

SMEFT studies of LHC data: status and perspective

Ilaria Brivio

Institut für Theoretische Physik,
Universität Heidelberg



What is SMEFT?

SMEFT = **S**tandard **M**odel **E**ffective **F**ield **T**heory

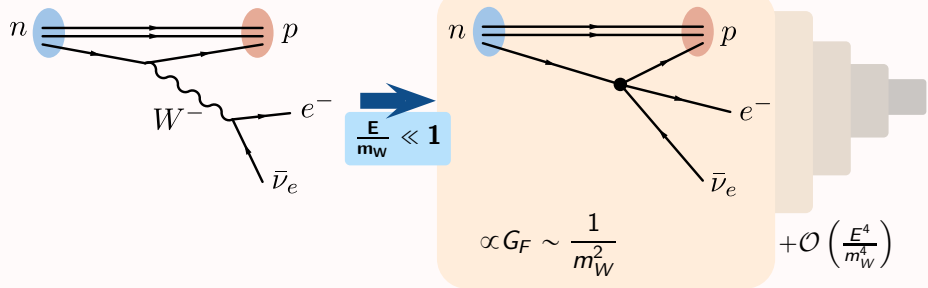
What is SMEFT?

SMEFT = **S**tandard **M**odel **E**ffective **F**ield **T**heory

What is an EFT?

a field theory that approximates another one in a specific kinematic/parameters limit

Example: Fermi interactions



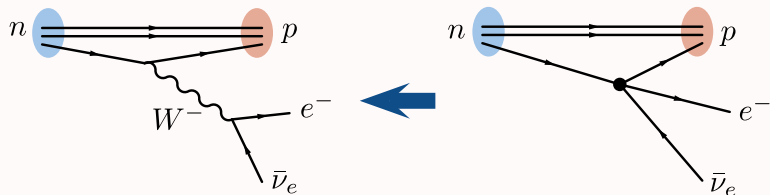
EFT = QFT valid in a regime $\mathbf{E}/m_W \ll 1$



Taylor expansion in \mathbf{E}/m_W

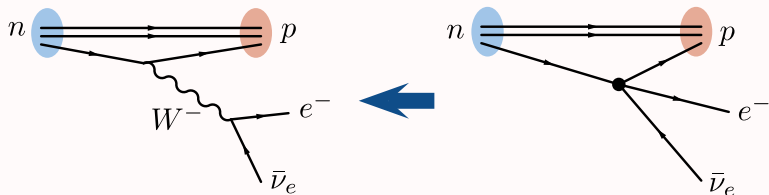
- ▶ **hierarchy** of physical effects determined by low-E fields and symmetries
 - ▶ **details** of underlying physics are irrelevant

Example: Fermi interactions



the EFT can also be formulated **from the bottom up**. One specifies:
fields & symmetries at the E of the measurement + **power counting**

Example: Fermi interactions



the EFT can also be formulated **from the bottom up**. One specifies:
fields & symmetries at the E of the measurement + **power counting**

$$\Rightarrow \mathcal{L}_{EFT} = \sum_i C_i \mathcal{O}_i$$

C_i = Wilson coefficients:

free, dimensionful parameters

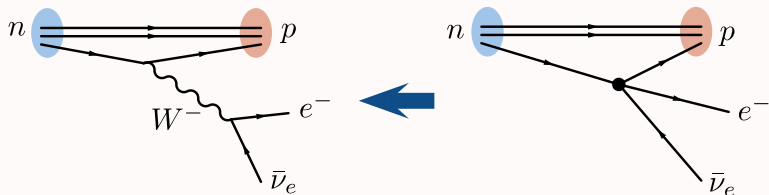
encode the (unknown) **UV** details

\mathcal{O}_i = operators

interactions terms classified in
expansion orders

encode the **IR** physics

Example: Fermi interactions



the EFT can also be formulated **from the bottom up**. One specifies:
fields & symmetries at the E of the measurement + **power counting**

$$\Rightarrow \mathcal{L}_{EFT} = \sum_i C_i \mathcal{O}_i$$

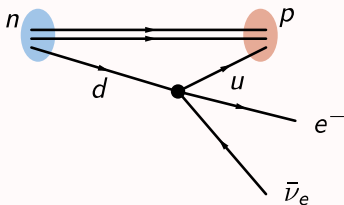
knowledge of the SM is **not required**.

the EFT can match **any model** leading to β decays.

Example: Fermi interactions, from the bottom up

1933 4-fermion interaction made of known fields ($+\nu$) and QED-invariant

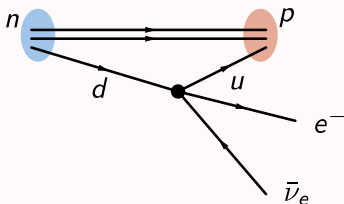
$$G_F(\bar{p}\gamma_\mu n)(\bar{e}\gamma_\mu \nu)$$



Example: Fermi interactions, from the bottom up

1933 4-fermion interaction made of known fields ($+\nu$) and QED-invariant

$$G_F(\bar{p}\gamma_\mu n)(\bar{e}\gamma_\mu\nu)$$



- ▶ energy spectra $\Rightarrow G_F \simeq (290 \text{ GeV})^{-2} \simeq v^{-2} \rightarrow$ scale of underlying physics!
- ▶ all β -decays have the same G_F
- ▶ angular distributions \Rightarrow the currents have *left-handed* chirality

