

Gravitational Waves

from the

EW Phase Transition

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22/10/20

The EW Phase Transition

WHY?

→ Yield Precise Understanding of EWSB in Early Universe

The EW Phase Transition

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→ Yield Precise Understanding of EWSB in Early Universe

→ (Possible) Answer to Open Mysteries at Interface of Particle Physics & Cosmology

Origin of Matter-Antimatter Asymmetry

▶ EW-scale Baryogenesis

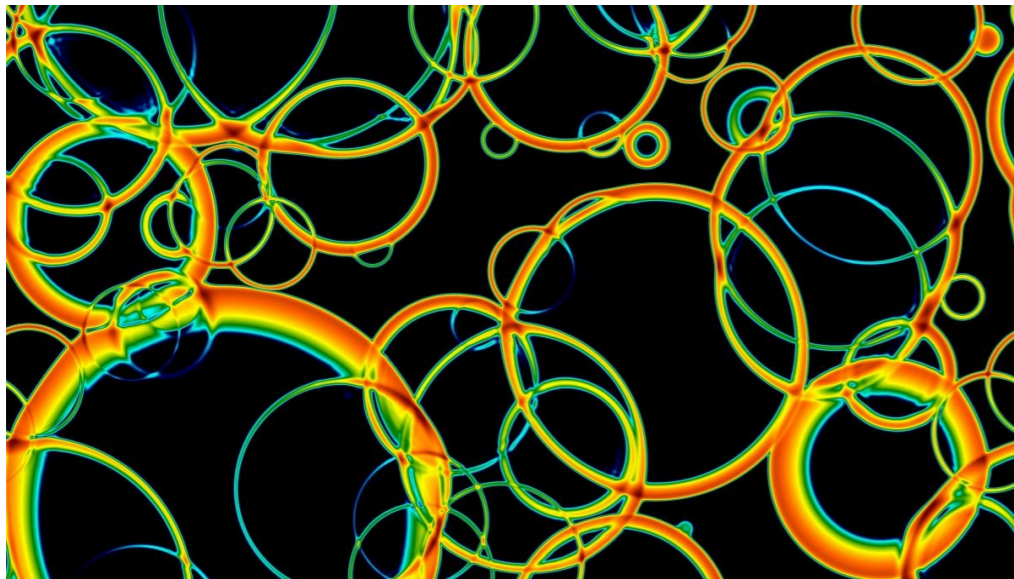
The EW Phase Transition

WHY?

- Yield Precise Understanding of EWSB in Early Universe
- (Possible) Answer to Open Mysteries at Interface of Particle Physics & Cosmology
- **(Possible) Cosmological Relics from the EW Epoch**

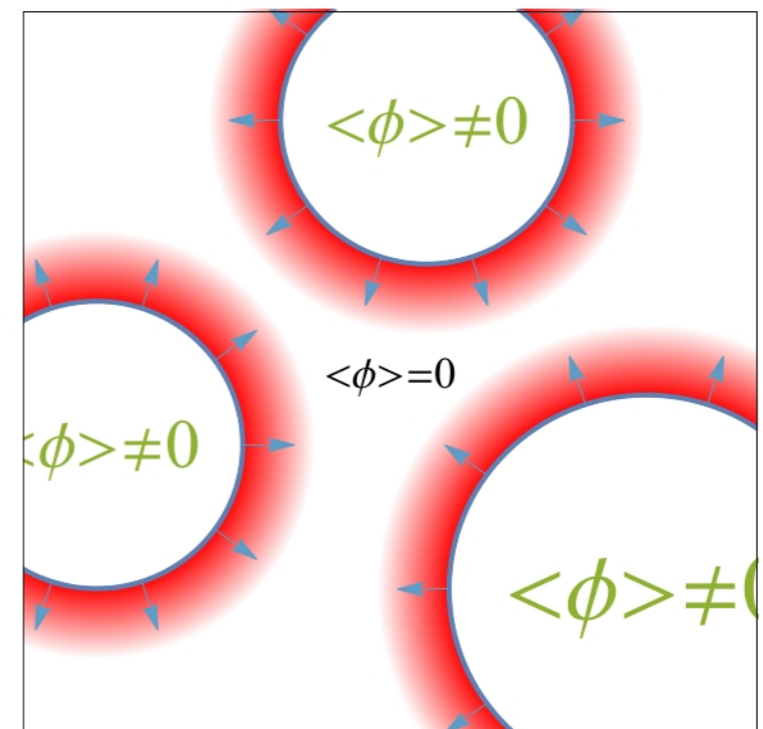
Gravitational Wave Signal

Sourced by Collisions of Higgs bubbles from a first order EW phase transition & subsequent plasma motions



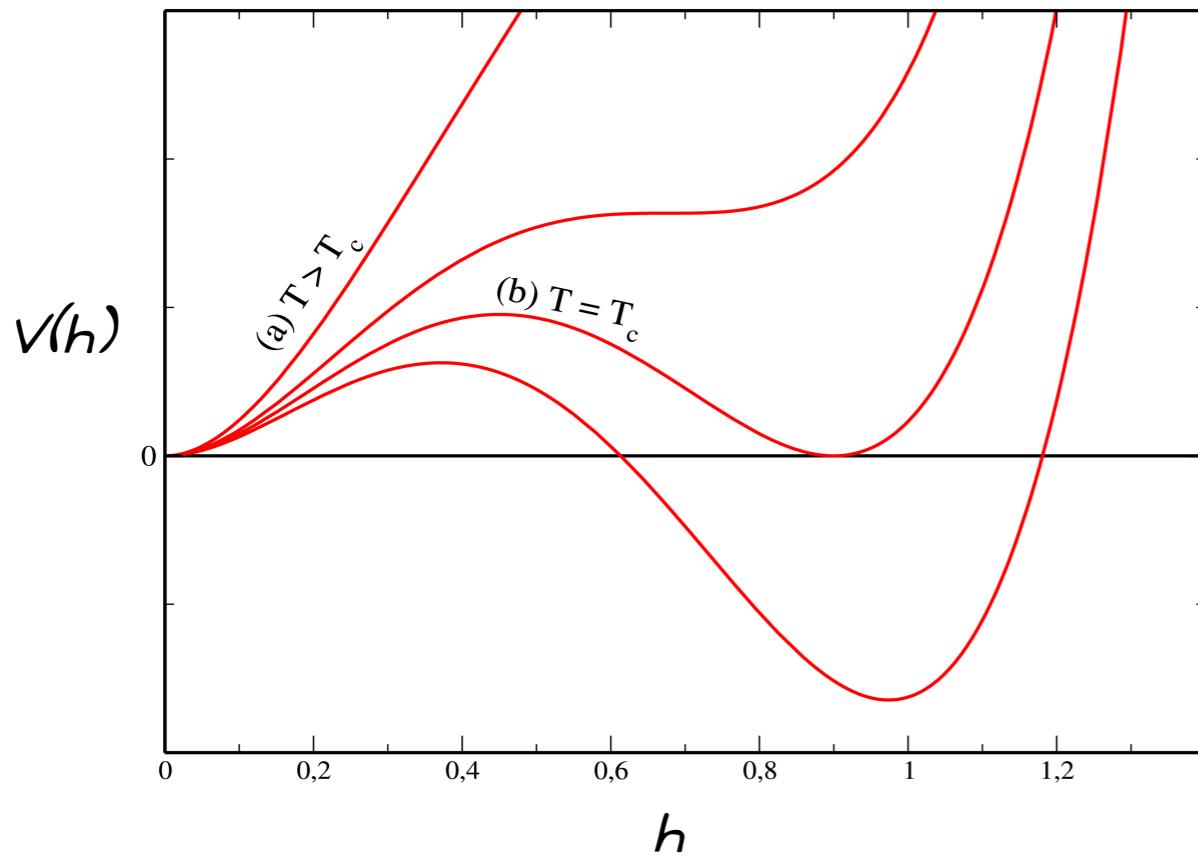
Courtesy of D. Weir (Helsinki)

Hindmarsh, Huber, Rummukainen, Weir, PRD **92** (2015) 123009



Assume a 1st Order EW Phase Transition...

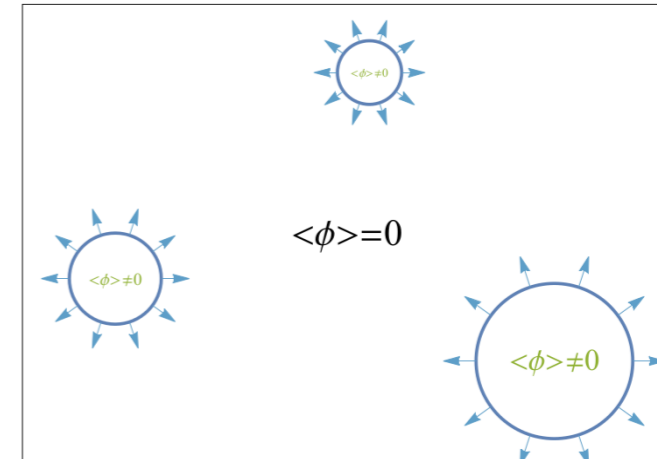
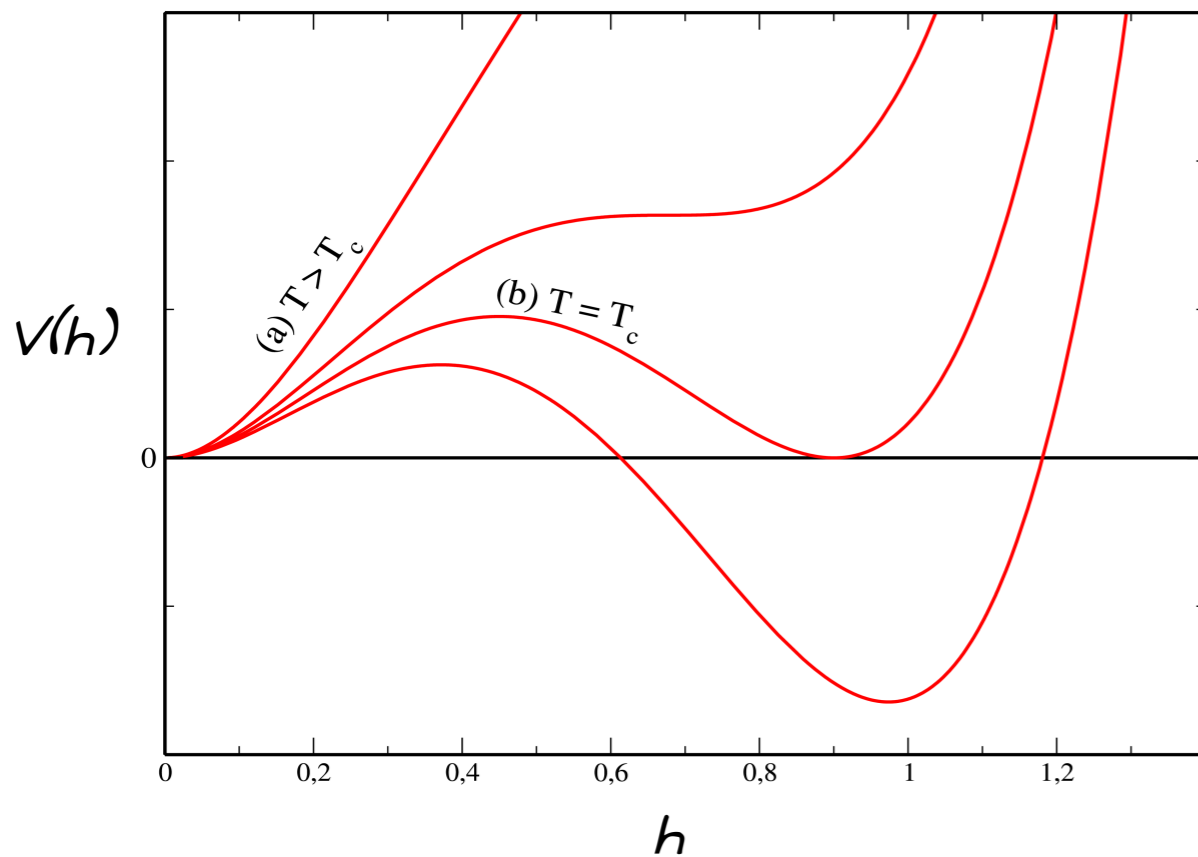
Higgs Effective Potential (finite T)



- Phases separated by potential barrier
- Broken phase bubbles nucleate, expand, merge

Assume a 1st Order EW Phase Transition...

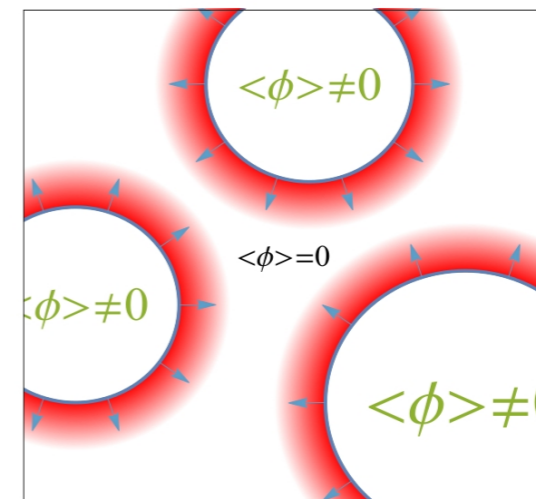
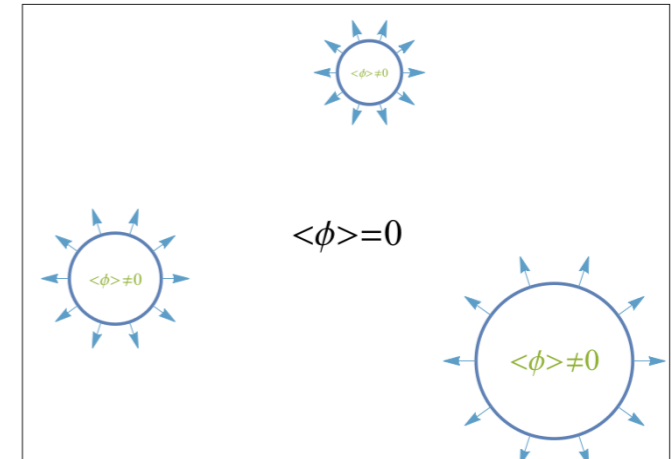
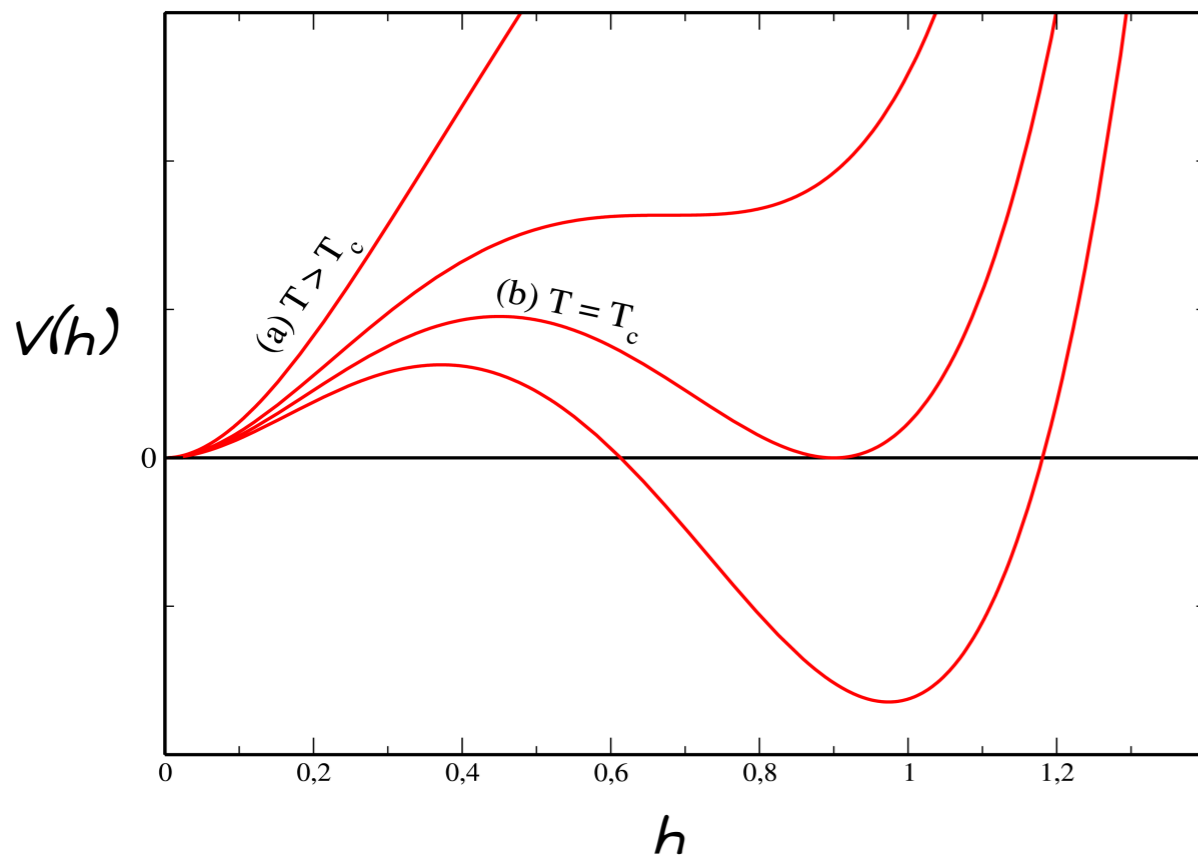
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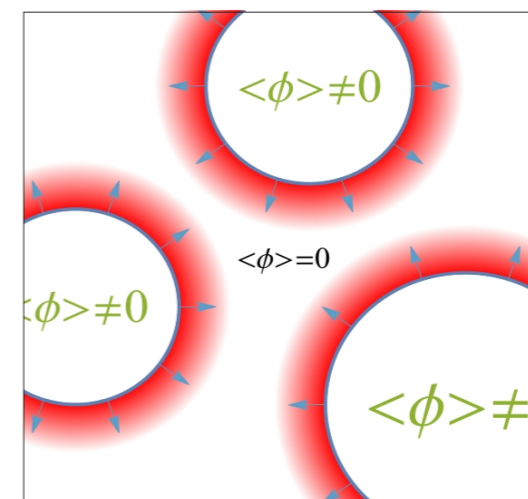
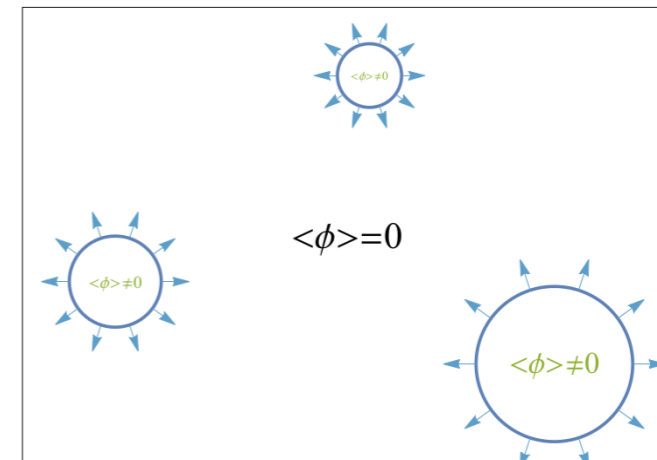
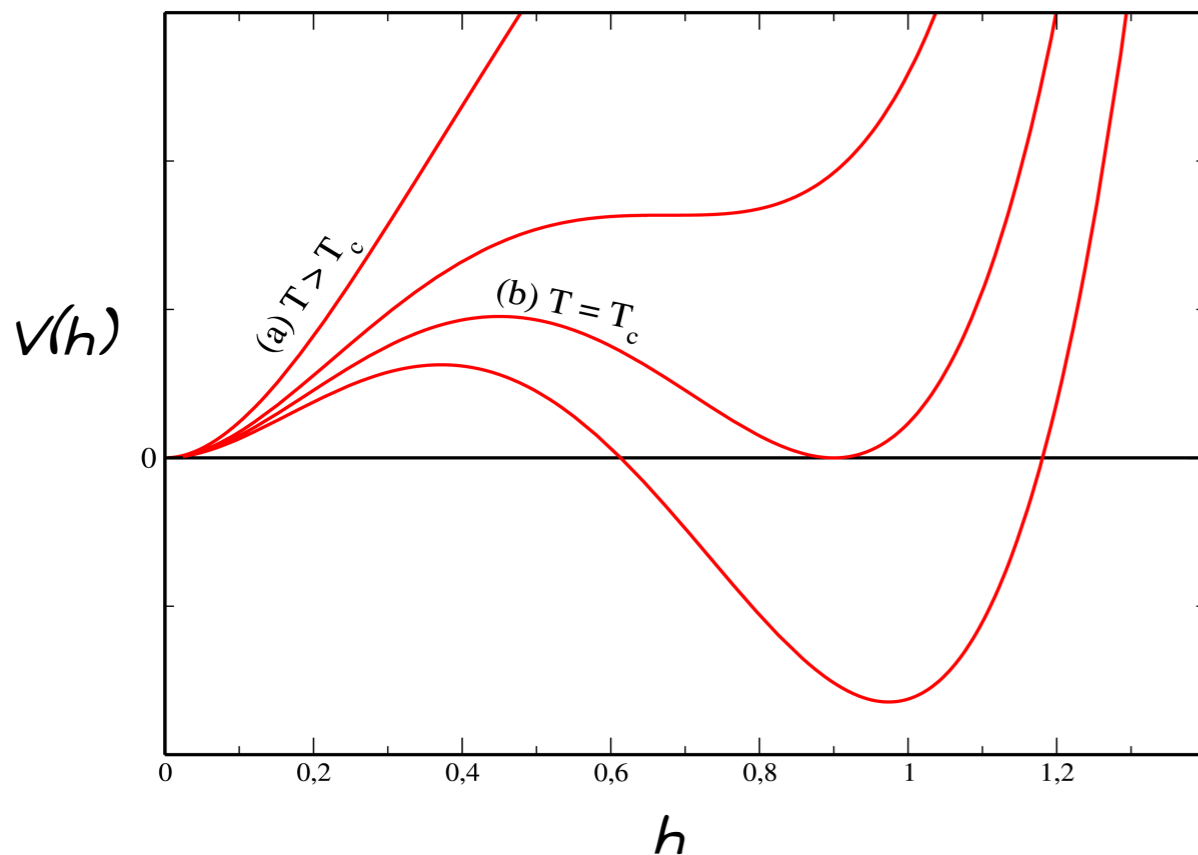


(if in plasma \rightarrow create fluid waves)

- Phases separated by potential barrier
- Broken phase **bubbles** nucleate, **expand**, merge

Assume a 1st Order EW Phase Transition...

