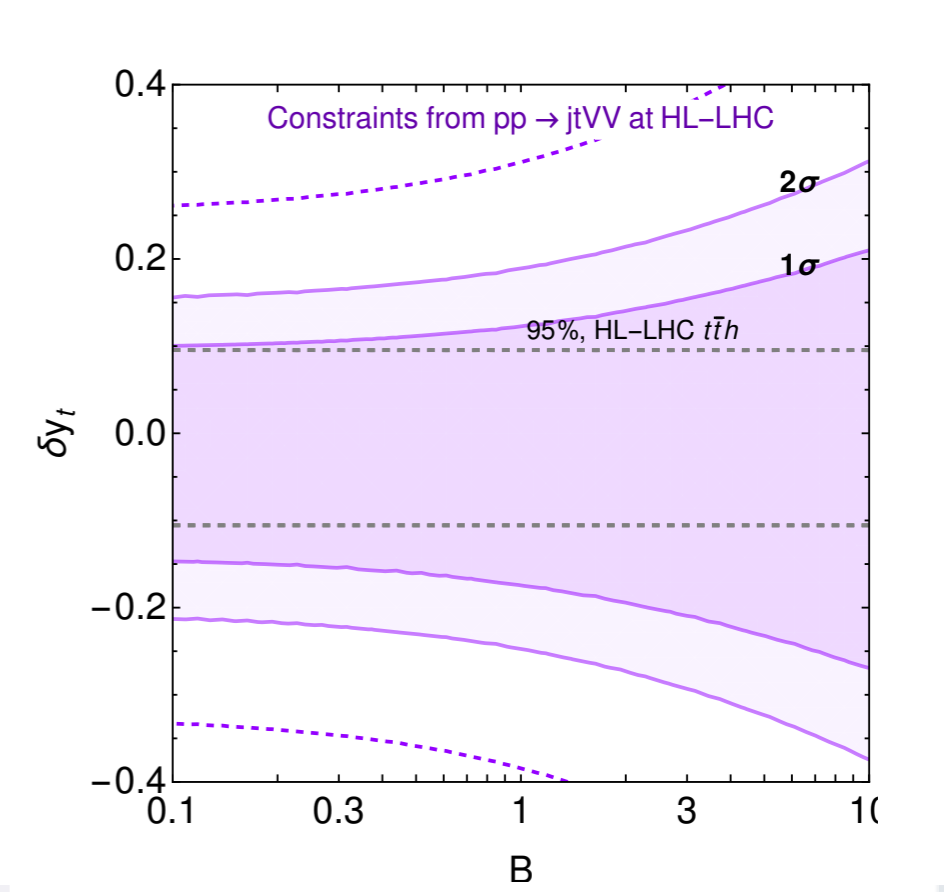
COLLIDER TESTS OF UNITARITY VIOLATION

Henning et.al. | 8 | 2.09299
 See also Kilian et.al. | 808.05534,
 & Stolarski, Wu 2006.09374



Searching for Unitarity violating processes (solid) has similar sensitivities to coupling measurement (dashed) for $t\bar{t}h$, $h\bar{h}h$

Extension to $t\bar{t}h\bar{h}$ and $VVh\bar{h}$?



CONCLUSIONS

- Precision Higgs couplings can discover a SM deviation; unitarity violation gives quantitative connection btw coupling deviations and bound on new physics
- Higgs self-couplings, hVV , htt current bounds allow new physics at LHC energies and future sensitivities can still place bounds below 10 TeV (with different sensitivities)
- Higher order couplings (e.g. $hhhh$, $hhVV$, $hhtt$) are SMEFT-like if new physics scale is well above TeV scale. If no new physics accompanies coupling deviation, evidence for SMEFT-like structure

CONCLUSIONS (CONT.)

- Di-Higgs searches test hhh , $hhVV$, $hh\tau\tau$ couplings, with interesting interplay for new physics bounds
- Future direction 1: Can we develop no lose theorems for the new physics accompanying a coupling deviations?
- Future direction 2: Are these amplitudes useful beyond unitarity violation? Are there better/stronger unitarity amplitudes?

THANK YOU

EXTRA SLIDES

V AND TOP COUPLINGS

Use a nonanalytic
Higgs doublet

$$P = \frac{H}{\sqrt{|H|^2}} = \begin{pmatrix} 0 \\ 1 \end{pmatrix} + \begin{pmatrix} \sqrt{2}G^+/v \\ iG^0/v \end{pmatrix} + \dots$$

$$\left(m_W^2 W^2 + \frac{1}{2} m_Z^2 Z^2 \right) \left(1 + 2(1 + \delta_{hVV}) \frac{h}{v} + (1 + \delta_{hhVV}) \frac{h^2}{v^2} + c_3 \frac{h^3}{v^3} \right)$$

$$\rightarrow |DP|^2 \left(\delta_{hVV} v X + \frac{1}{2} \delta_{hhVV} X^2 + \frac{c_3}{2v} X^3 \right)$$

$$-m_t \bar{T} T \left[1 + (1 + \delta\kappa_t) \frac{h}{v} + \frac{1}{2} c_2 \left(\frac{h}{v} \right)^2 + \frac{1}{6} c_3 \left(\frac{h}{v} \right)^3 \right]$$

$$\rightarrow -m_t \bar{T}_R P \epsilon \begin{pmatrix} T_L \\ B_L \end{pmatrix} \left[\delta\kappa_t \frac{X}{v} + \frac{1}{2} c_2 \left(\frac{X}{v} \right)^2 + \frac{1}{6} c_3 \left(\frac{X}{v} \right)^3 + \dots \right] + h.c.$$