Higgs Effective Potential (finite T)



○ Phases separated by potential barrier

○ Broken phase **bubbles** nucleate, expand, **merge/collide**

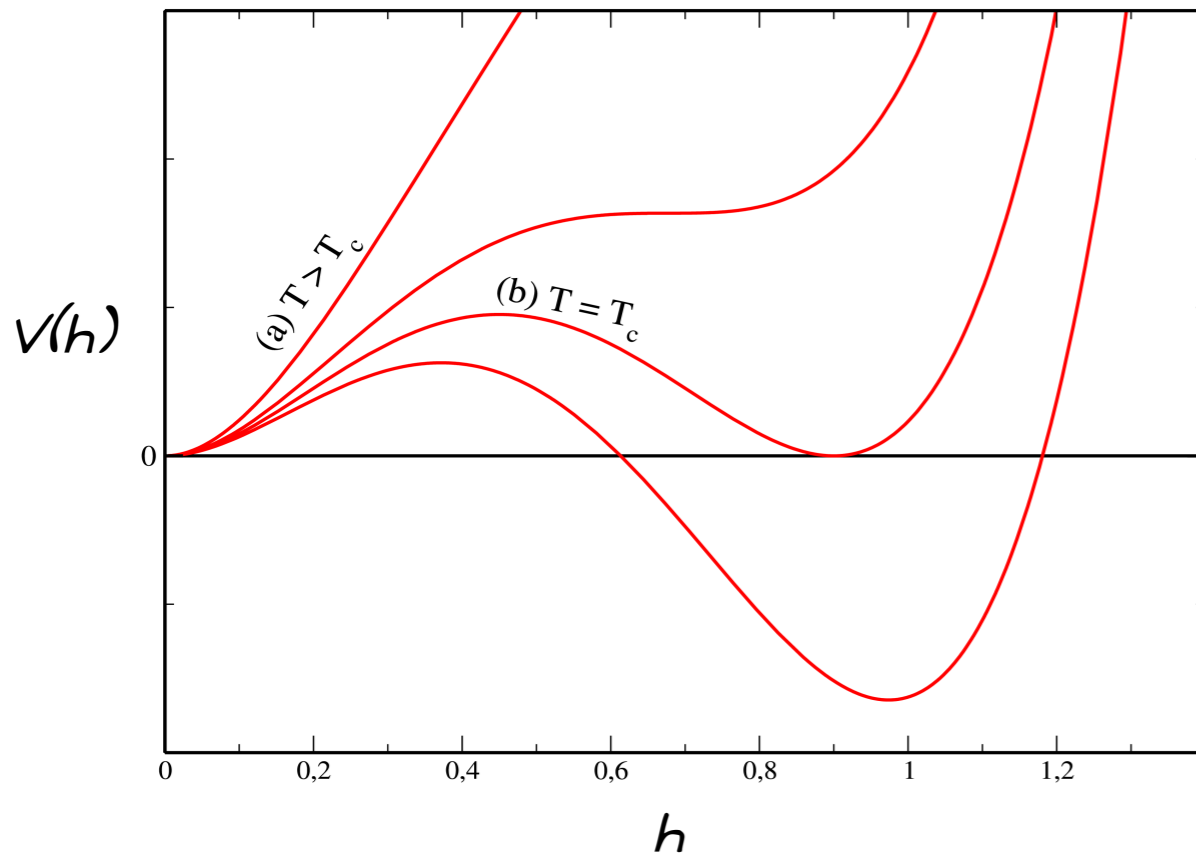
→ Anisotropic Stress



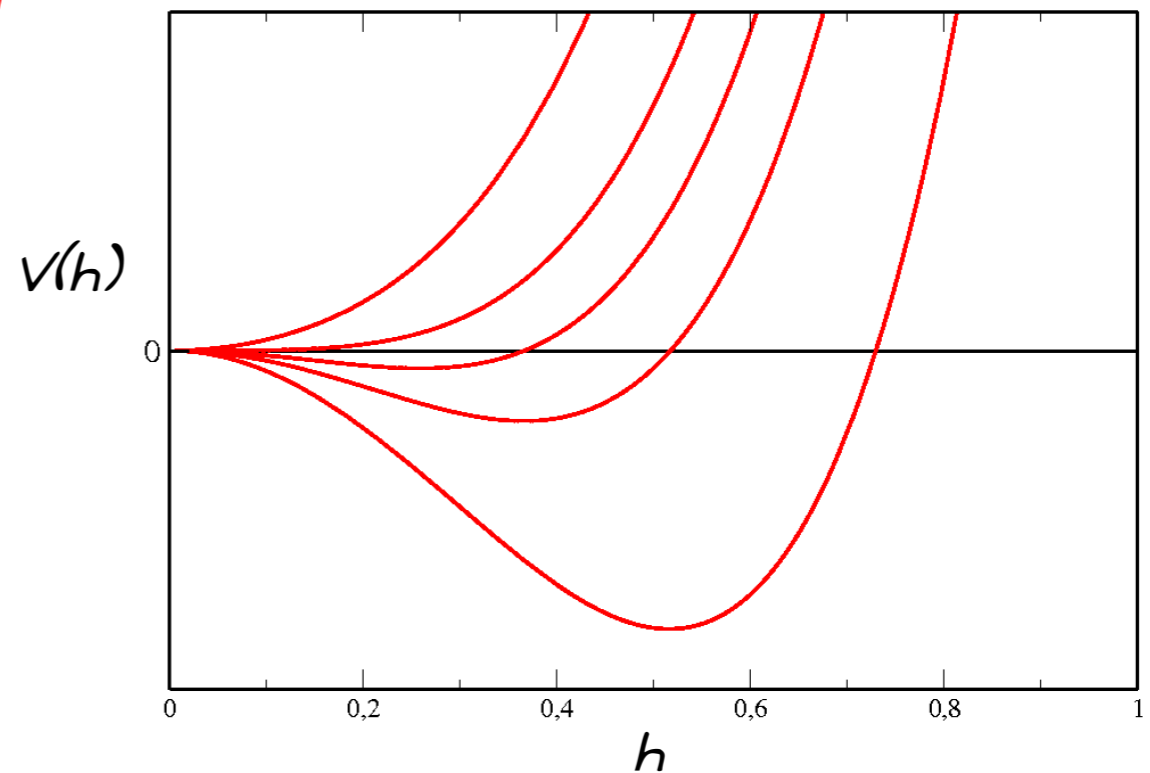
Sources Gravitational Wave Production

Assume a 1st Order EW Phase Transition...

Higgs Effective Potential (finite T)



In the Standard Model...



NO 1st Order EW Phase Transition

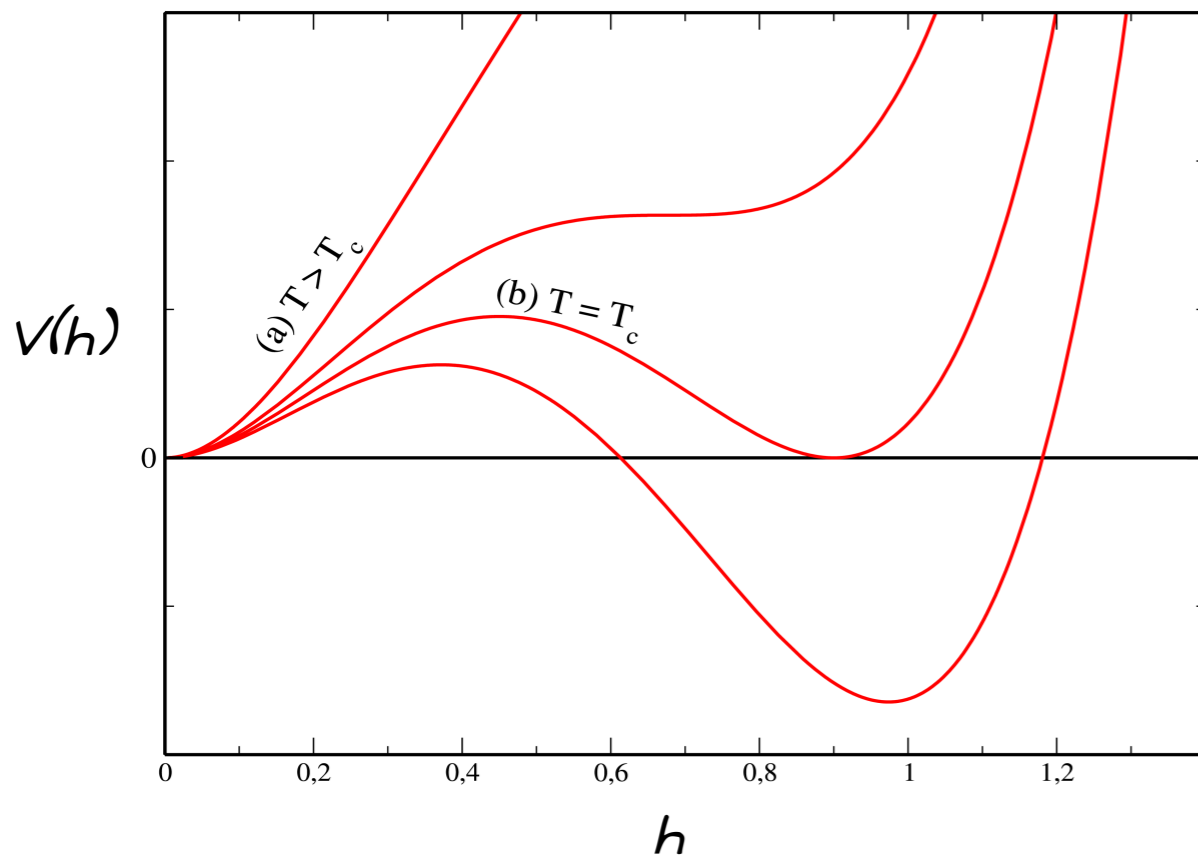
EWPT (non-perturbatively) is smooth cross-over

Kajantie, Laine, Rummukainen, Shaposhnikov, Phys. Rev. Lett. 77 (1996) 2887

Physics Beyond the SM can induce a 1st Order EW Phase Transition

Assume a 1st Order Phase Transition...

Effective Potential (finite T)



Two “Types” of Cosmological 1st Order PTs

○ “Vacuum” Transitions

Fluid/plasma effects negligible

(either plasma is very diluted or coupling between transition field and plasma small/non-existent)

Bubble walls accelerate until collision

Energy of PT stored in bubble walls

○ Thermal Transitions

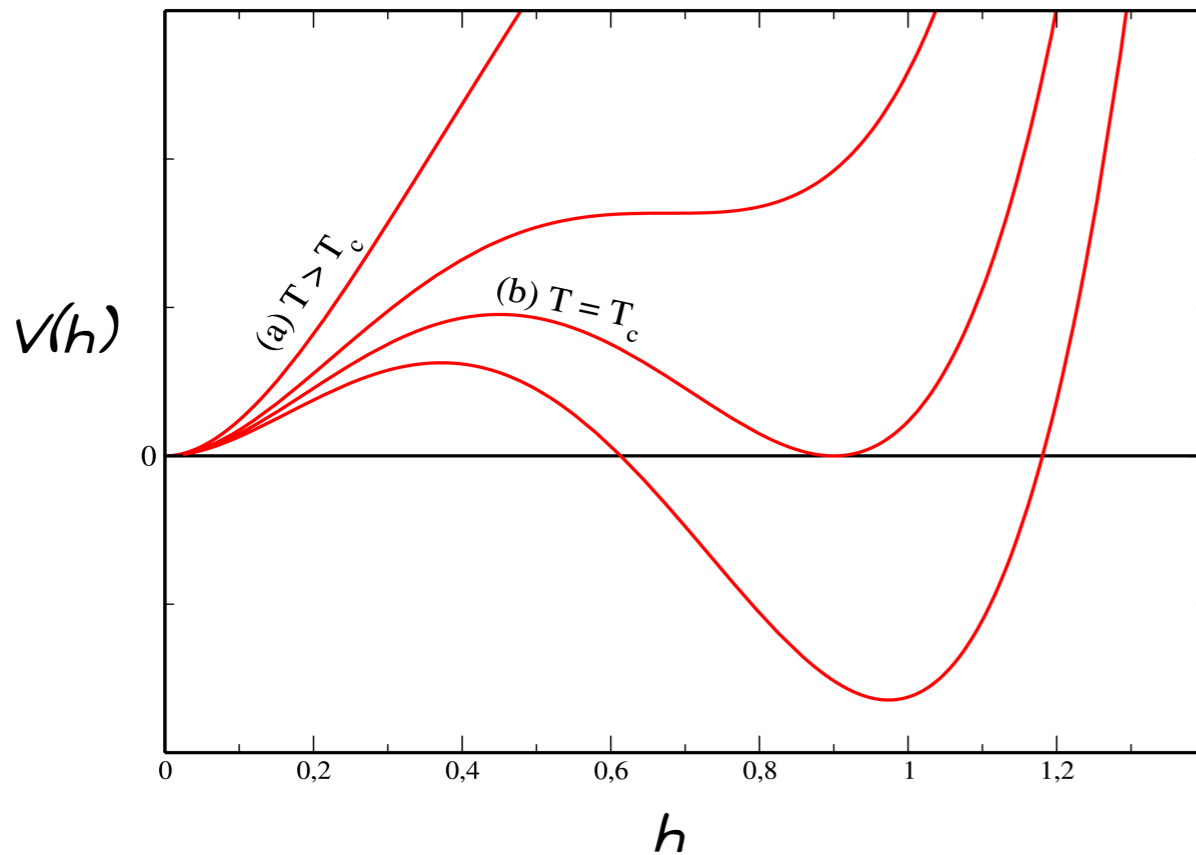
Energy of PT transferred to plasma

Plasma exerts friction on bubble wall

Terminal bubble wall velocity
(steady state)

Assume a 1st Order Phase Transition...

Effective Potential (finite T)



- Decay rate $\Gamma(T) \approx T^4 \exp\left(-\frac{S_3(T)}{T}\right)$

- O(3) symmetric action

$$S_3(T) = 4\pi \int dr r^2 \left[\frac{1}{2} \left(\frac{d\phi}{dr} \right)^2 + V(\phi, T) \right]$$

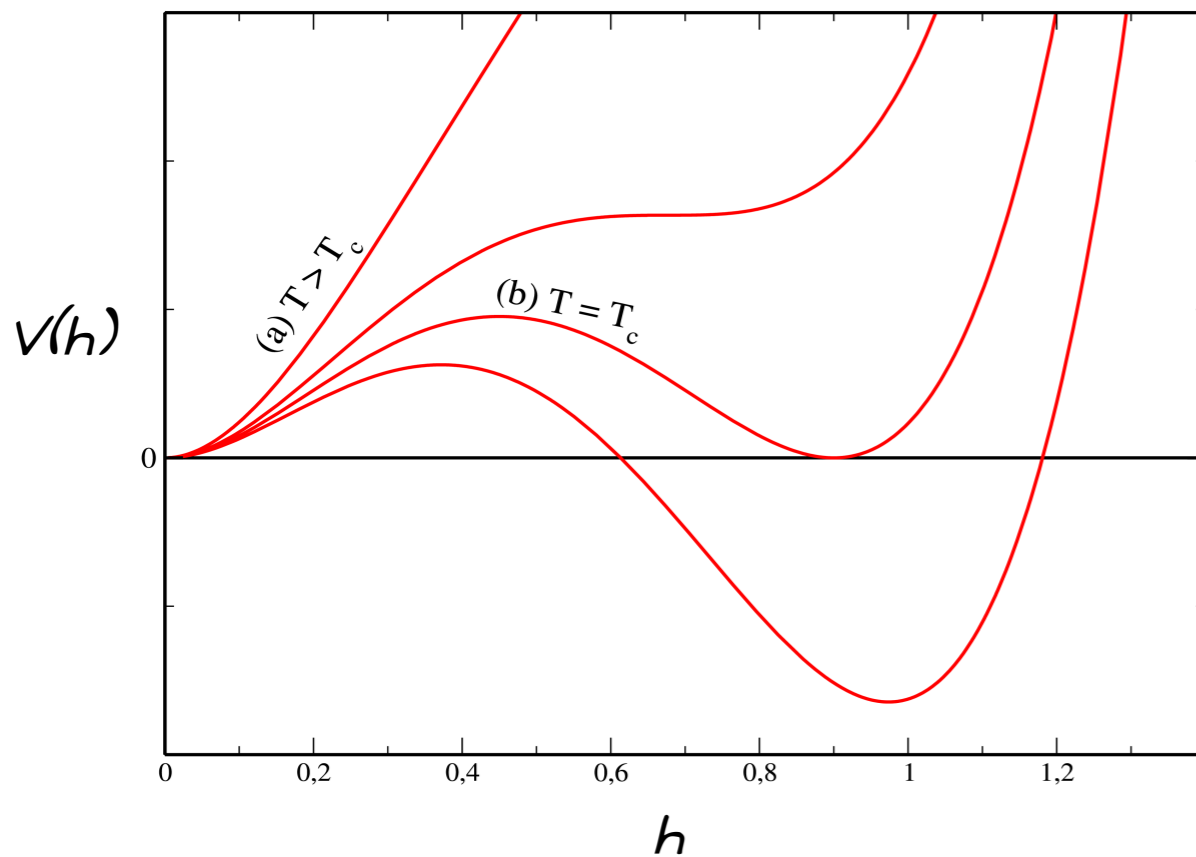
- Bubble profile (bounce)

$$\frac{d^2 \phi}{dr^2} + \frac{2}{r} \frac{d\phi}{dr} - \frac{\partial V(\phi, T)}{\partial \phi} = 0$$

$$\phi(r \rightarrow \infty) = 0 \quad \text{and} \quad \dot{\phi}(r = 0) = 0$$

Assume a 1st Order Phase Transition...

Effective Potential (finite T)



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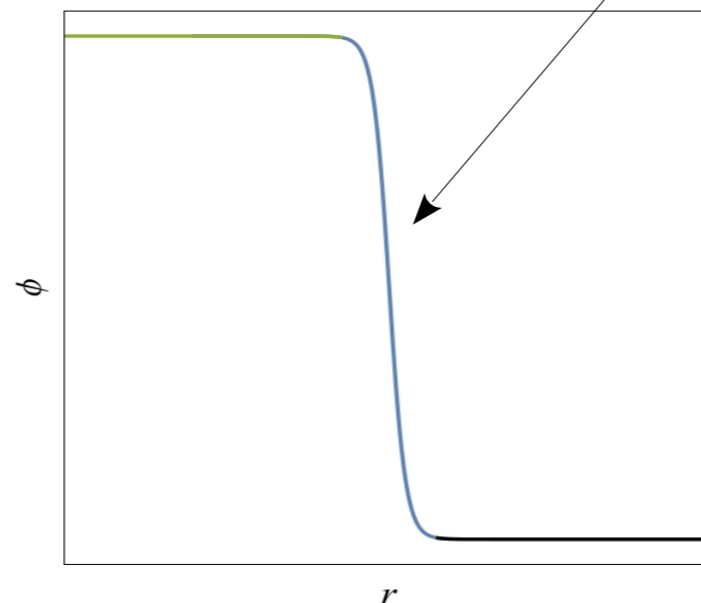
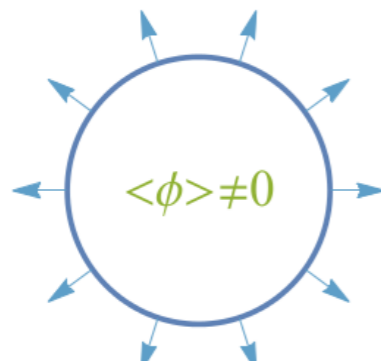
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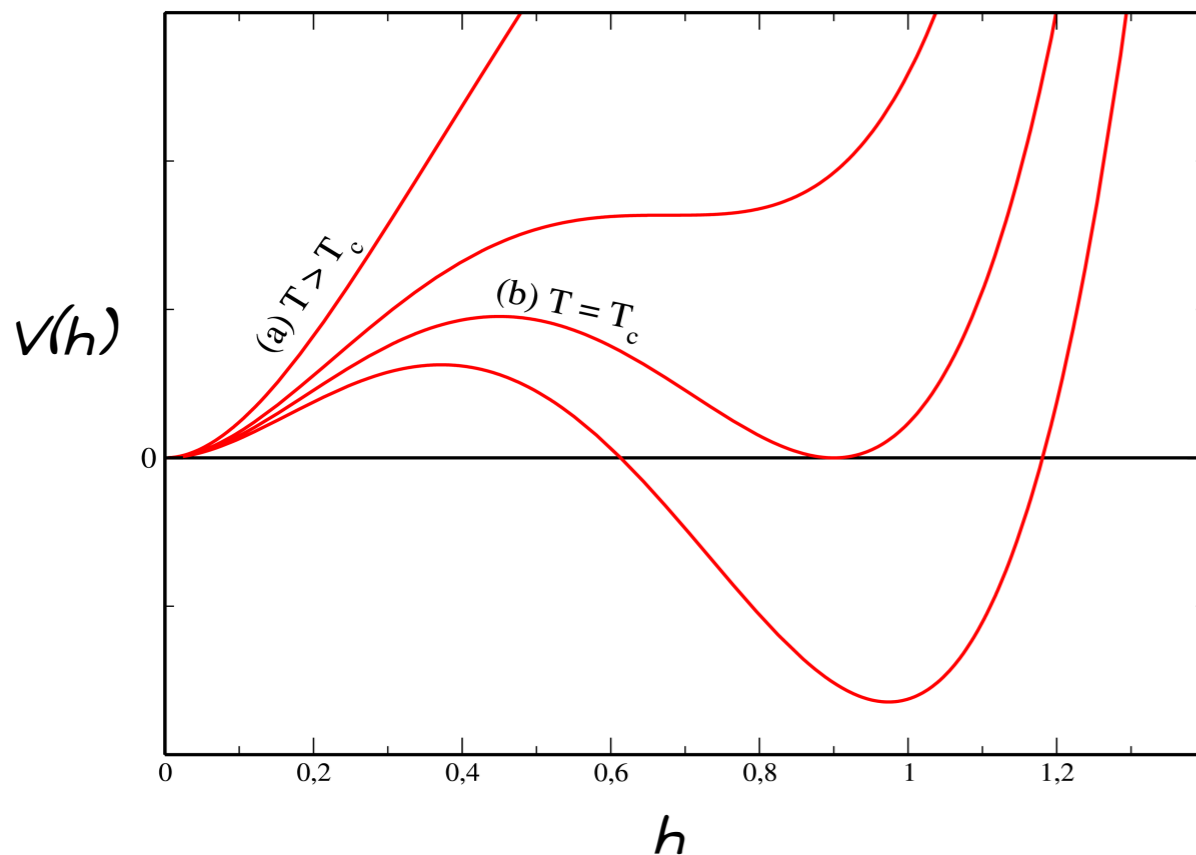
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$$\langle \phi \rangle = 0$$



Assume a 1st Order Phase Transition...

Effective Potential (finite T)



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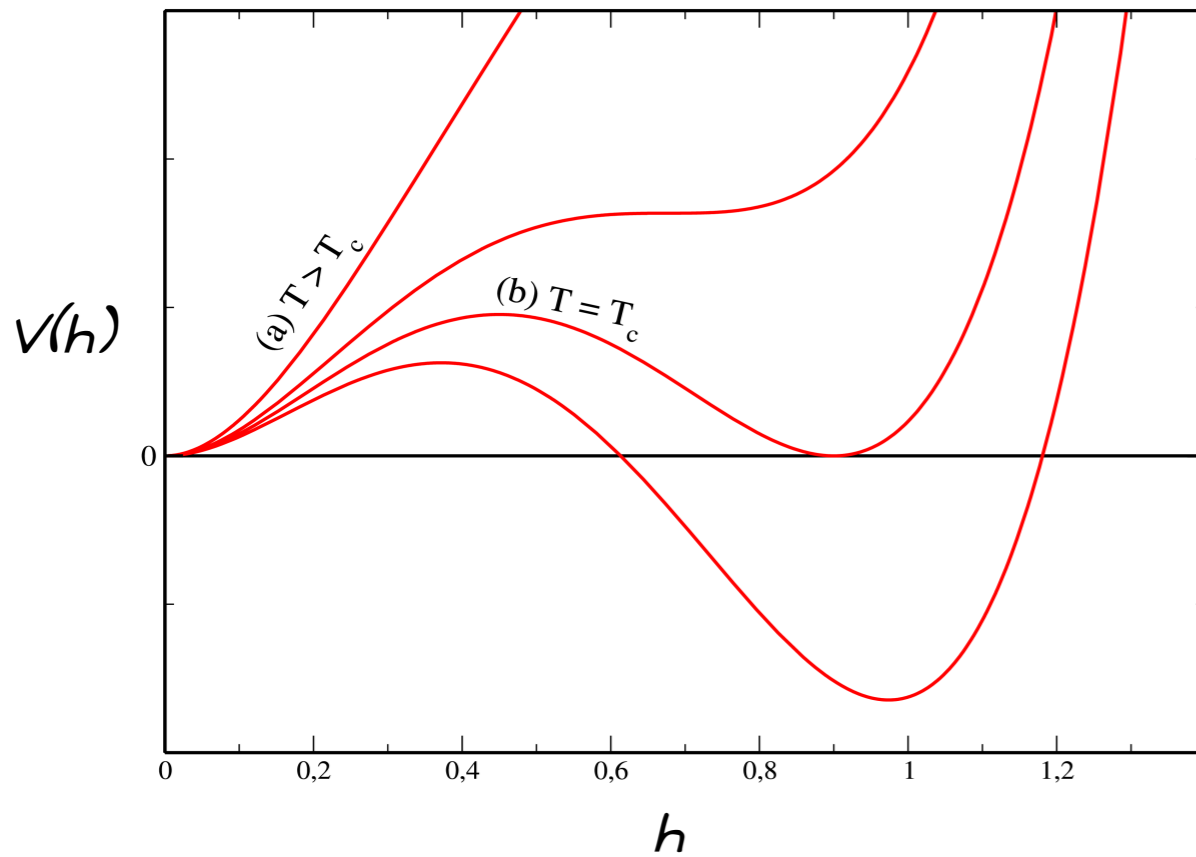
Nucleation temperature:

One Higgs bubble per Horizon volume (on average)

$$N(T_n) = \int_{t_c}^{t_n} dt \frac{\Gamma(t)}{H(t)^3} = \int_{T_n}^{T_c} \frac{dT}{T} \frac{\Gamma(T)}{H(T)^4} = 1$$

Assume a 1st Order Phase Transition...

Effective Potential (finite T)



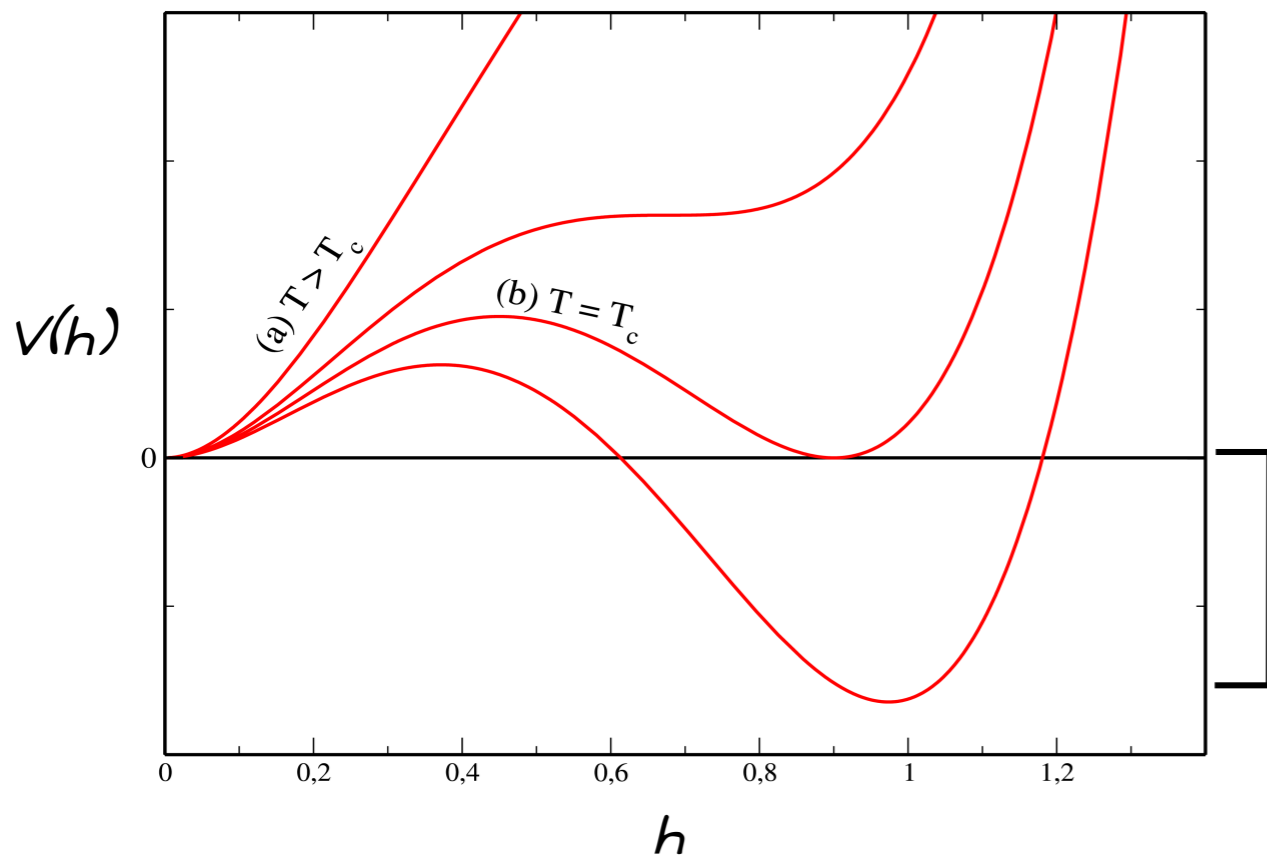
Two KEY Phase Transition Quantities:

○ (Available) Transition Energy (normalized)

$$\alpha = \frac{\epsilon}{a_+ T^4}$$

Assume a 1st Order Phase Transition...

Effective Potential (finite T)



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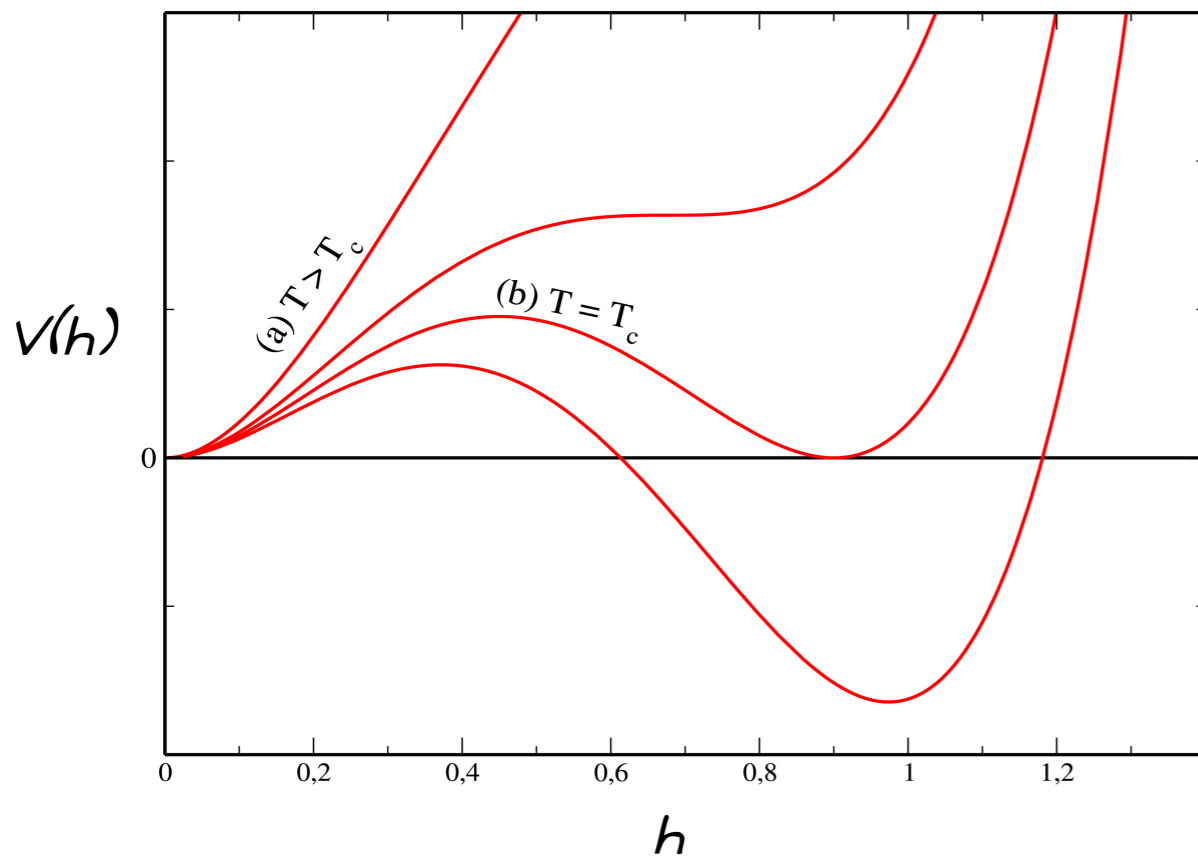
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Radiation Energy Density

Assume a 1st Order Phase Transition...

Effective Potential (finite T)



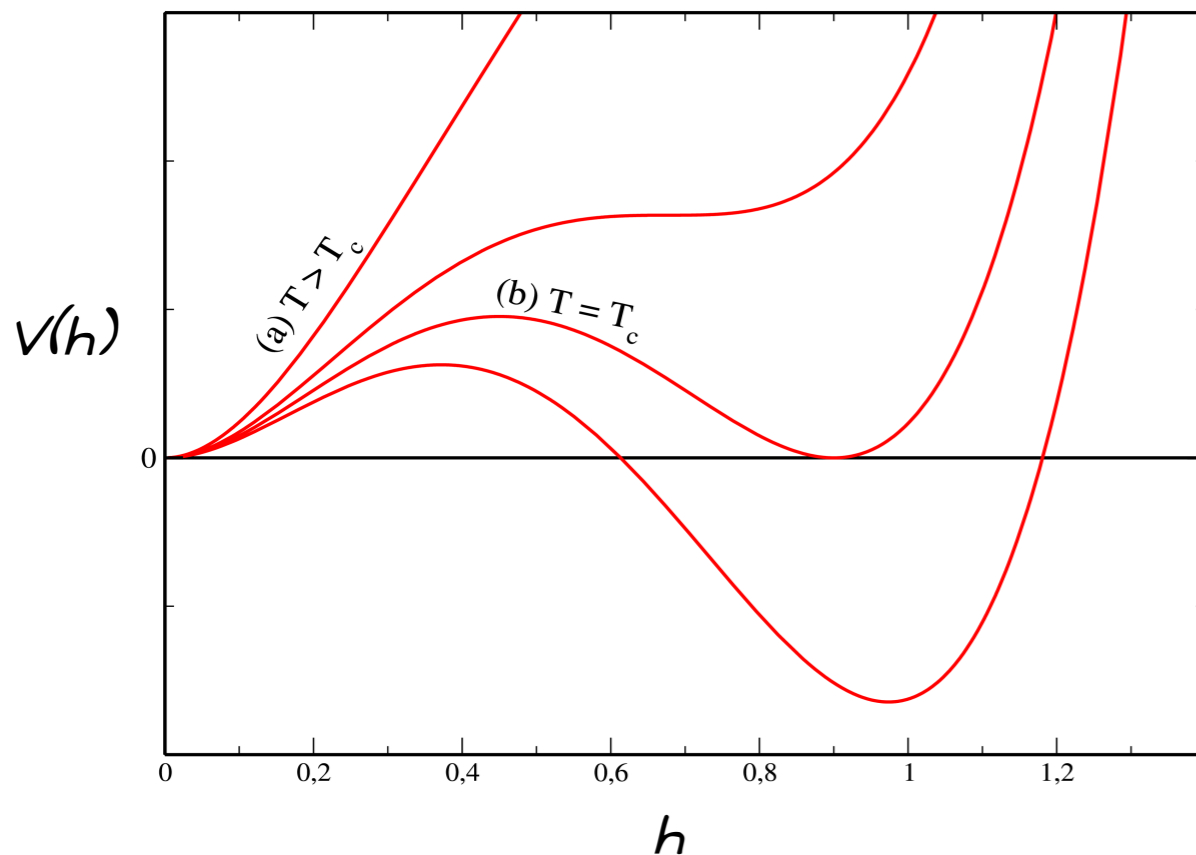
Two KEY Phase Transition Quantities:

○ (Available) Transition Energy (normalized)

$$\alpha_e \equiv \frac{4}{3} \frac{\Delta e(T_n)}{w_+(T_n)}$$

Assume a 1st Order Phase Transition...

Effective Potential (finite T)



Two KEY Phase Transition Quantities:

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$$\alpha_e \equiv \frac{4 \Delta e(T_n)}{3 w_+(T_n)}$$

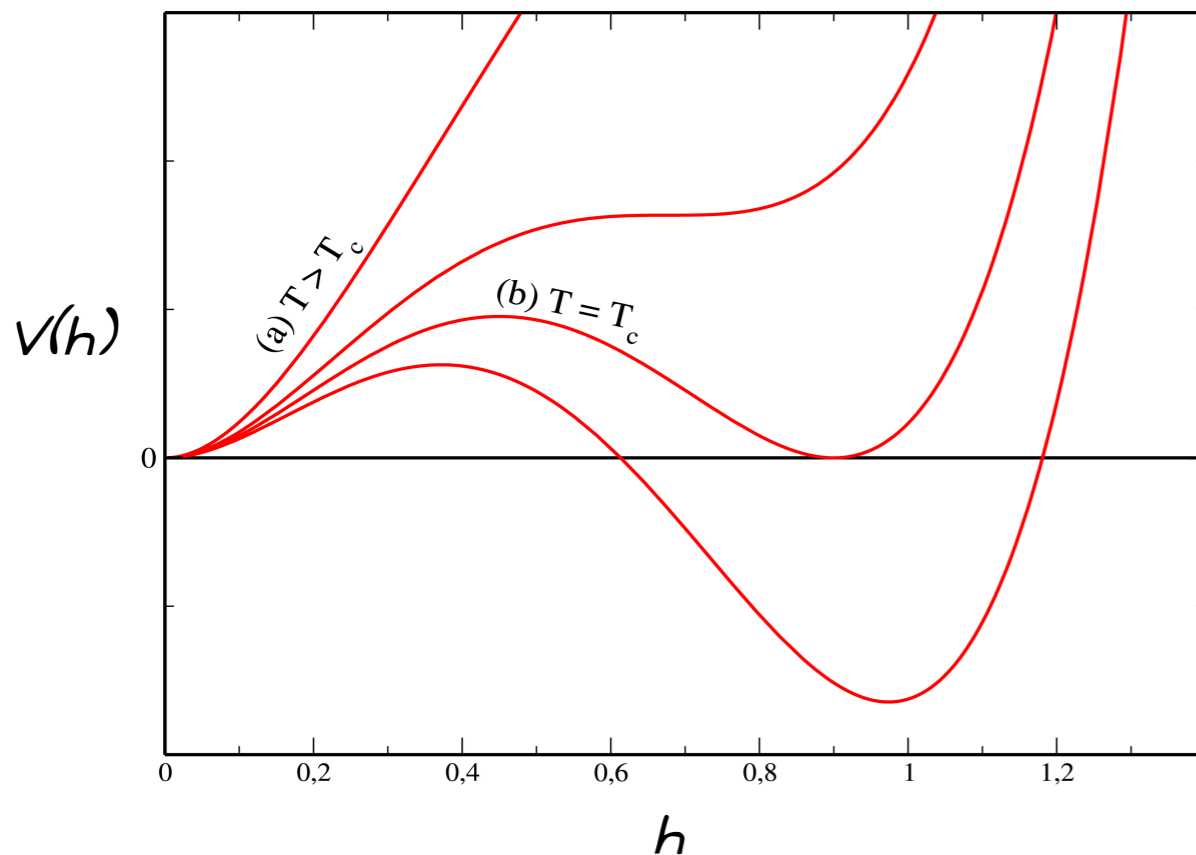
- Duration of the Transition (-1)

$$\frac{\beta}{H} \equiv - \left. \frac{dS_3}{dt} \right|_{t=t_n} \approx T \left. \frac{d(S_3/T)}{dT} \right|_{T=T_n}$$

(Related to the change of the Decay Rate)

Assume a 1st Order Phase Transition...

Effective Potential (finite T)



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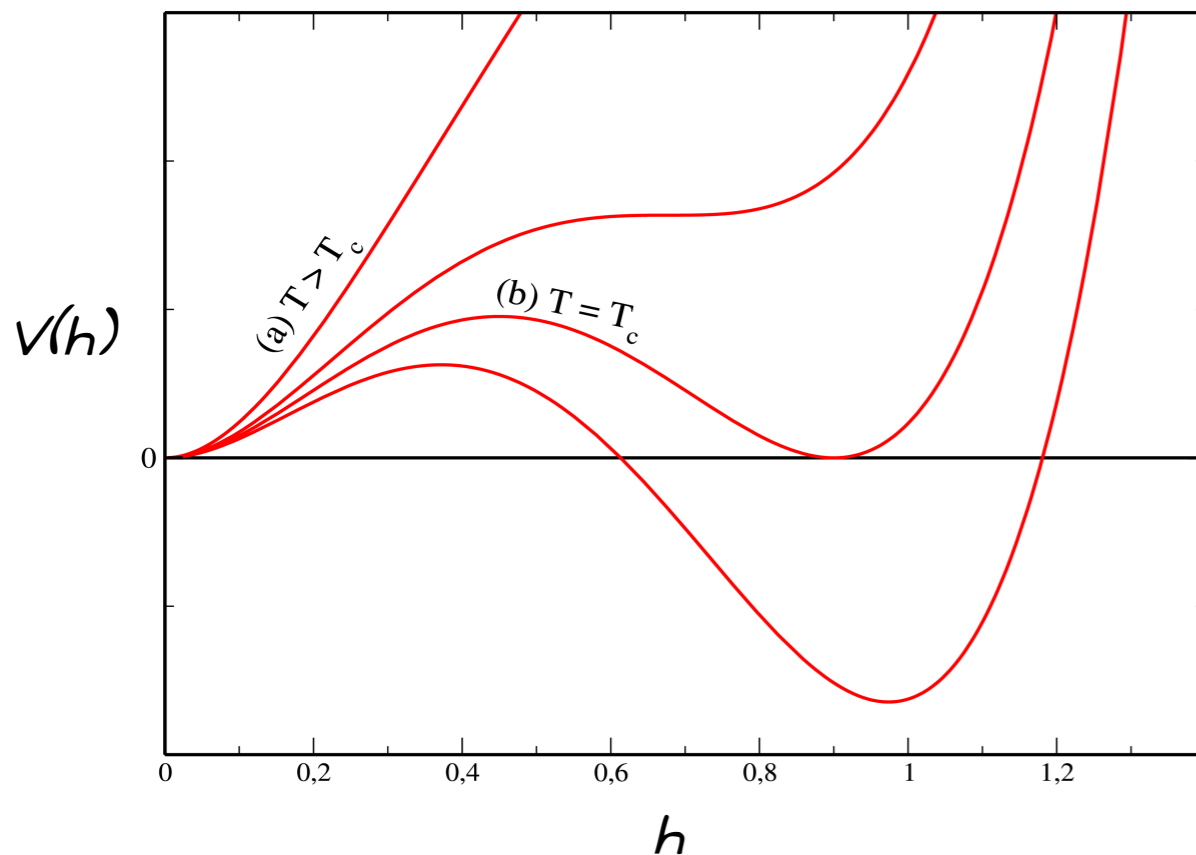
Average number of bubbles per horizon at the time of bubble coalescence/percolation

(Transition Completes, T_*)

H_*

Assume a 1st Order Phase Transition...

Effective Potential (finite T)



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(Related to the change of the Decay Rate)

▶ GW frequency ~ size of bubbles @ collision

▶ For $T_* \sim 100$ GeV and $\frac{\beta}{H_*} \sim 100$, GW frequency (redshifted to today!) ~ **mHz**

Average number of bubbles per horizon at the time of bubble coalescence/percolation

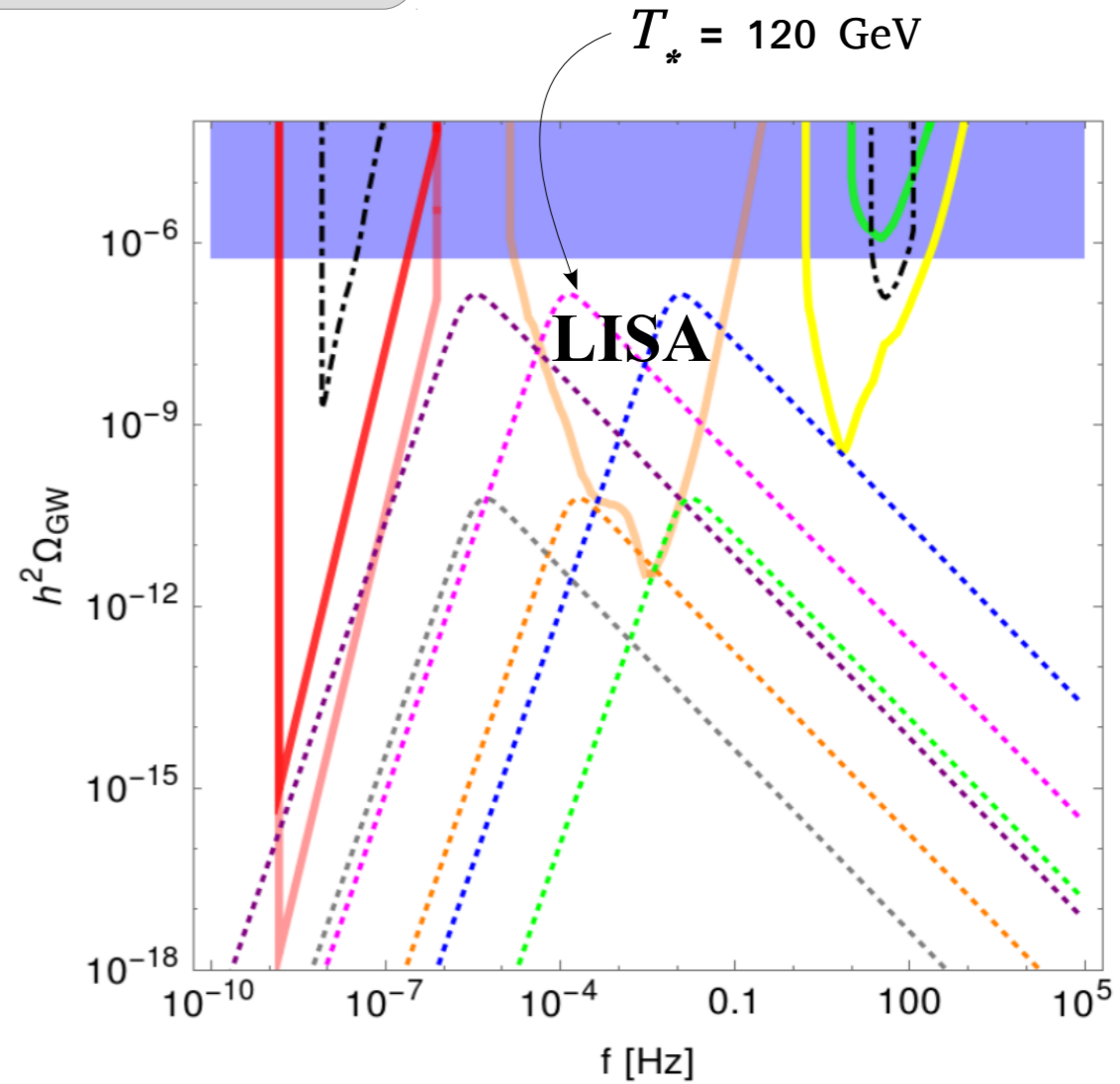
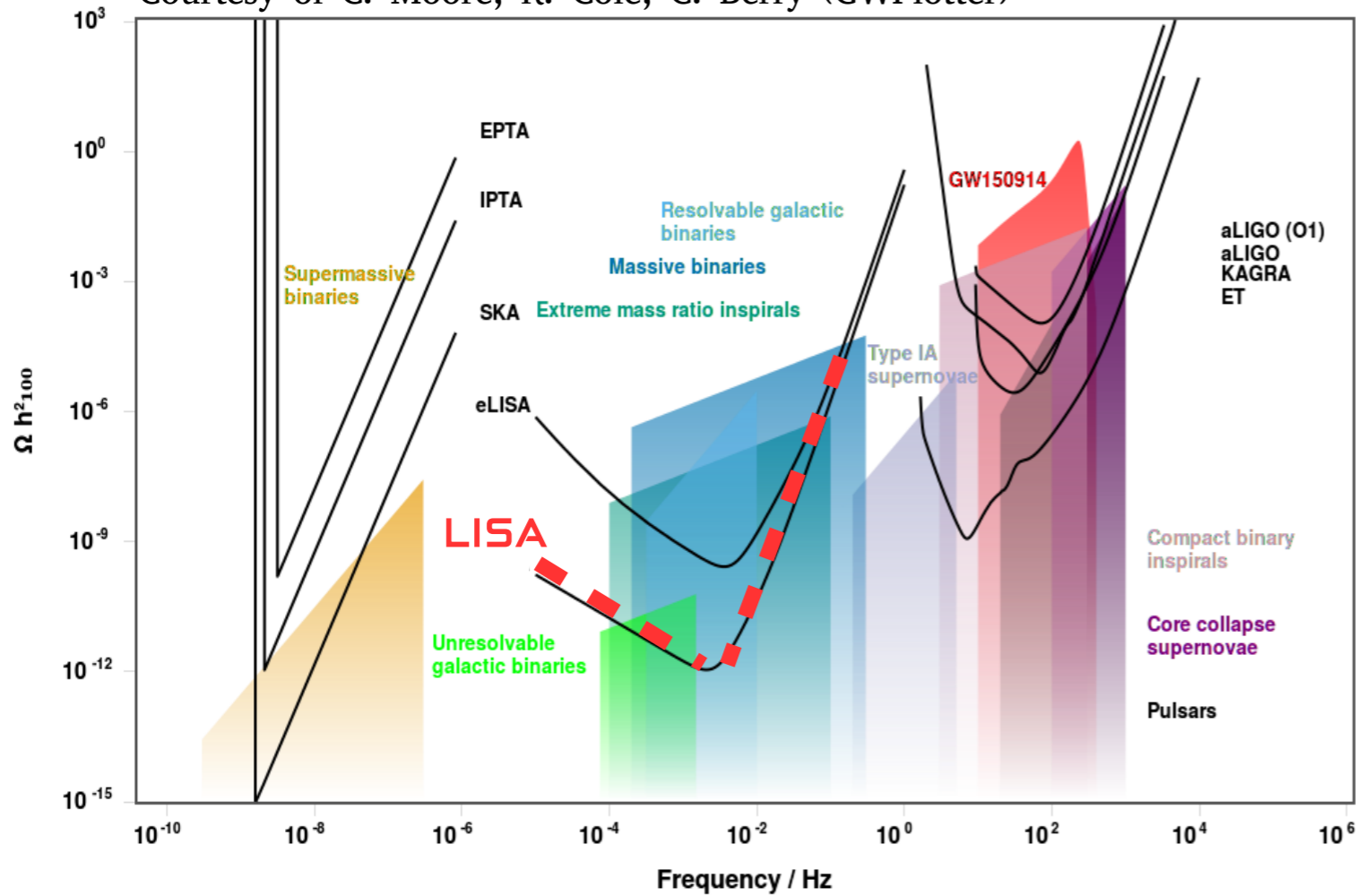
(Transition Completes, T_*)

H_*

1st Order (EW) Phase Transition

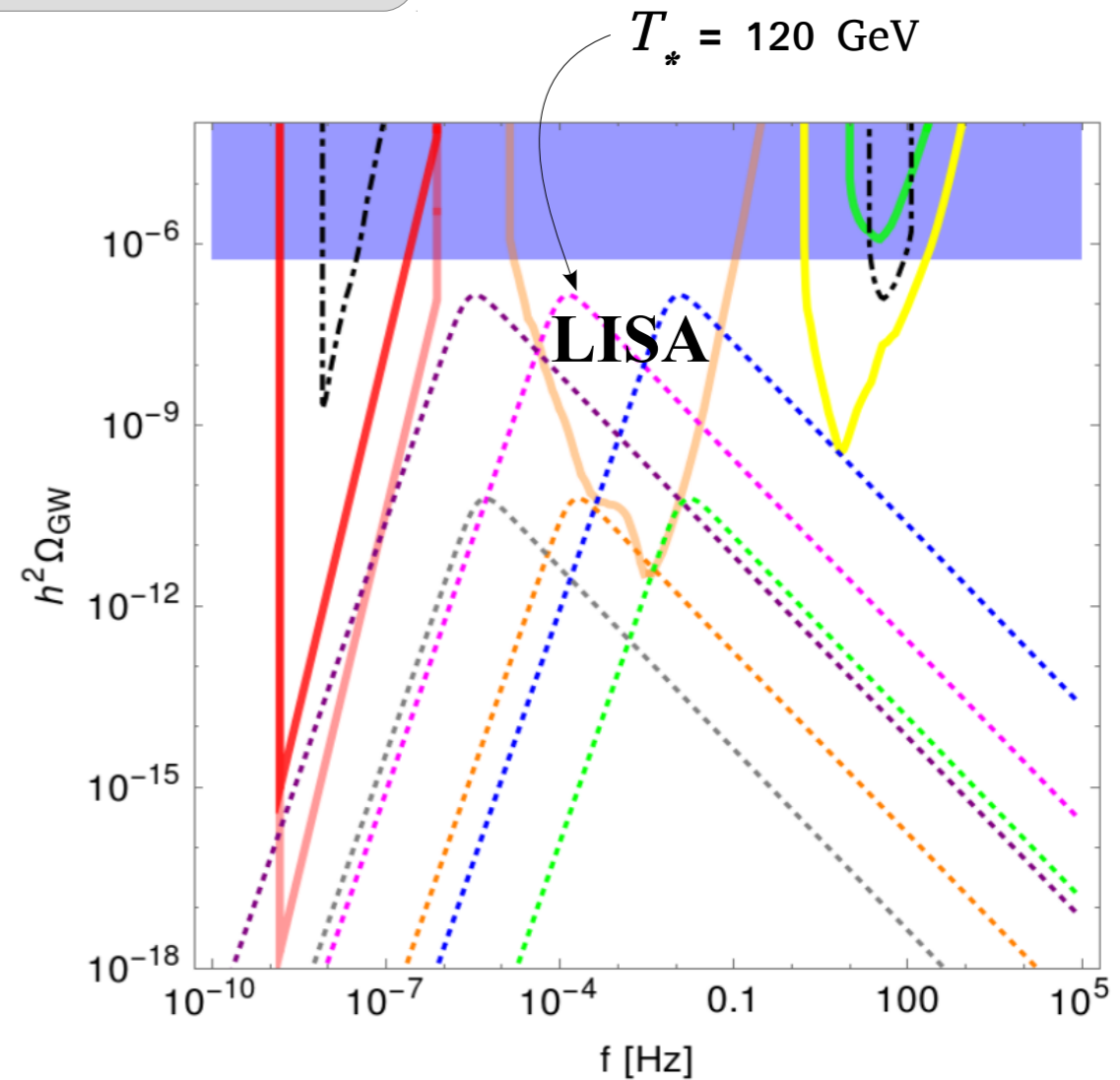
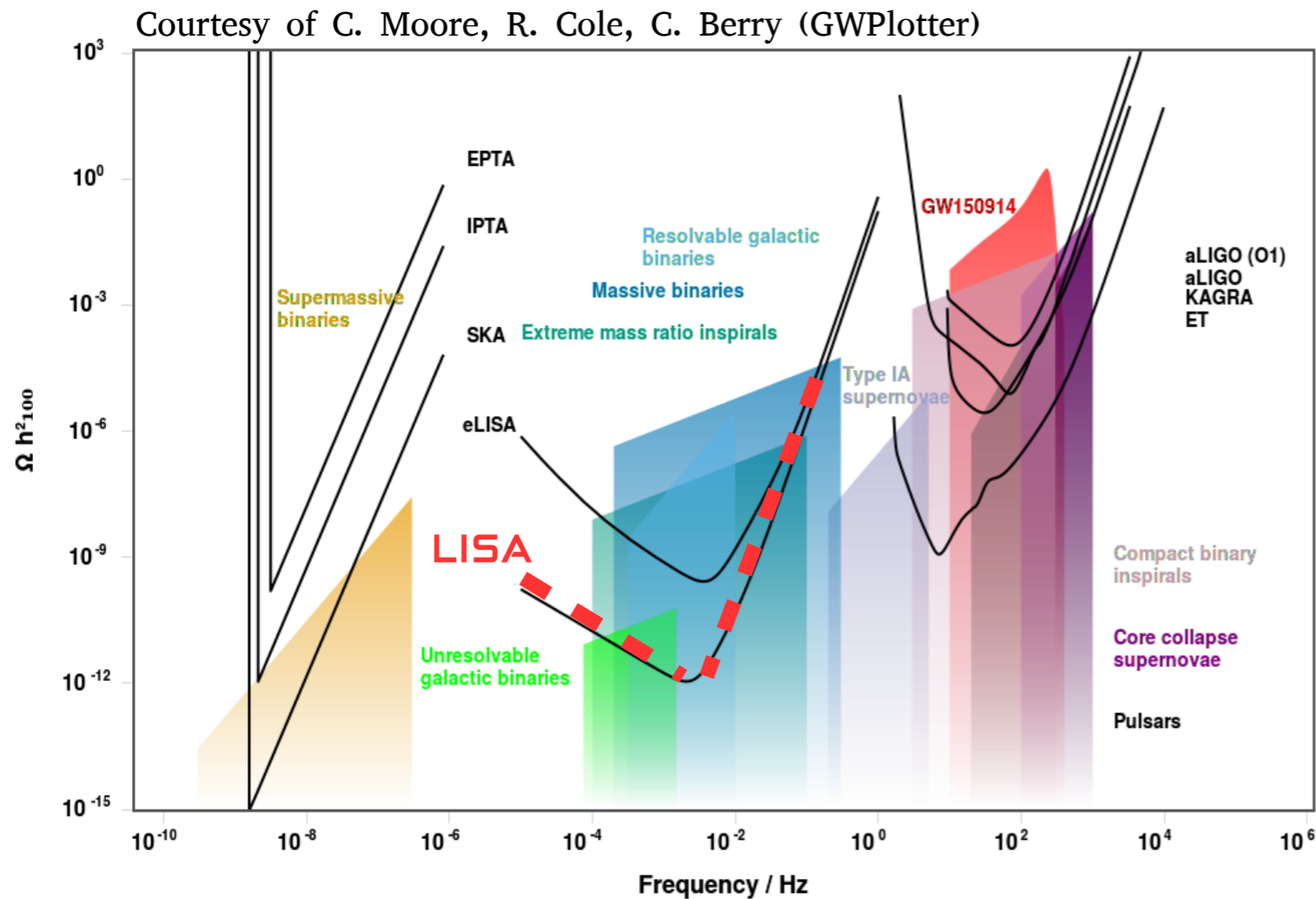
mHz GW Signal in the sensitivity band of future space-based GW detector **LISA**

Courtesy of C. Moore, R. Cole, C. Berry (GWPlotter)



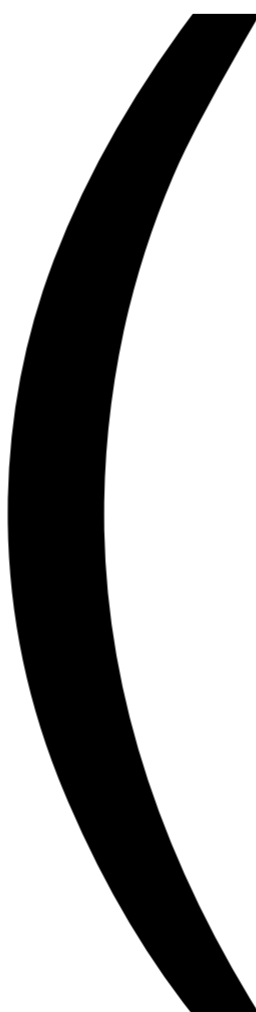
1st Order (EW) Phase Transition

mHz GW Signal in the sensitivity band of future space-based GW detector **LISA**



Figuroa et al., PoS GRASS2018 (2018) 036

LISA can probe the **EW** epoch of the early Universe



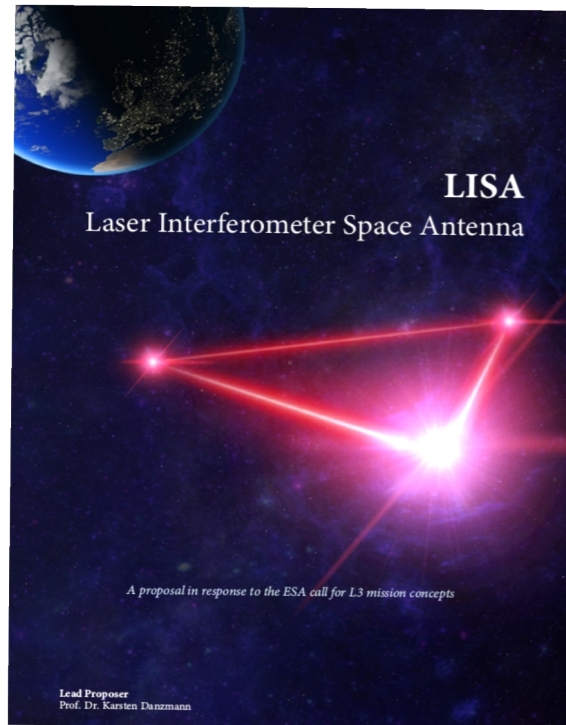
The LISA Mission

(Laser Interferometer Space Antenna)

A brief status report

Thanks to G. Nardini

2017: LISA proposal to ESA

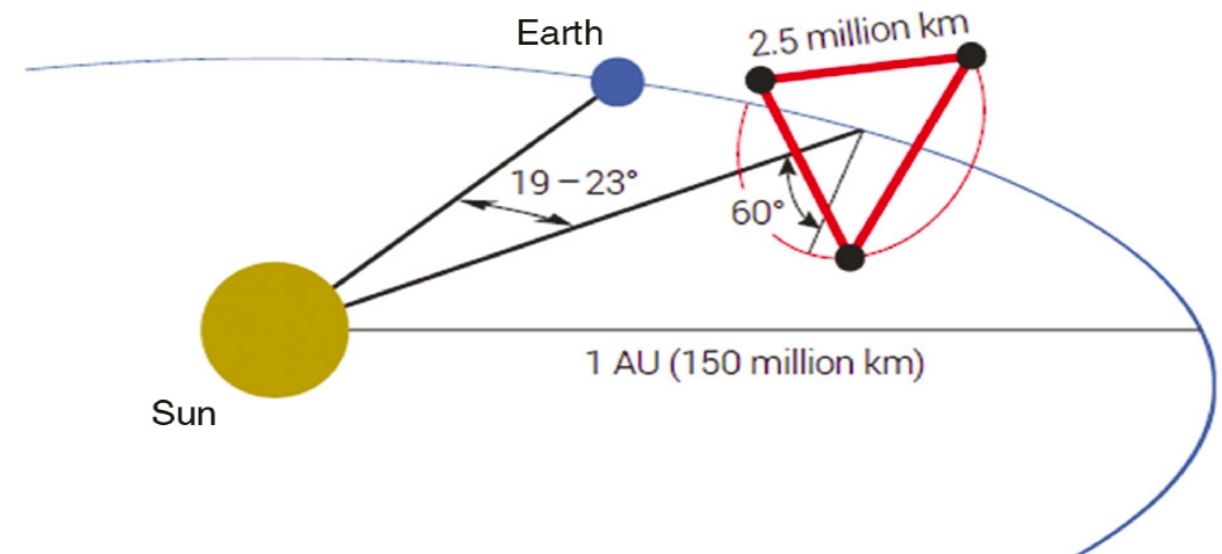


LISA Collaboration, 1702.00786

Launch date 2030-2034

LISA Mission selected by ESA (Summer 2017)

+ (On Jan 22 2018, LISA passed ESA's Mission Definition Review)



4 years of lifetime (w. consumables up to 10 years)

2.5 MKm (arm length)

From the proposal:

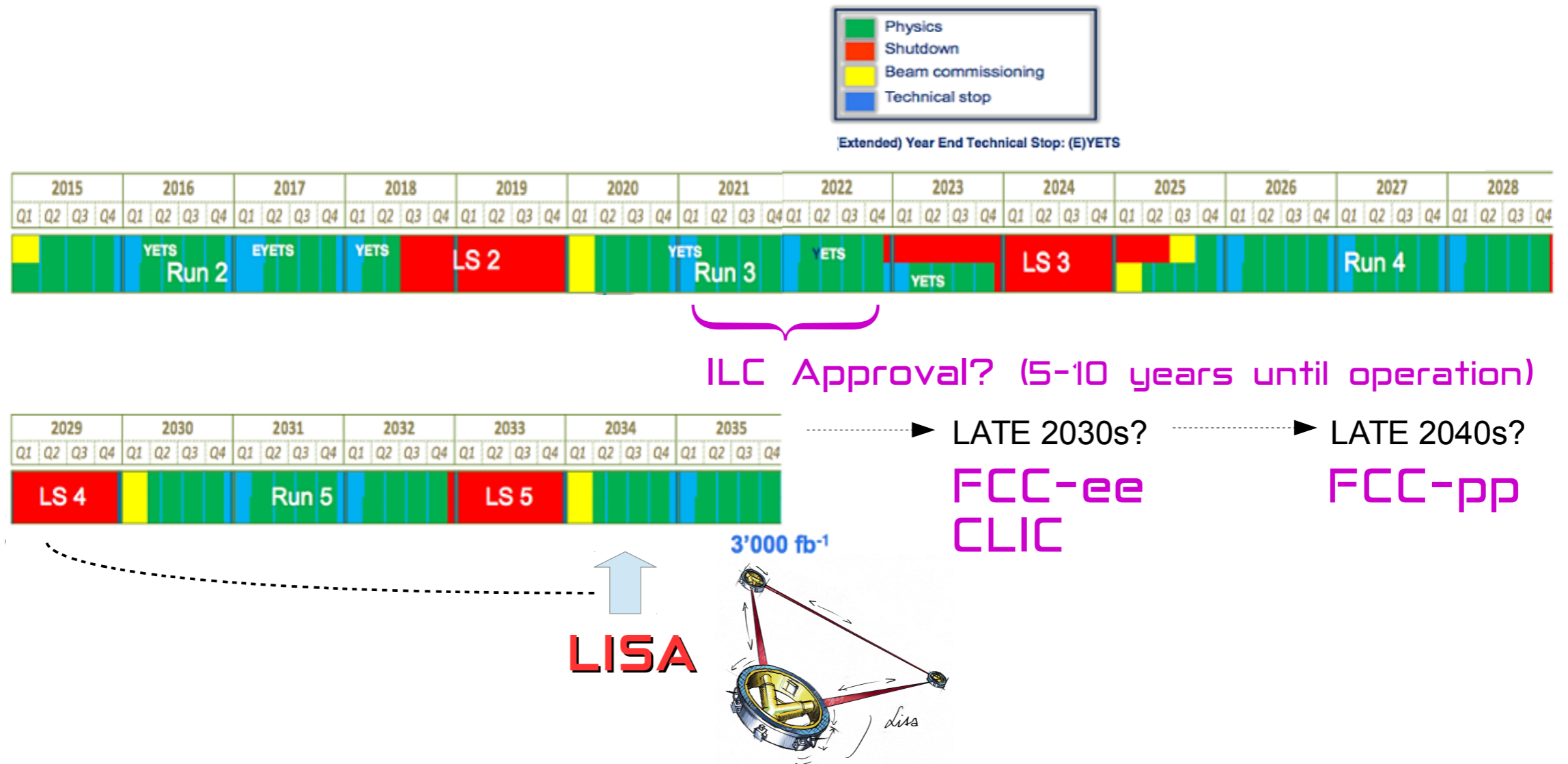
Audley et al, arXiv:1702.00786

SI7.2 : Measure, or set upper limits on, the spectral shape of the cosmological stochastic GW background

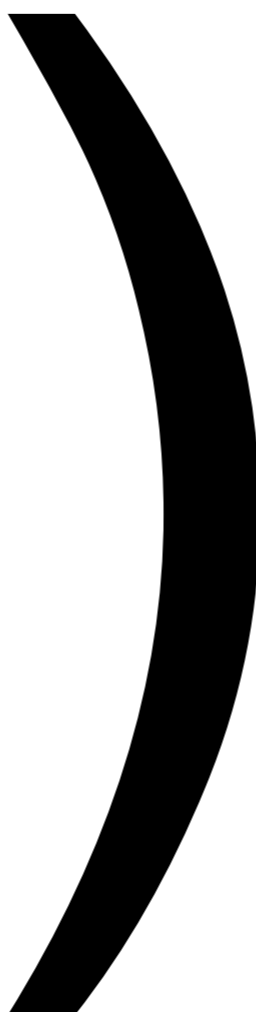
OR7.2: Probe a broken power-law stochastic background from the early Universe as predicted, for example, by first order phase transitions [21] (other spectral shapes are expected, for example, for cosmic strings [22] and inflation [23]). Therefore, we need the ability to measure $\Omega = 1.3 \times 10^{-11} (f/10^{-4} \text{ Hz})^{-1}$ in the frequency ranges $0.1 \text{ mHz} < f < 2 \text{ mHz}$ and $2 \text{ mHz} < f < 20 \text{ mHz}$, and $\Omega = 4.5 \times 10^{-12} (f/10^{-2} \text{ Hz})^3$ in the frequency ranges $2 \text{ mHz} < f < 20 \text{ mHz}$ and $0.02 < f < 0.2 \text{ Hz}$.

GW – Collider complementarity

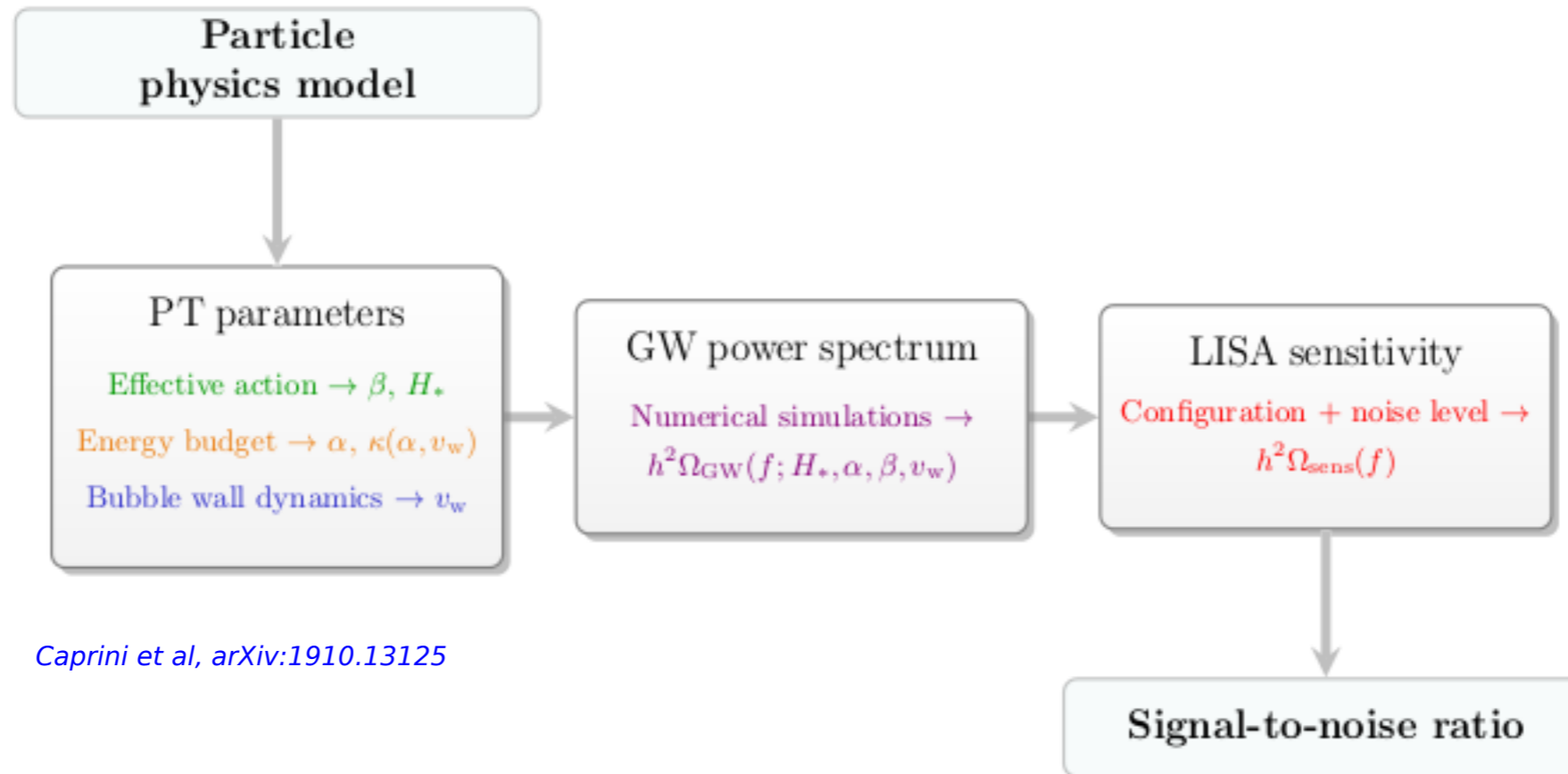
Timeline: LISA GW Observatory in the Context of High-Energy Colliders



After LHC, LISA is next step in exploration of EW scale physics

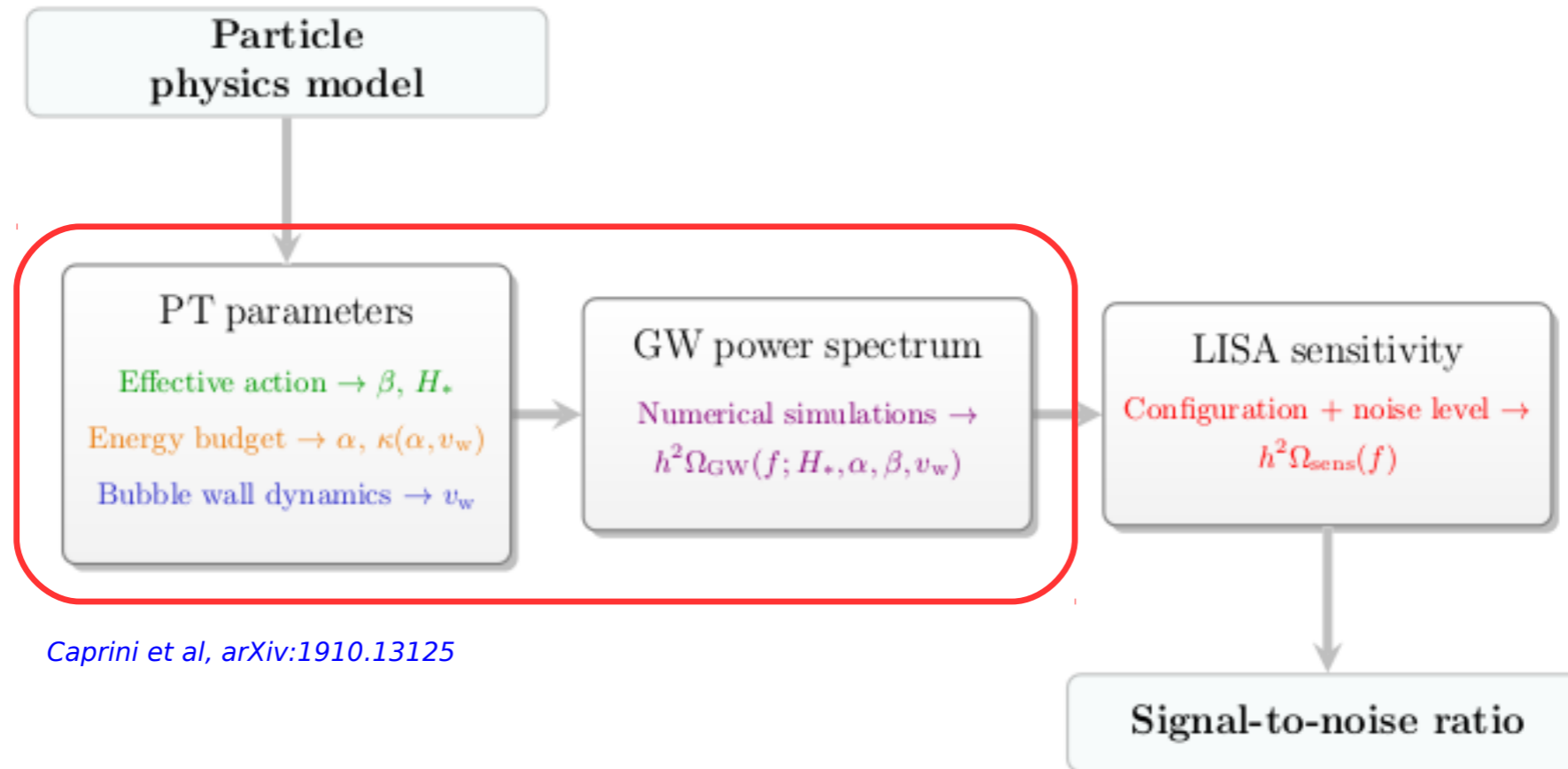


GW from the EW Phase Transition with LISA



Assess the capability of LISA to probe GW signal from EW epoch \Rightarrow BSM physics

GW from the EW Phase Transition with LISA



Assess the capability of LISA to probe GW signal from EW epoch \Rightarrow BSM physics

Need to predict GW signal as robustly as possible

Thermal EW Phase Transition

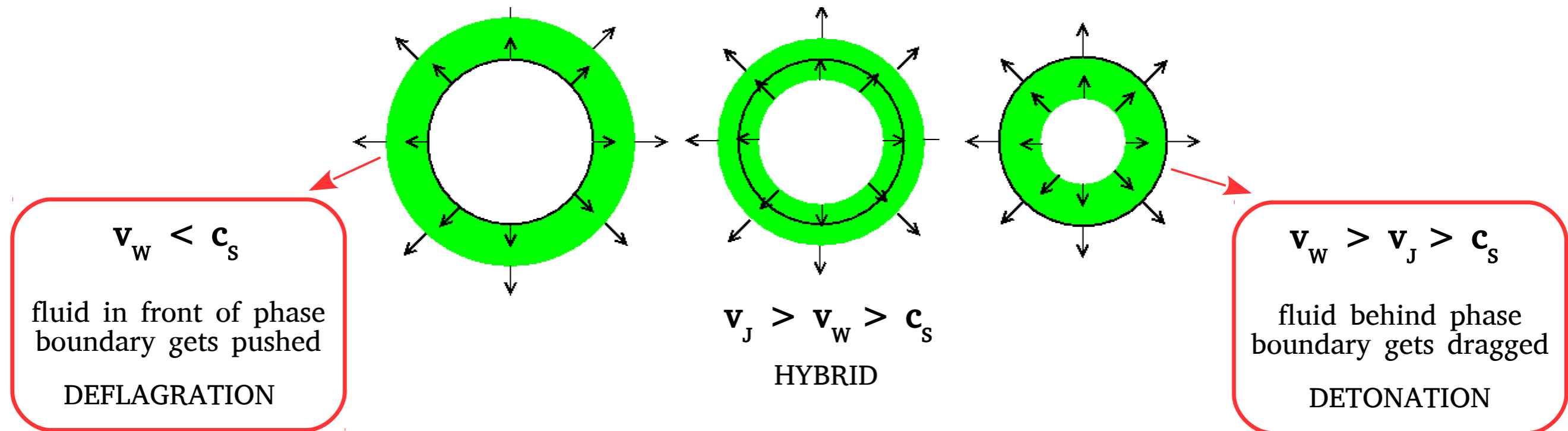
Energy liberated from phase change transferred (mostly) to plasma

- Kinetic energy \Rightarrow Thermal plasma bulk motion
- Thermal energy \Rightarrow Thermal plasma gets heated up

Depending on Higgs bubble wall velocity, energy transfer to plasma creates different types of **expanding fluid shells**

Laine, Phys. Rev. D **49** (1994) 3847

Espinosa, Konstandin, No, Servant, JCAP **1006** (2010) 028



Thermal EW Phase Transition

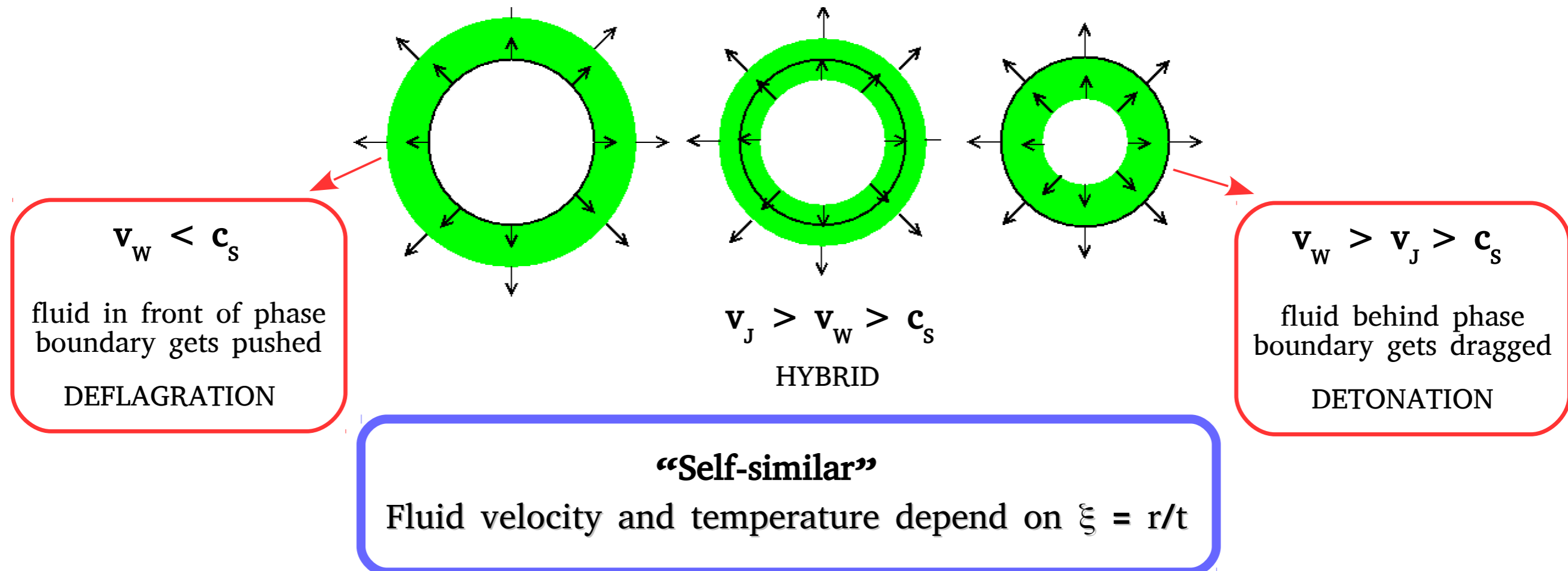
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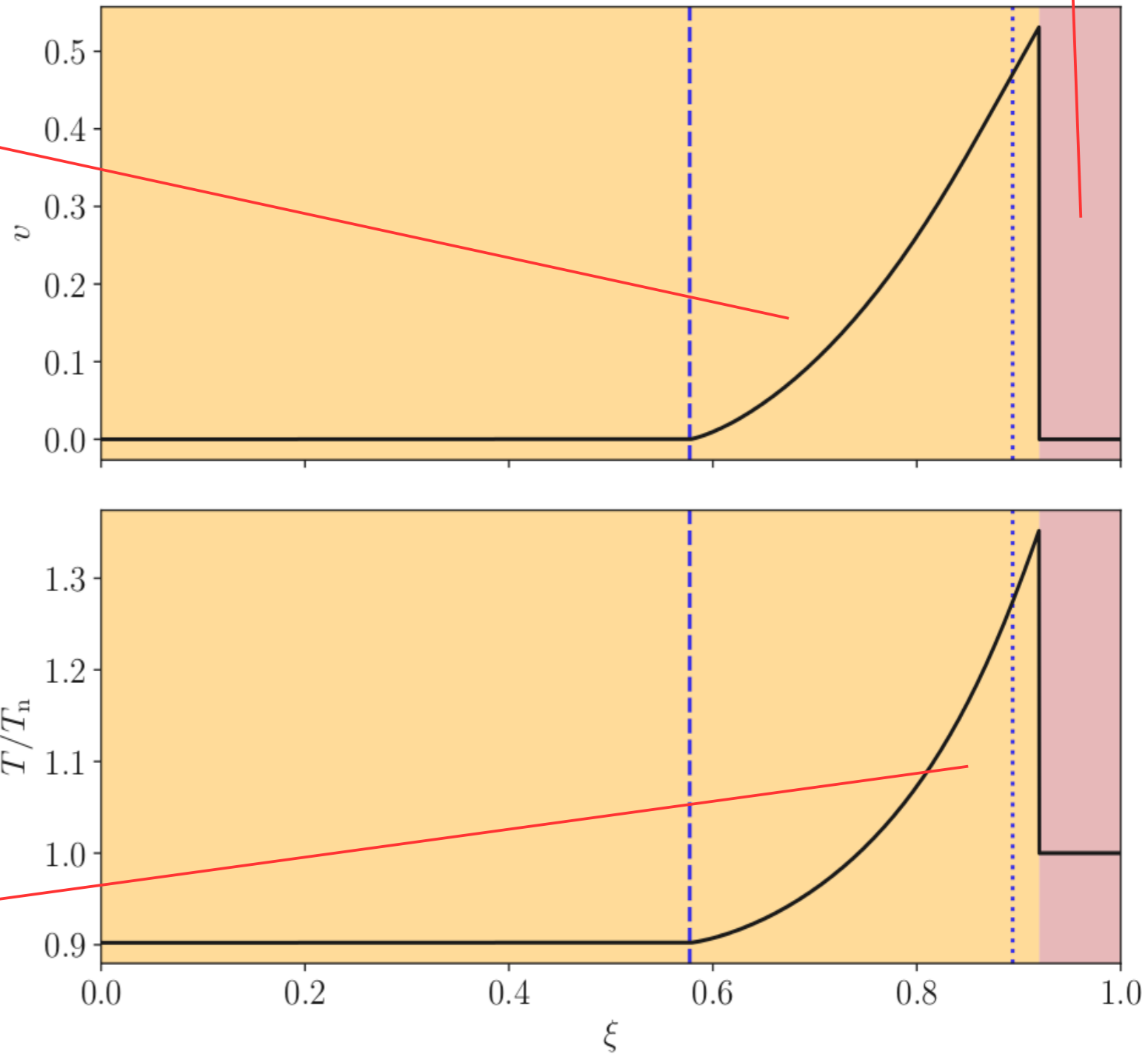
Espinosa, Konstandin, No, Servant, JCAP **1006** (2010) 028



Detonations ($v_w > v_J > c_s$)

fluid at rest in front of bubble

fluid behind phase boundary gets dragged, until it becomes at rest again

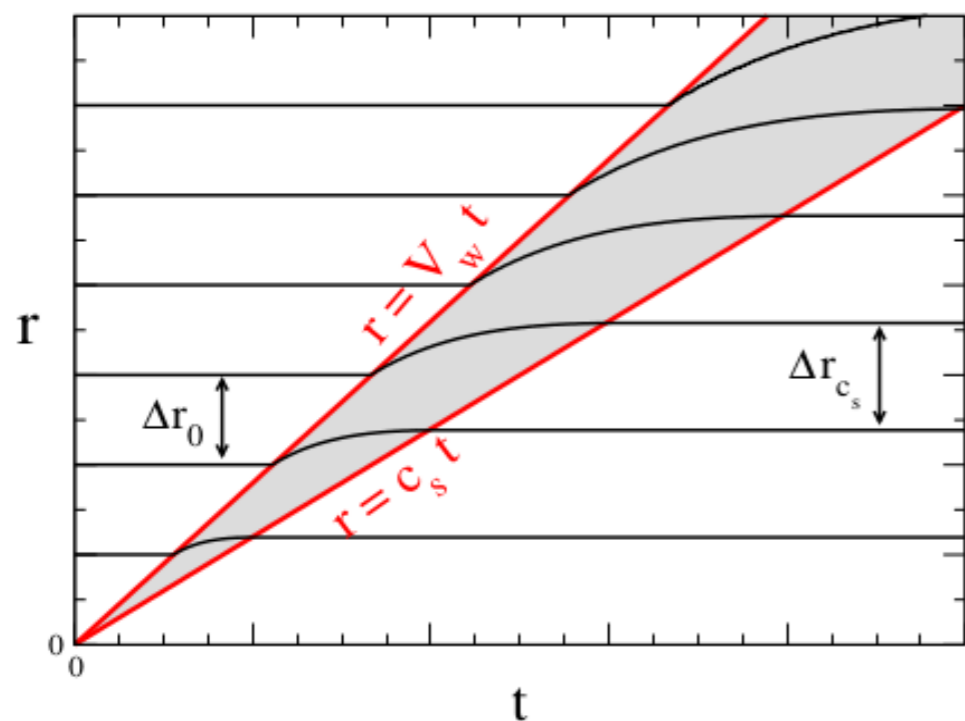


fluid heated inside bubble

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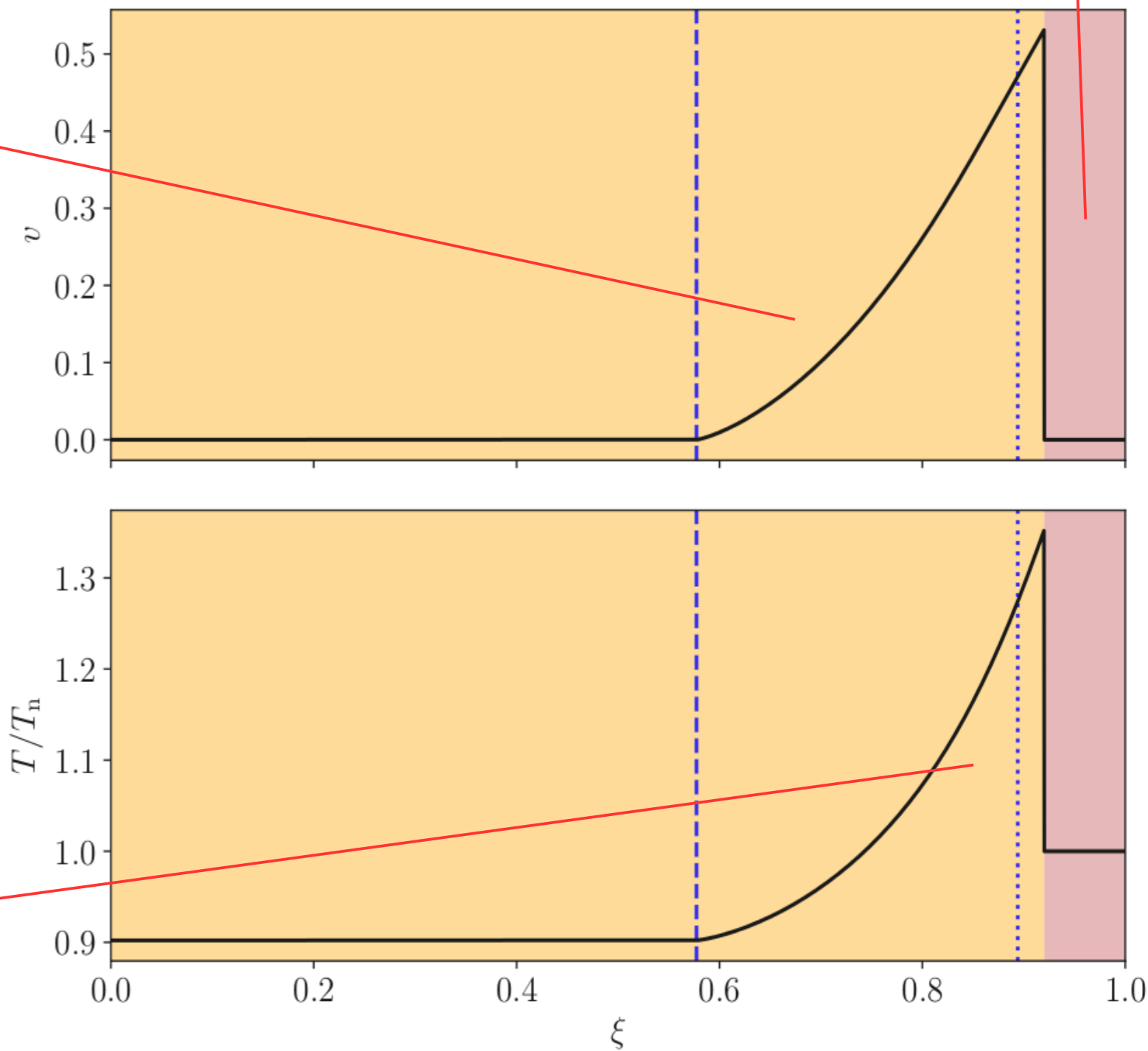
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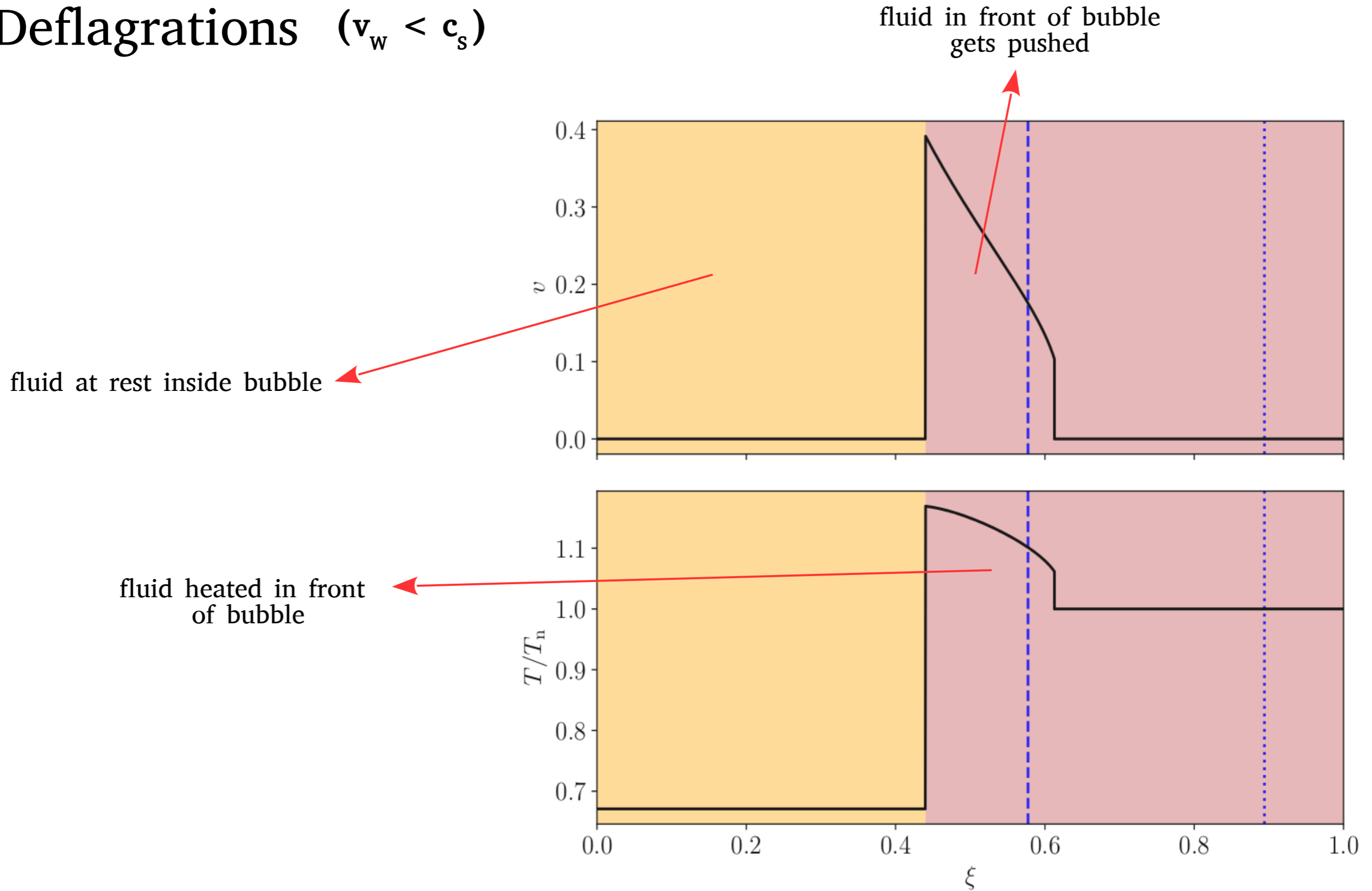
Caprini, No, JCAP **1201** (2012) 031

fluid heated inside bubble



Courtesy of D. Cutting (Sussex)

Deflagrations ($v_w < c_s$)



Courtesy of D. Cutting (Sussex)

Fluid shell Profiles

$$\partial^\mu T_{\mu\nu}^{\text{plasma}} = 0 \quad (\text{with appropriate boundary conditions on bubble wall})$$

Local Thermal Equilibrium

$$T_{\mu\nu}^{\text{plasma}} = w u_\mu u_\nu - g_{\mu\nu} p$$

$$w = e + p$$

$$u_\mu = \frac{(1, \mathbf{v})}{\sqrt{1 - \mathbf{v}^2}} = (\gamma, \gamma \mathbf{v})$$

Self-similarity

$$v(r, t) = v(\xi = r/t)$$

Estimate of Energy available for GW production (fluid bulk motion for one bubble)

$$\overline{U}_f^2 = \frac{3}{e v_w^3} \int w(\xi) v^2 \gamma^2 \xi^2 d\xi = \frac{\kappa \alpha}{1 + \alpha}$$

(enthalpy weighted) plasma
RMS four velocity

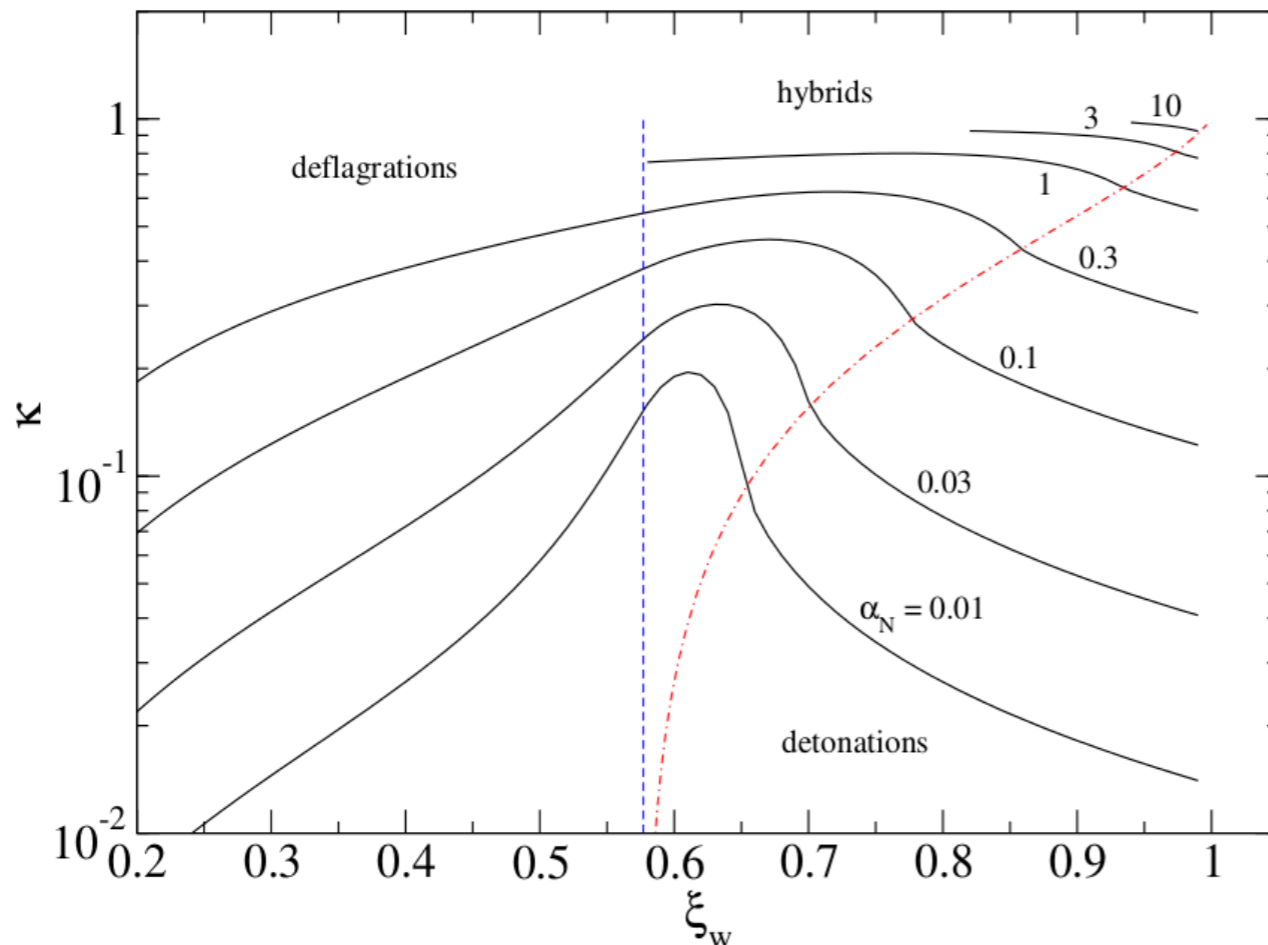
Hindmarsh, Huber, Rummukainen, Weir, PRD 96 (2017) 103520

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Hindmarsh, Huber, Rummukainen, Weir, PRD 96 (2017) 103520



Efficiency coefficient
(PT Energy Budget)

Kamionkowski, Kosowsky, Turner, PRD 49 (1994) 2837

Espinosa, Konstandin, No, Servant, JCAP 1006 (2010) 028

Gravitational Waves from Phase Transitions

- Gravitational waves (GWs) produced by several sources in a PT:

$$h^2\Omega_{\text{gw}} = h^2\Omega_{\phi} + h^2\Omega_{\text{sw}} + h^2\Omega_{\text{turb}}$$

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LISA Cosmology Working Group effort to provide state-of-art:

2015 CosWG Review

*Caprini et al, JCAP **1604** (2016) 001*

+ recent update

*Caprini et al, JCAP **2003** (2020) 024*

Science with the space-based interferometer eLISA.
II: Gravitational waves from cosmological phase transitions

Chiara Caprini^a, Mark Hindmarsh^{b,c}, Stephan Huber^b,
Thomas Konstandin^d, Jonathan Kozaczuk^e, Germano Nardini^f,
Jose Miguel No^b, Antoine Petiteau^g, Pedro Schwaller^d,
Géraldine Servant^{d,h}, David J. Weirⁱ



Detecting gravitational waves from cosmological phase transitions with LISA: an update

Chiara Caprini^a, Mikael Chala^{b,c,†}, Glauber C. Dorsch^d, Mark Hindmarsh^{e,f},
Stephan J. Huber^f, Thomas Konstandin^{g,‡}, Jonathan Kozaczuk^{h,i,j,§},
Germano Nardini^k, Jose Miguel No^{l,m}, Kari Rummukainen^e, Pedro Schwallerⁿ,
Geraldine Servant^{g,o}, Anders Tranberg^k, David J. Weir^{e,p,¶}

For the LISA Cosmology Working Group

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- $h^2\Omega_{\phi}$ sourced by collisions of bubble walls

*Kosowsky, Turner, Watkins, PRL **69** (1992) 2026; PRD **45** (1992) 4514*

*Huber, Konstandin, JCAP **0809** (2008) 022*

*Weir, PRD **93** (2016) 124037*

*Cutting, Hindmarsh, Weir, PRD **97** (2018) 123513*

In general, negligible expect for very strong supercooling $\Rightarrow \alpha \gg 1$

Gravitational Waves from Phase Transitions

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*Kosowsky, Turner, Watkins, PRL **69** (1992) 2026; PRD **45** (1992) 4514*

*Huber, Konstandin, JCAP **0809** (2008) 022*

*Weir, PRD **93** (2016) 124037*

*Cutting, Hindmarsh, Weir, PRD **97** (2018) 123513*

In general, negligible expect for very strong supercooling $\Rightarrow \alpha \gg 1$

Such amount of supercooling incompatible with PT completion...

*Ellis, Lewicki, No, JCAP **1904** (2019) 003*

...except for conformal scalar potentials

Gravitational Waves from Phase Transitions

- Gravitational waves (GWs) produced by several sources in a PT:

$$h^2\Omega_{\text{gw}} = h^2\Omega_{\phi} + h^2\Omega_{\text{sw}} + h^2\Omega_{\text{turb}}$$

- $h^2\Omega_{\text{sw}}$ sourced by plasma sounds waves (longitudinal modes)

*Hindmarsh, Huber, Rummukainen, Weir, PRL **112** (2014) 041301; PRD **92** (2015) 123009; PRD **96** (2017) 103520*

*Hindmarsh, PRL **120** (2018) 071301*

*Konstandin, JCAP **1803** (2018) 047*

Hindmarsh, Hijazi, arXiv:1909.10040

Typically dominant signal

GW power spectrum (numerical simulations)

$$\frac{d\Omega_{\text{gw},0}}{d\ln(f)} = 0.687 F_{\text{gw},0} K^2 (H_* R_* / c_s) \tilde{\Omega}_{\text{gw}} C \left(\frac{f}{f_{\text{p},0}} \right)$$

Peak amplitude
(at maximum)

Peak frequency
(at maximum)

Spectral Shape
(Broken Power Law)

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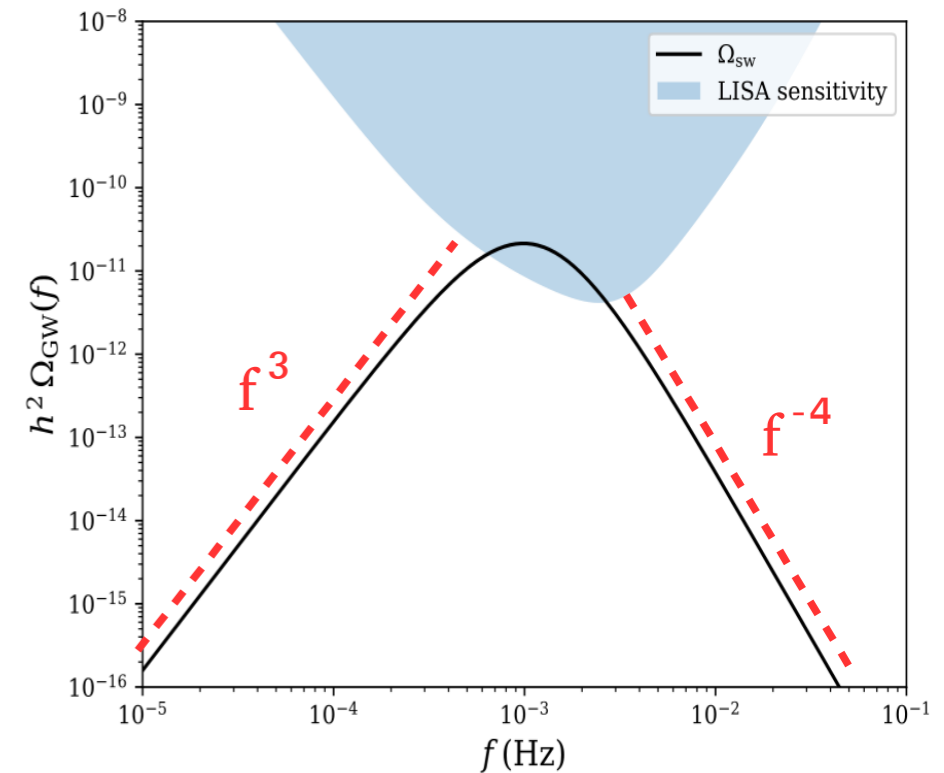
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Peak amplitude
(at maximum)

Peak frequency
(at maximum)

**Spectral Shape
(Broken Power Law)**

$$C(s) = s^3 \left(\frac{7}{4 + 3s^2} \right)^{\frac{7}{2}}$$



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$$F_{\text{gw},0} = (3.57 \pm 0.05) \times 10^{-5} \left(\frac{100}{g_*} \right)^{\frac{1}{3}}$$

$$\tilde{\Omega}_{\text{gw}} \sim 10^{-2}$$

(from simulation)

$$K = \frac{\langle w \gamma^2 v^2 \rangle}{\bar{e}} = \Gamma \bar{U}_f^2$$

Kinetic Energy Fraction

Gravitational Waves from Phase Transitions

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Typically dominant signal

After $\tau_{\text{sh}} \sim L_f / \bar{U}_f$, fluid becomes nonlinear (shock formation)

characteristic fluid length scale \swarrow \Downarrow Sound wave GW source shuts-off

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$$H_* \tau_{\text{sh}} = H_* R_* / K^{1/2} < 1$$

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*Hindmarsh, PRL **120** (2018) 071301*

*Konstandin, JCAP **1803** (2018) 047*

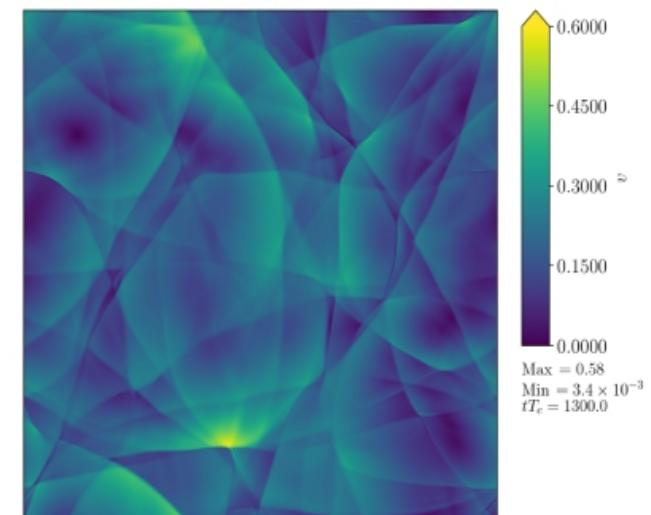
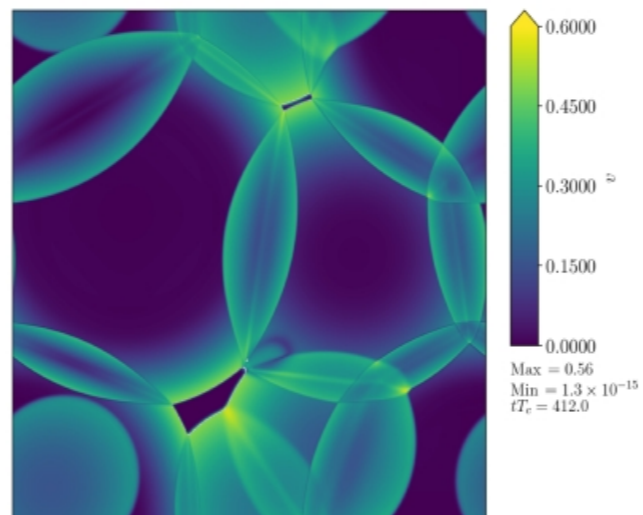
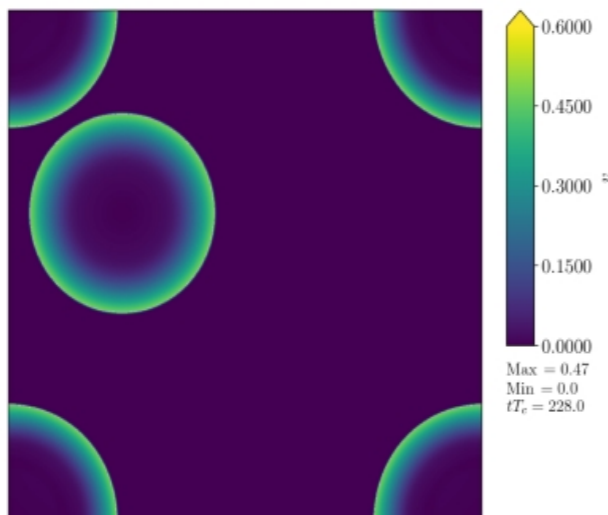
Hindmarsh, Hijazi, arXiv:1909.10040

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*Konstandin, JCAP **1803** (2018) 047*

Hindmarsh, Hijazi, arXiv:1909.10040

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Sound wave GW source shuts-off

- $h^2\Omega_{\text{turb}}$ sourced by plasma turbulence (vortical modes)

*Gogoberidze, Kahniashvili, Kosowsky, PRD **76** (2007) 083002*

*Caprini, Durrer, Servant, JCAP **0912** (2009) 024*

Roper Pol, Mandal, Brandenburg, Kahniashvili, Kosowsky, arXiv:1903.08585

→ Turbulent flow expected to develop when sound waves shut-off

→ **Vorticity can also coexist with sound waves for deflagrations and $\alpha > 0.1$**

Cutting, Hindmarsh, Weir, arXiv:1906.00480

Gravitational Waves from Phase Transitions

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→ Turbulence is expected to develop when sound waves shut-off

→ Vorticity can also coexist with sound waves for reionizations and $\alpha > 0.1$

Cutting, Hindmarsh, Weir, arXiv:1906.00480

Numerical simulations still ongoing

Gravitational Waves from Phase Transitions

Duration of sound wave GW source

Initially assumed linear fluid regime lasts approx. a Hubble time

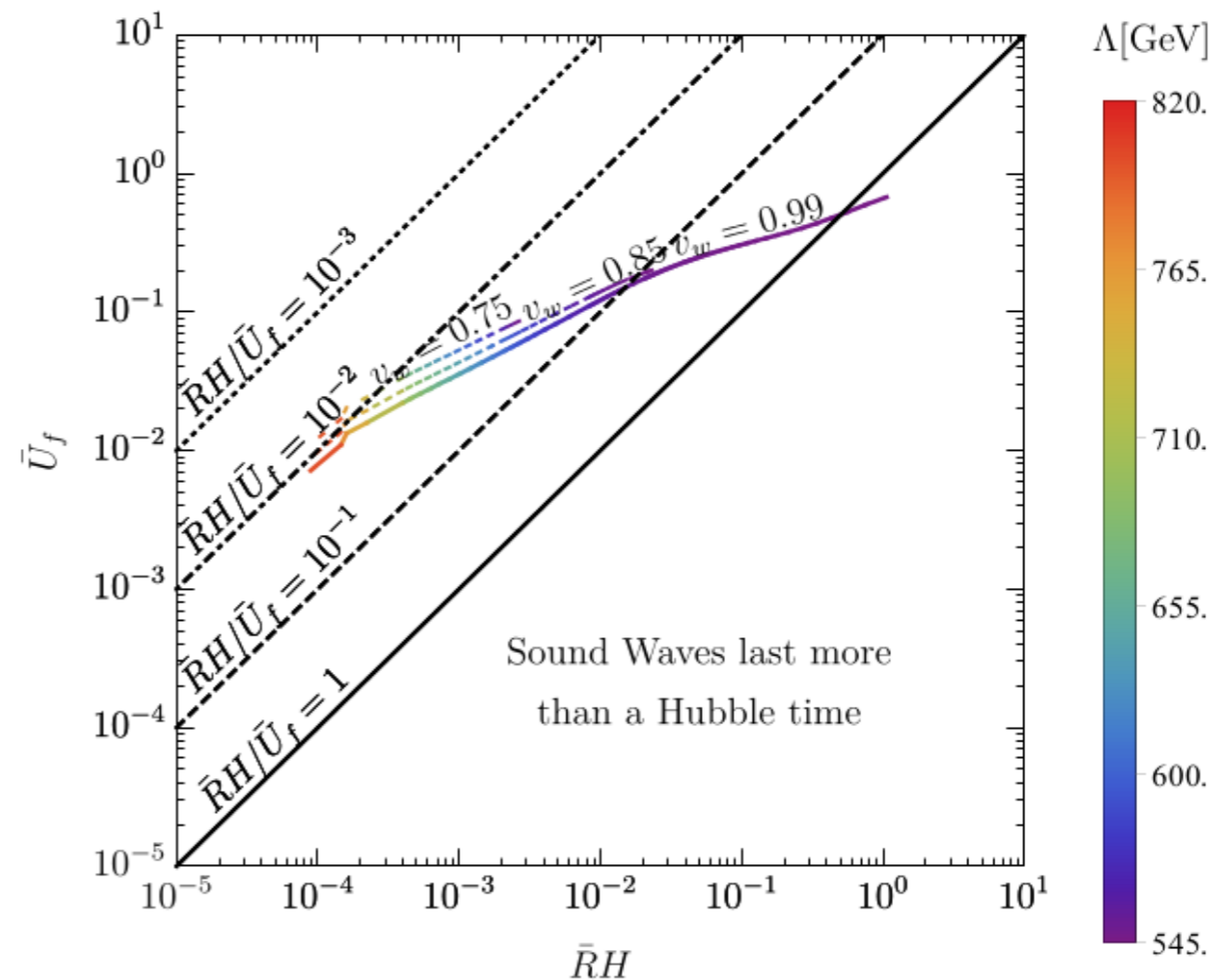
$$\tau_{\text{sh}} \gtrsim H_*^{-1}$$

But non-linearities generally “cut short” the sound wave GW source:

Ellis, Lewicki, No, JCAP 1904 (2019) 003

Concrete BSM example:

$$V(H) = -m^2|H|^2 + \lambda|H|^4 + \frac{1}{\Lambda^2}|H|^6$$



Gravitational Waves from Phase Transitions

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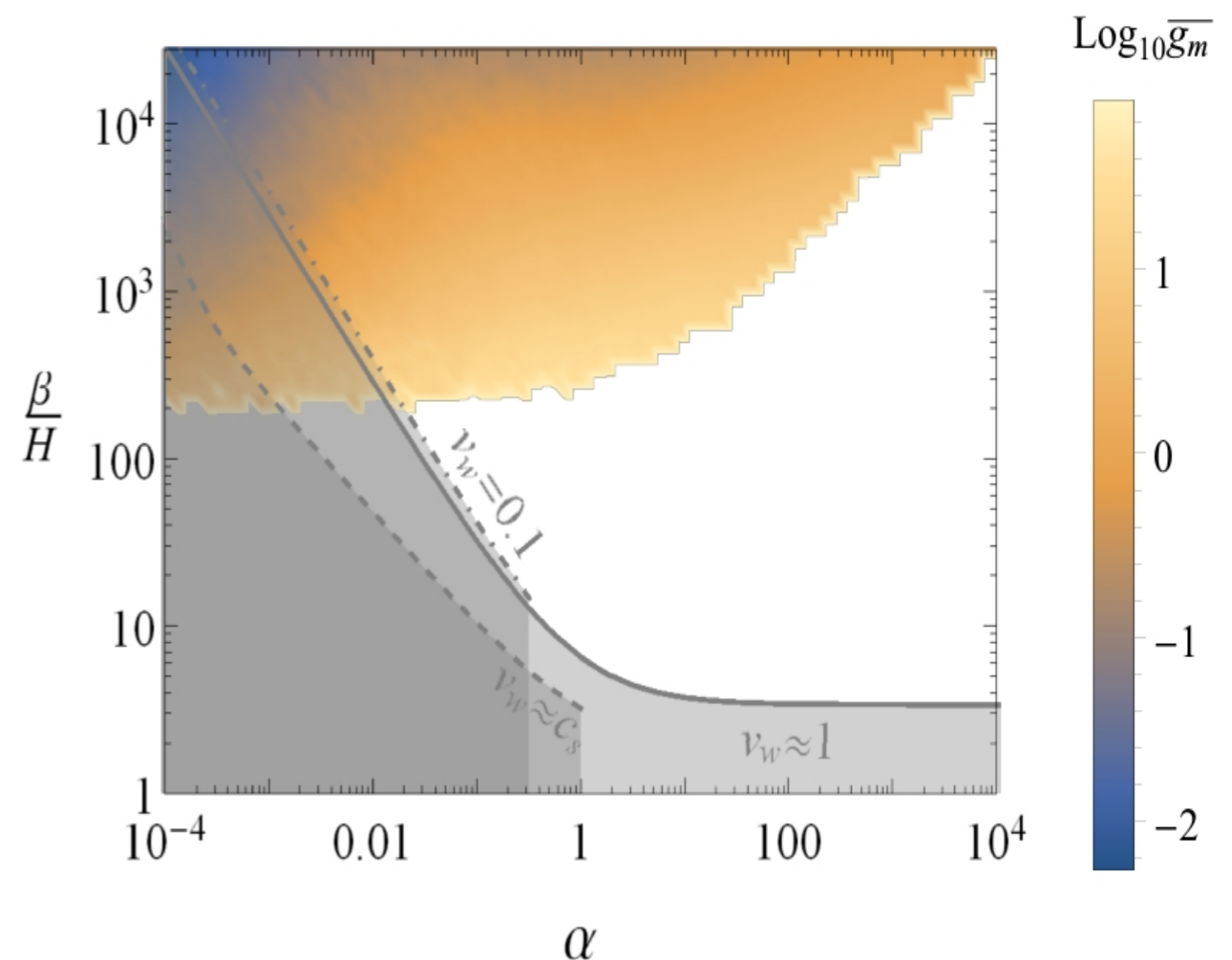
*Ellis, Lewicki, No, JCAP **1904** (2019) 003*

General analysis

*Ellis, Lewicki, No, JCAP **2007** (2020) 050*

$$V(\phi, T) = \frac{g_m^2}{24} (T^2 - T_0^2) \phi^2 - \frac{g_m}{12\pi} T \phi^3 + \lambda \phi^4$$

(purely thermal potential barrier)

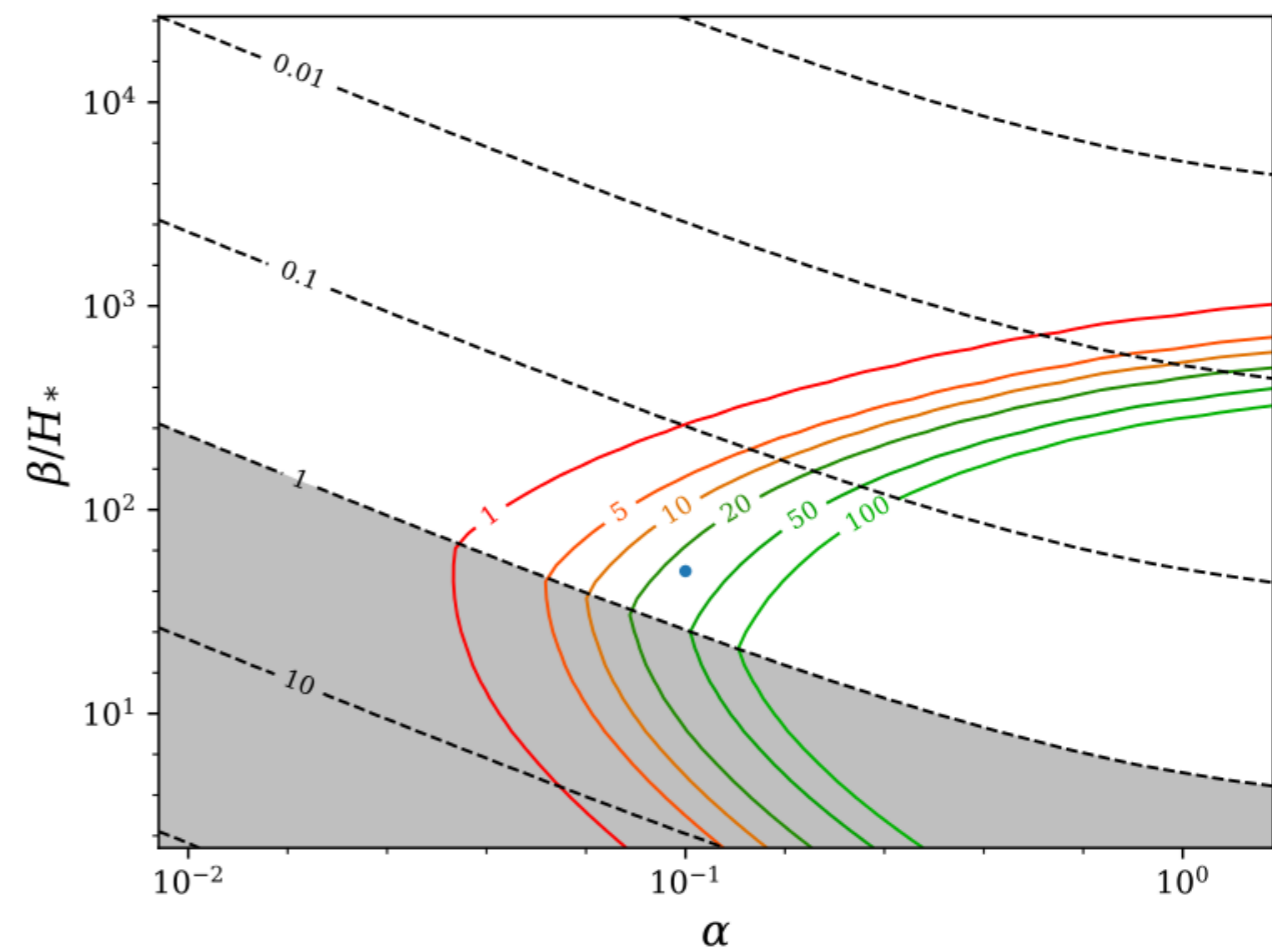
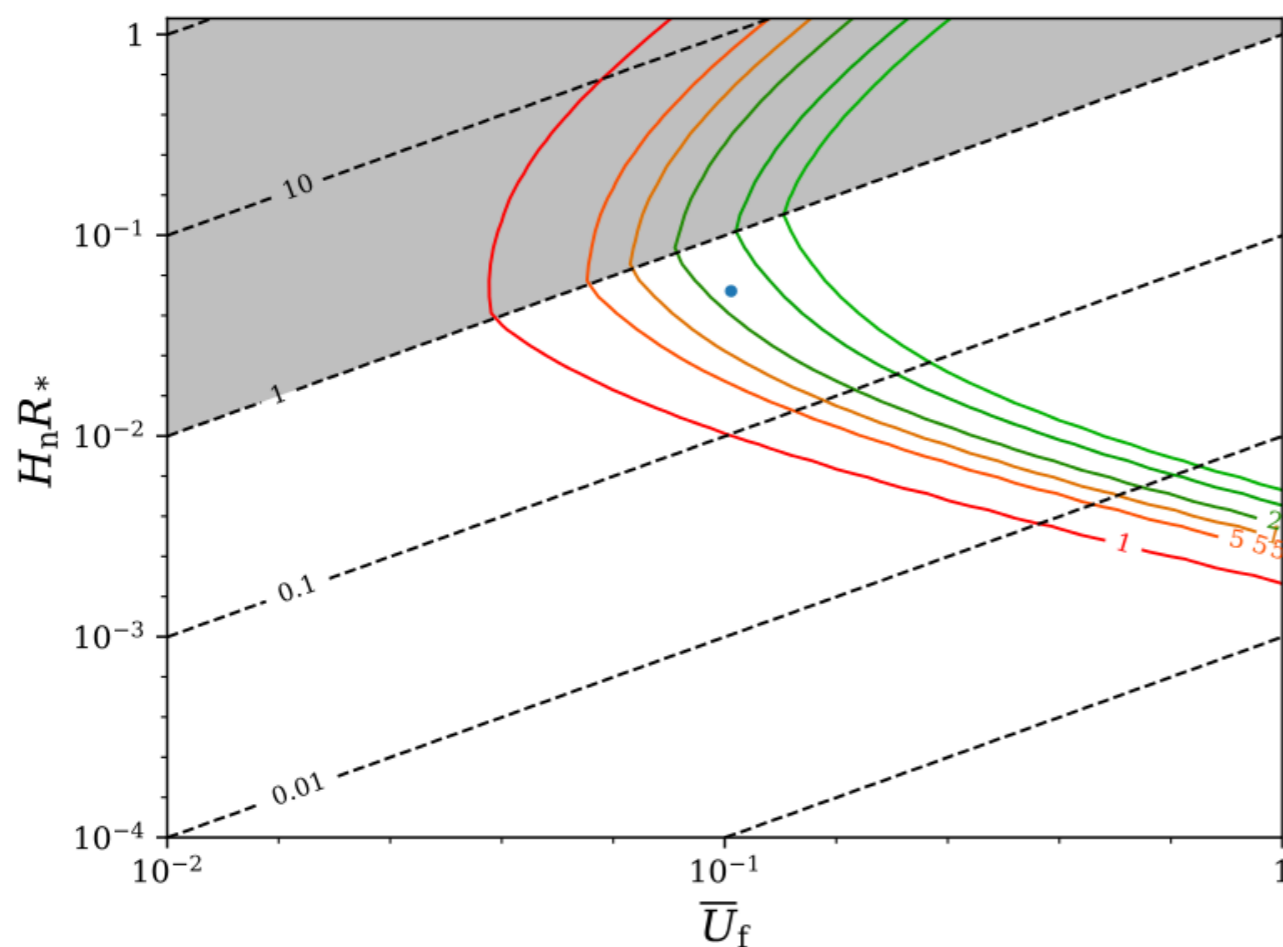


Gravitational Waves from Phase Transitions

LISA signal to noise

$$\text{SNR} = \sqrt{\mathcal{T} \int_{f_{\min}}^{f_{\max}} df \left[\frac{h^2 \Omega_{\text{GW}}(f)}{h^2 \Omega_{\text{Sens}}(f)} \right]^2}$$

$$h^2 \Omega_{\text{Sens}}(f) = \frac{2\pi^2}{3H_0^2} f^3 S_h(f)$$

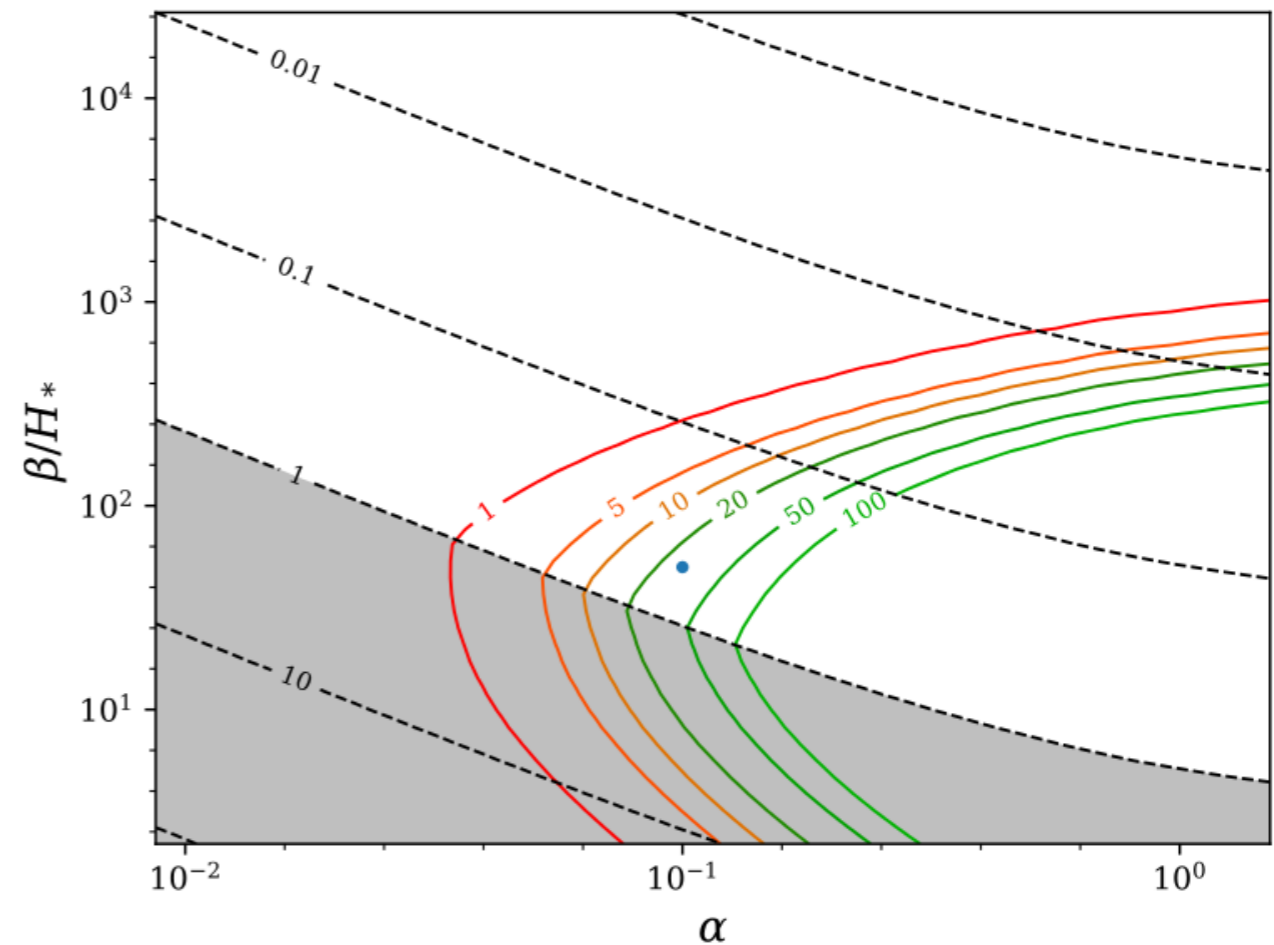
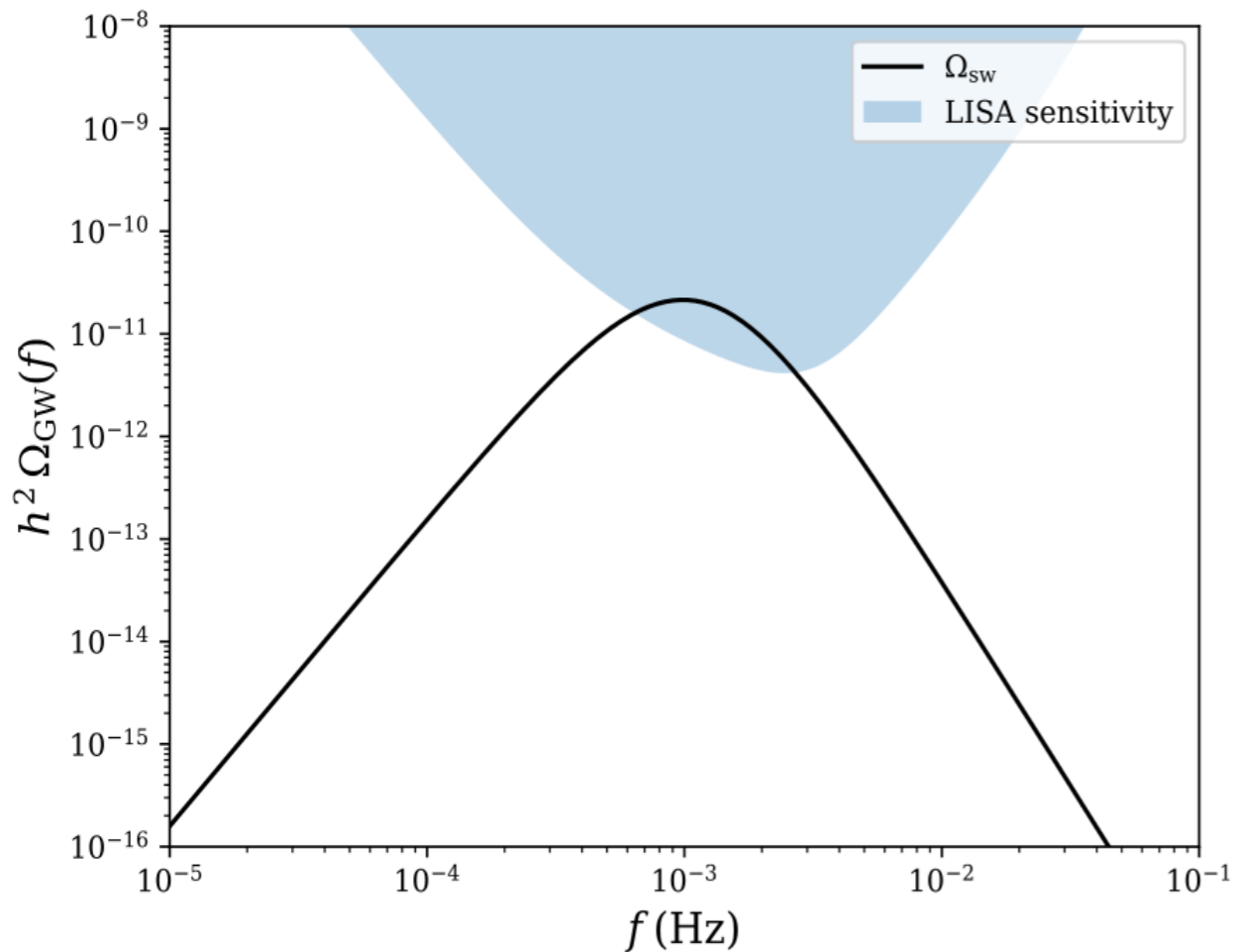


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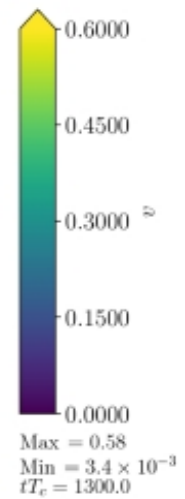
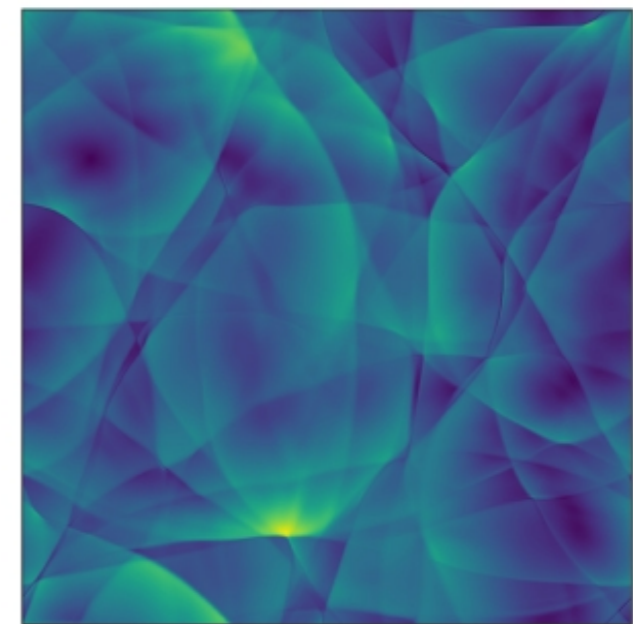
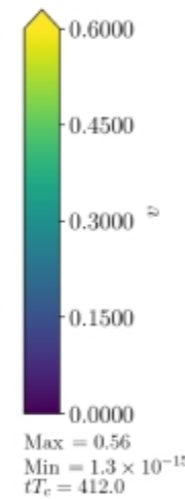
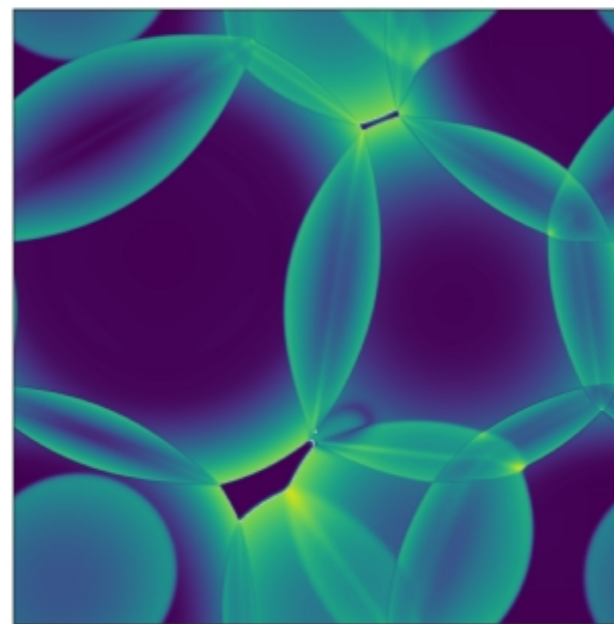
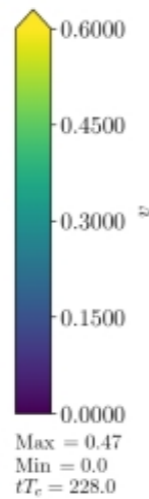
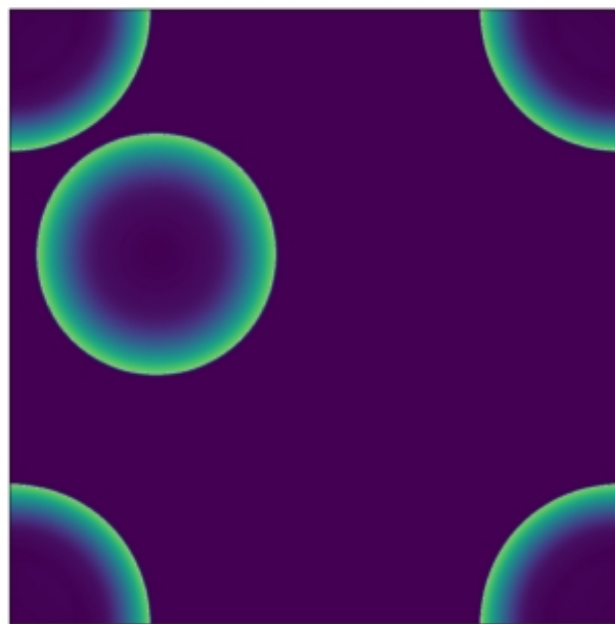


Gravitational Waves from Phase Transitions

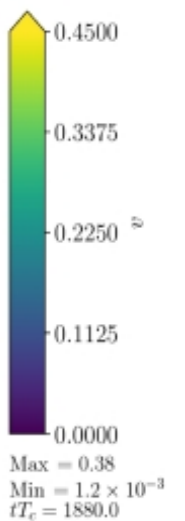
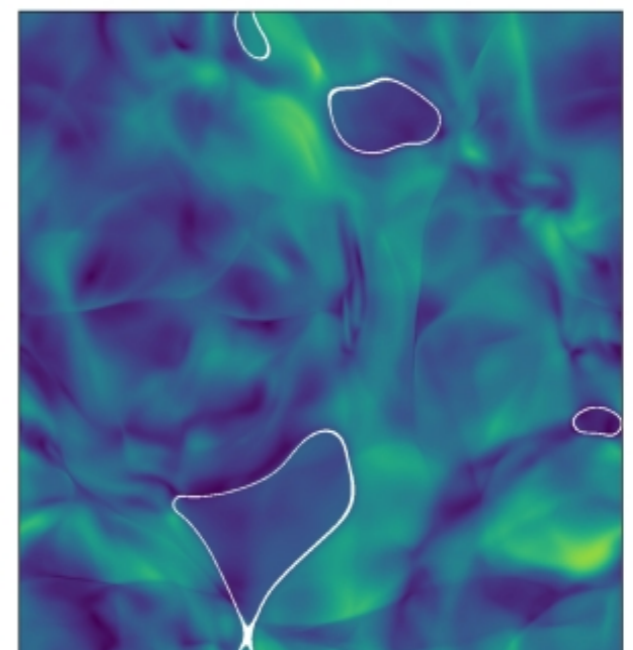
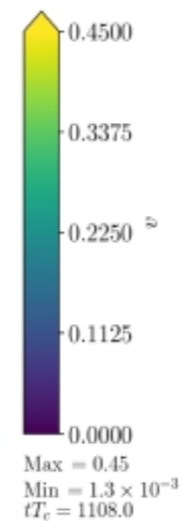
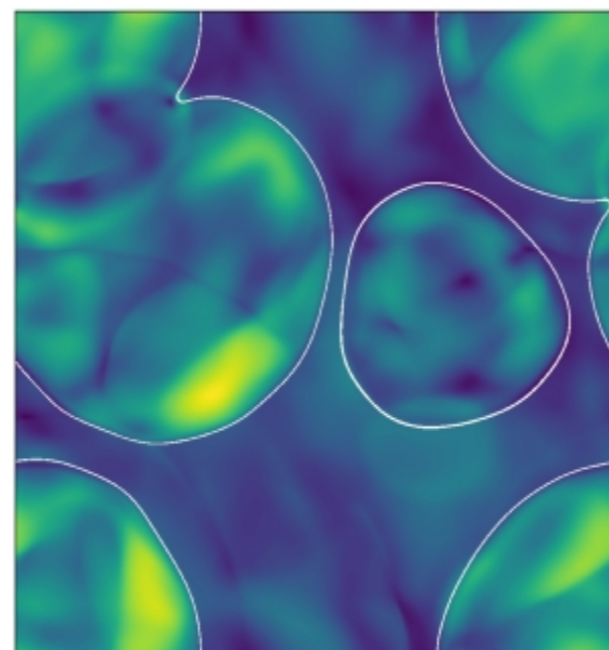
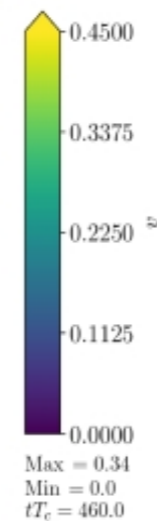
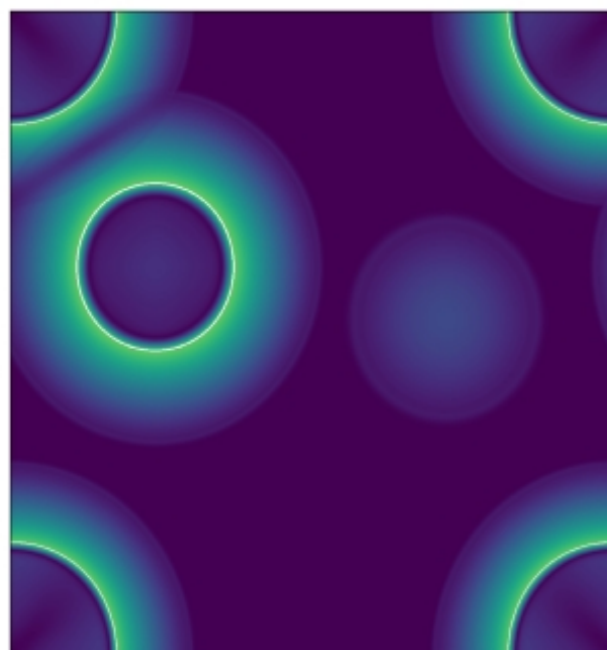
Understanding of vorticity generation is ongoing...

Cutting, Hindmarsh, Weir, arXiv:1906.00480

Detonations ($\alpha > 0.1$)



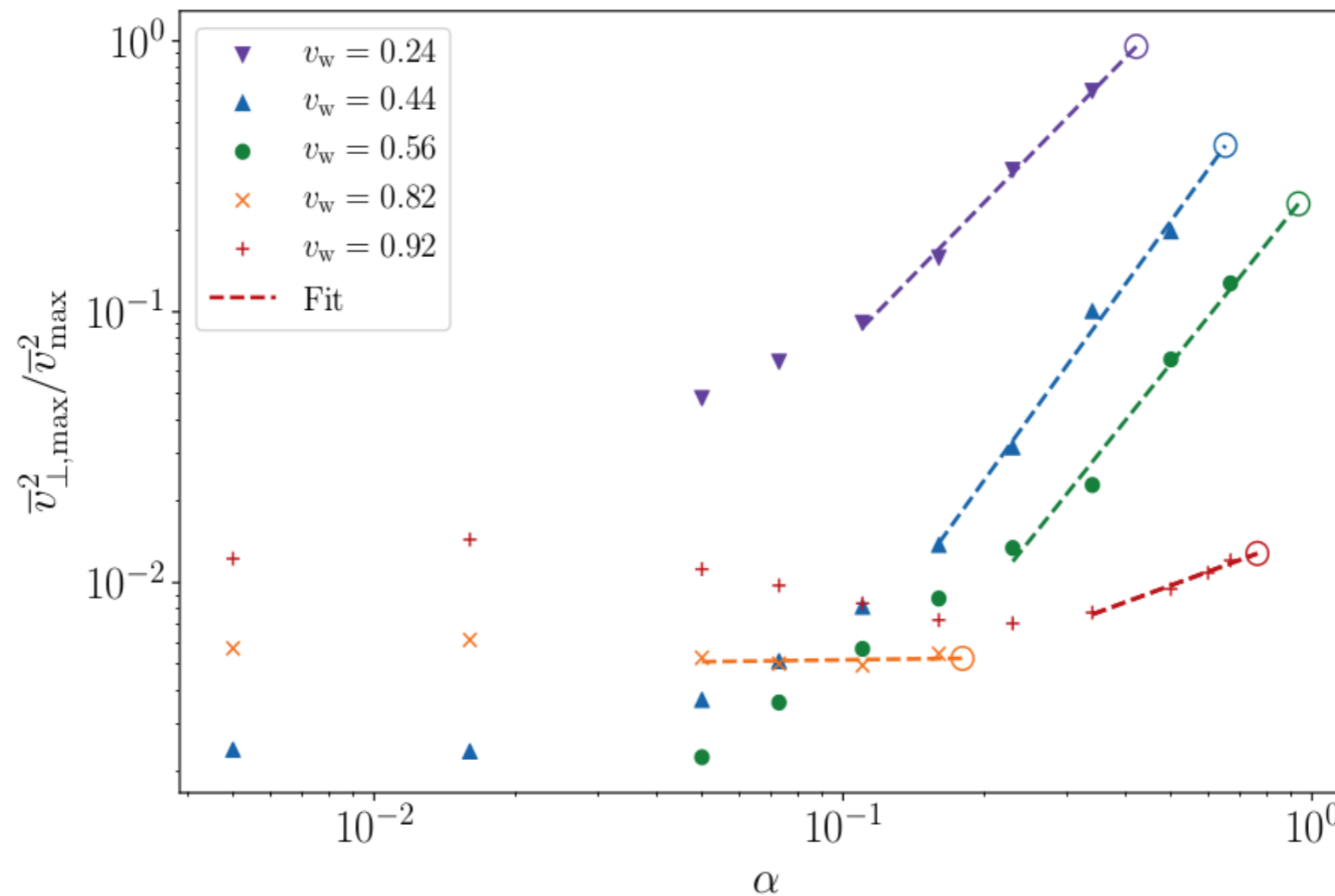
Deflagrations ($\alpha > 0.1$)



Gravitational Waves from Phase Transitions

Understanding of vorticity generation is ongoing...

Cutting, Hindmarsh, Weir, arXiv:1906.00480



Deflagrations with large α (> 0.1) generate significant vorticity coexisting with sound waves!

In the last
couple of
min...

GW generation vs EW Baryogenesis in 1st Order EW Phase Transition

GWs: Sizable plasma bulk motion \Rightarrow **Sizable** v_w

EWBG: Velocities $\sim 0.05 - 0.1$ preferred
(efficient transport)

Incompatible?

GW generation vs EW Baryogenesis in 1st Order EW Phase Transition

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Incompatible?



NO!

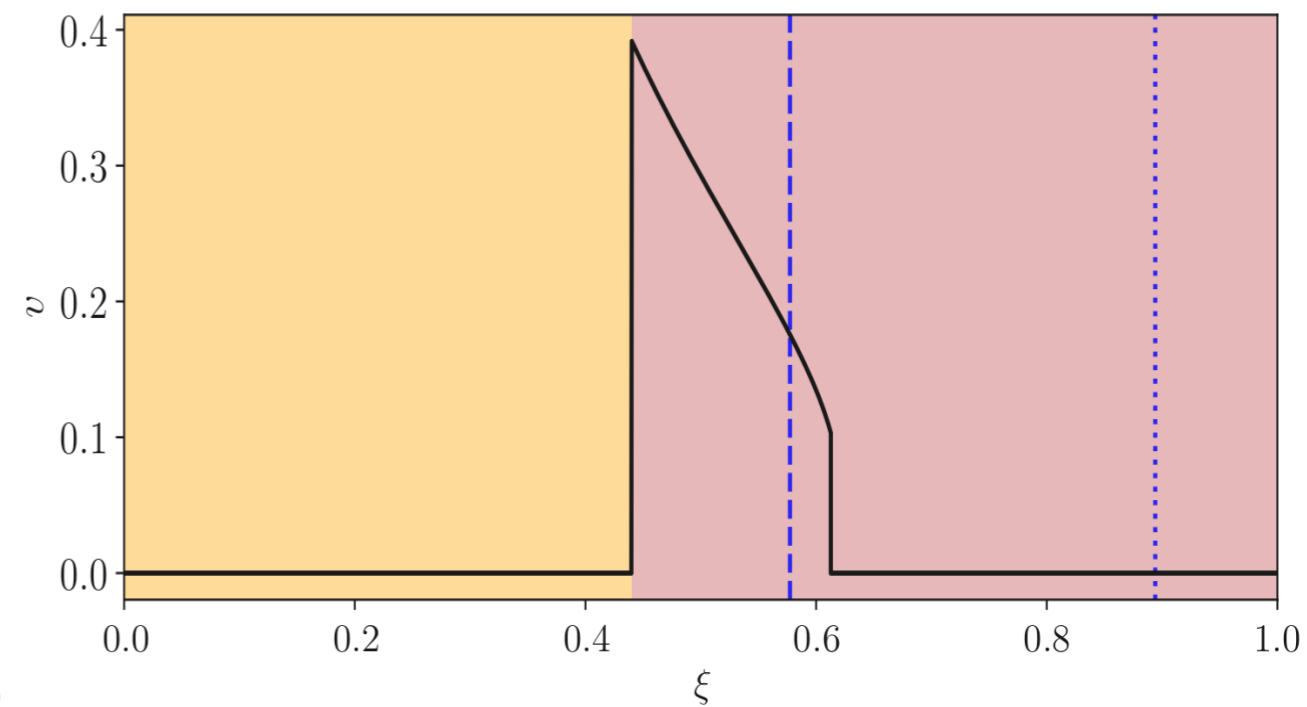
Relevant velocities are not the same...
(for deflagrations)

No, PRD 84 (2011) 124025

GWs: Fluid velocity (bubble rest frame)

EWBG: Fluid velocity (bubble wall rest frame)

(relative velocity between bubble wall and plasma in front)



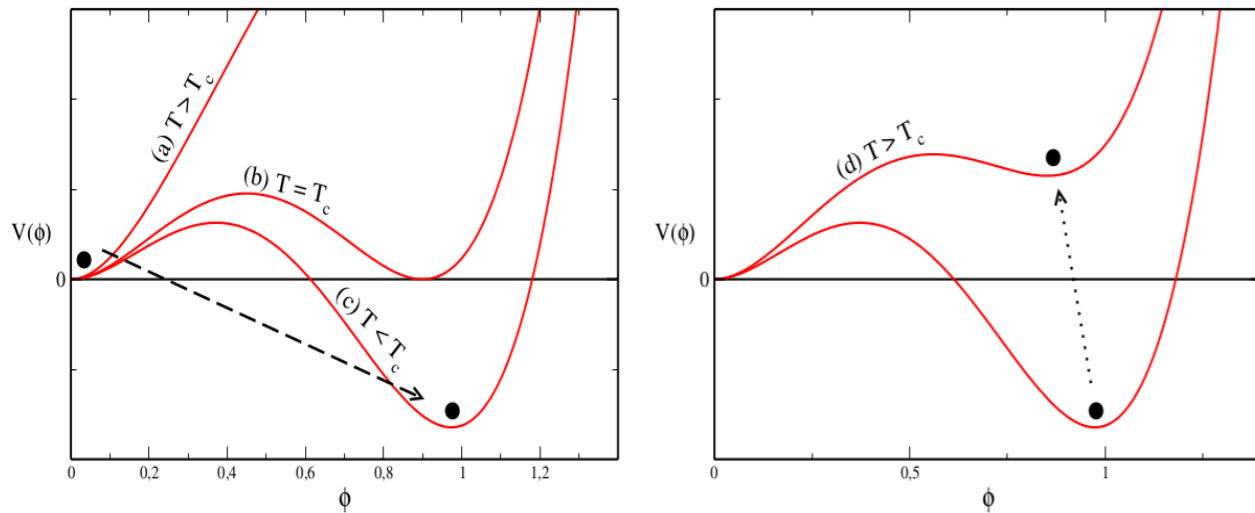
GW generation vs EW Baryogenesis in 1st Order EW Phase Transition

For detonations: EWBG would naively not work (inefficient transport)

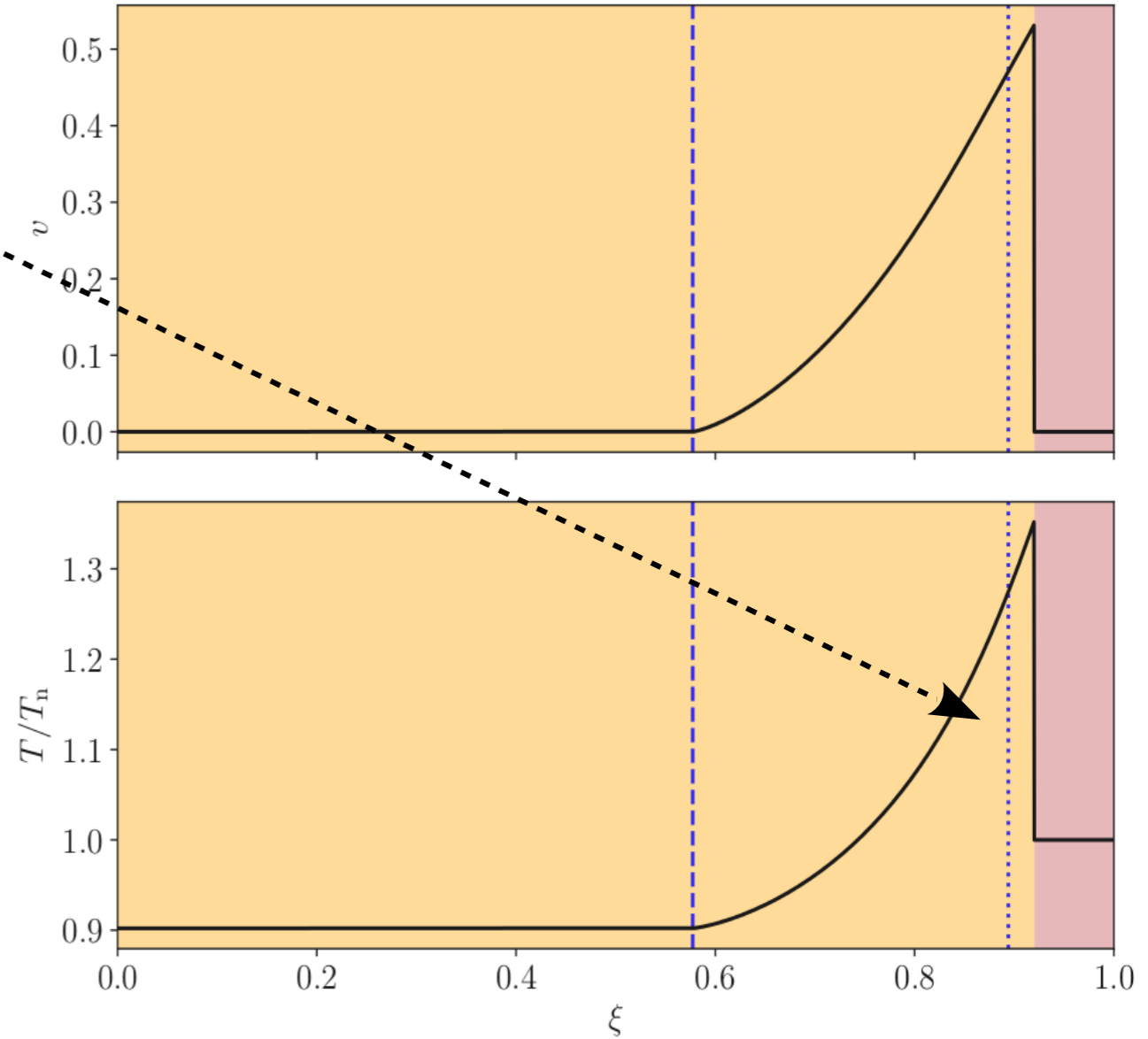
However... for detonations plasma is reheated behind bubble wall

Possible to do EWBG from local back-tunneling (if $T > T_c$) in broken phase due to reheating!

Caprini, No, JCAP 1201 (2011) 031



“Supersonic EWBG”



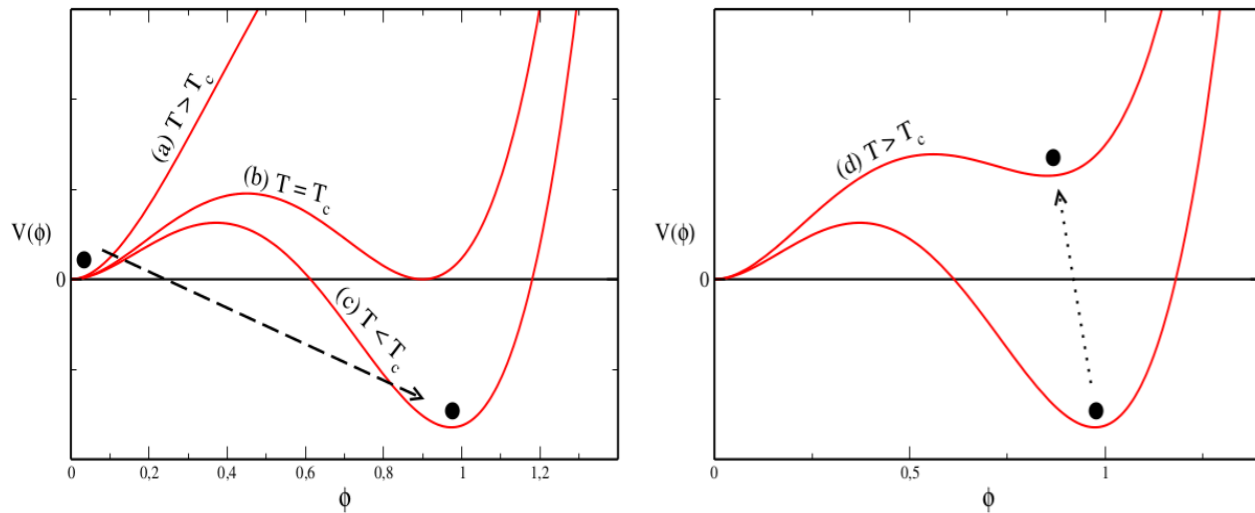
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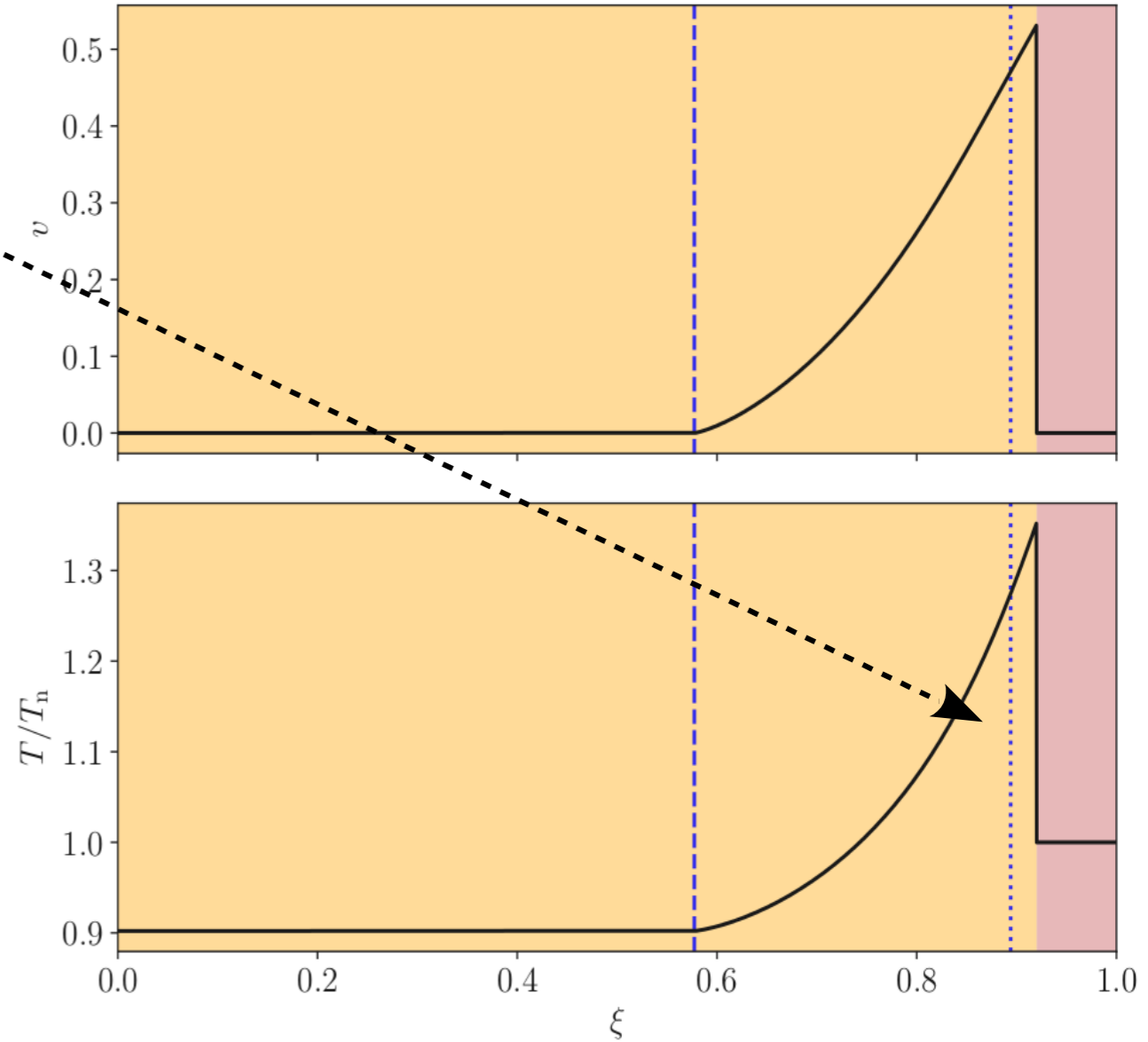
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Caprini, No, JCAP **1201** (2011) 031



“Supersonic EWBG”



Recent works suggest transport may also be efficient at high v_w

Cline, Kainulainen, PRD **101** (2020) 063525

Thank you!



KEEP
CALM
AND
BACKUP
YOUR
WORK

Higgs Evolution in Early Universe



FINITE-TEMPERATURE EFFECTIVE POTENTIAL

$$V_{\text{eff}}(h, T) = V_0(h) + V_0^{\text{loop}}(h) + V_T(h, T)$$

Tree-level potential

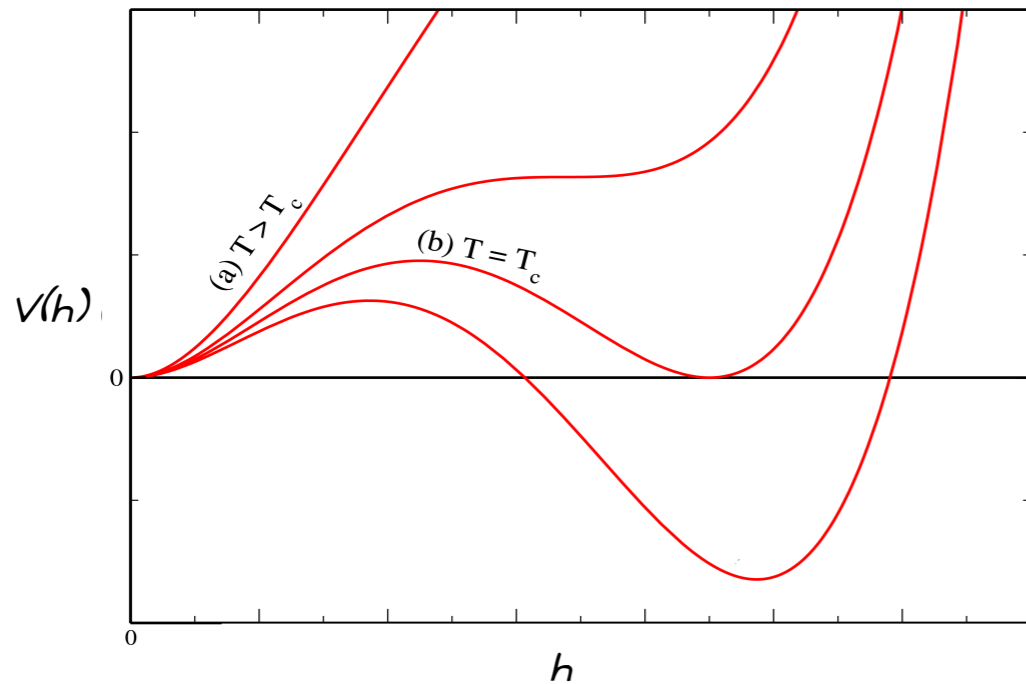
Loop corrections

Thermal corrections

(Perturbative) Nature of EWPT

1st Order:

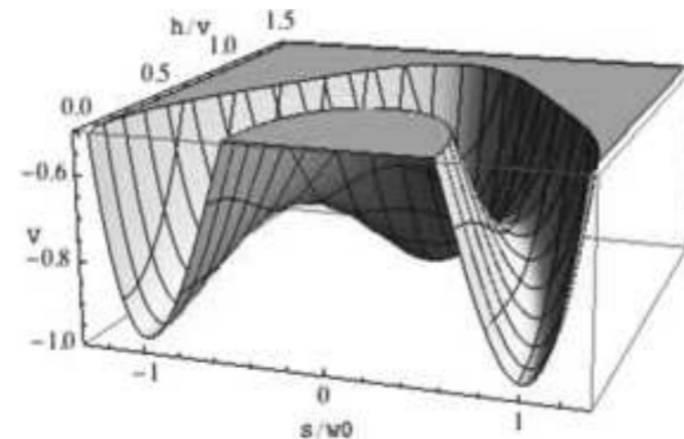
$\langle h \rangle = 0 \rightarrow \langle h \rangle = h(T)$ Discontinuous



Non-analytic term $(m^2)^{3/2}$ in $V(h, T)$ from Matsubara Zero-modes

(only present for bosons)

Multiple fields involved in the EWPT may allow for tree-level potential barrier

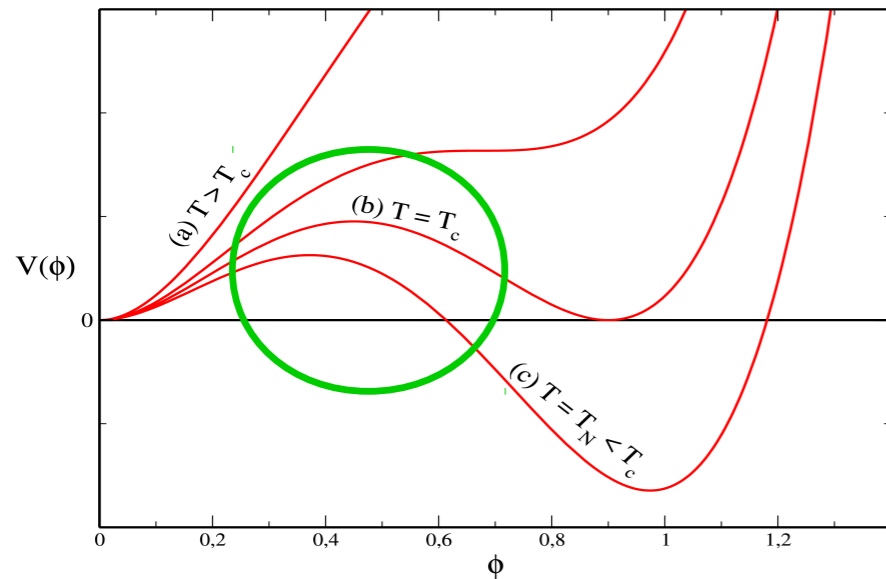


BSM: New Physics sizeably coupled to Higgs can drastically change the EWPT nature

- ▶ New Physics should induce deviations in Higgs couplings
- ▶ New Physics needed close to EW scale

Some further aspects of the EW Phase Transition

Effective Potential (finite T)



$$V_{\text{eff}}(h, T) = V_0(h) + V_0^{\text{loop}}(h) + V_T(h, T)$$

$$V_1^T(h, T) = \frac{T^4}{2\pi^2} \left[\sum_i \pm n_i J_{\pm} \left(\frac{m_i^2(h)}{T^2} \right) \right]$$

$$J_{\pm}(x) = \int_0^{\infty} dy y^2 \log \left[1 \mp \exp(-\sqrt{x^2 + y^2}) \right]$$

High-T expansion:

$$T^4 J_+ \left(\frac{m^2}{T^2} \right) = -\frac{\pi^4 T^4}{45} + \frac{\pi^2 m^2 T^2}{12} - \frac{T \pi (m^2)^{3/2}}{6} - \frac{(m^4)}{32} \log \frac{m^2}{a_b T^2}$$

$$T^4 J_- \left(\frac{m^2}{T^2} \right) = \frac{7\pi^4 T^4}{360} - \frac{\pi^2 m^2 T^2}{24} - \frac{(m^4)}{32} \log \frac{m^2}{a_f T^2},$$

$$V_{\text{eff}}(h, T) \approx (a T^2 - \mu^2) h^2 - E(T) h^3 + \lambda_{\text{eff}}(T) h^4$$