\rightarrow very strong hints at the nature of EW interactions

SMEFT is the new Fermi theory

Standard **M**odel **E**ffective **F**ield **T**heory:
The EFT constructed with **Standard Model** field & symmetries

→ expansion in **canonical dimensions** d :

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda} \mathcal{L}_5 + \frac{1}{\Lambda^2} \mathcal{L}_6 + \frac{1}{\Lambda^3} \mathcal{L}_7 + \frac{1}{\Lambda^4} \mathcal{L}_8 + \dots$$

$$\mathcal{L}_d = \sum_i C_i \mathcal{O}_i^{(d)}$$

Introducing explicitly a new-physics scale parameter Λ , all C_i are dimensionless.

→ Taylor series in v/Λ or E/Λ (multi-scale!)

At each order, $\mathcal{O}_i^{(d)}$ form a complete, non-redundant **basis**

First relevant order: \mathcal{L}_6 → up to 2499 parameters

d=6: the Warsaw basis

Grzadkowski, Iskrzynski, Misiak, Rosiek 1008.4884

X^3		φ^6 and $\varphi^4 D^2$		$\psi^2 \varphi^3$	
Q_G	$f^{ABC} G_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	Q_φ	$(\varphi^\dagger \varphi)^3$	$Q_{e\varphi}$	$(\varphi^\dagger \varphi)(\bar{l}_p e_r \varphi)$
$Q_{\tilde{G}}$	$f^{ABC} \tilde{G}_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	$Q_{\varphi\Box}$	$(\varphi^\dagger \varphi)\Box(\varphi^\dagger \varphi)$	$Q_{u\varphi}$	$(\varphi^\dagger \varphi)(\bar{q}_p u_r \tilde{\varphi})$
Q_W	$\varepsilon^{IJK} W_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$	$Q_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^* (\varphi^\dagger D_\mu \varphi)$	$Q_{d\varphi}$	$(\varphi^\dagger \varphi)(\bar{q}_p d_r \varphi)$
$Q_{\tilde{W}}$	$\varepsilon^{IJK} \tilde{W}_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$				
$X^2 \varphi^2$		$\psi^2 X \varphi$		$\psi^2 \varphi^2 D$	
$Q_{\varphi G}$	$\varphi^\dagger \varphi G_{\mu\nu}^A G^{A\mu\nu}$	Q_{eW}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\varphi l}^{(1)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{l}_p \gamma^\mu l_r)$
$Q_{\varphi \tilde{G}}$	$\varphi^\dagger \varphi \tilde{G}_{\mu\nu}^A G^{A\mu\nu}$	Q_{eB}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \varphi B_{\mu\nu}$	$Q_{\varphi l}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi)(\bar{l}_p \tau^I \gamma^\mu l_r)$
$Q_{\varphi W}$	$\varphi^\dagger \varphi W_{\mu\nu}^I W^{I\mu\nu}$	Q_{uG}	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \tilde{\varphi} G_{\mu\nu}^A$	$Q_{\varphi e}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{e}_p \gamma^\mu e_r)$
$Q_{\varphi \tilde{W}}$	$\varphi^\dagger \varphi \tilde{W}_{\mu\nu}^I W^{I\mu\nu}$	Q_{uW}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \tilde{\varphi} W_{\mu\nu}^I$	$Q_{\varphi q}^{(1)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{q}_p \gamma^\mu q_r)$
$Q_{\varphi B}$	$\varphi^\dagger \varphi B_{\mu\nu} B^{\mu\nu}$	Q_{uB}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tilde{\varphi} B_{\mu\nu}$	$Q_{\varphi q}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi)(\bar{q}_p \tau^I \gamma^\mu q_r)$
$Q_{\varphi \tilde{B}}$	$\varphi^\dagger \varphi \tilde{B}_{\mu\nu} B^{\mu\nu}$	Q_{dG}	$(\bar{q}_p \sigma^{\mu\nu} T^A d_r) \varphi G_{\mu\nu}^A$	$Q_{\varphi u}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{u}_p \gamma^\mu u_r)$
$Q_{\varphi WB}$	$\varphi^\dagger \tau^I \varphi W_{\mu\nu}^I B^{\mu\nu}$	Q_{dW}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\varphi d}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{d}_p \gamma^\mu d_r)$
$Q_{\varphi \tilde{W}B}$	$\varphi^\dagger \tau^I \varphi \tilde{W}_{\mu\nu}^I B^{\mu\nu}$	Q_{dB}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \varphi B_{\mu\nu}$	$Q_{\varphi ud}$	$i(\tilde{\varphi}^\dagger D_\mu \varphi)(\bar{u}_p \gamma^\mu d_r)$

d=6: the Warsaw basis

Grzadkowski, Iskrzynski, Misiak, Rosiek 1008.4884

$(\bar{L}L)(\bar{L}L)$		$(\bar{R}R)(\bar{R}R)$		$(\bar{L}L)(\bar{R}R)$	
Q_{ll}	$(\bar{l}_p \gamma_\mu l_r)(\bar{l}_s \gamma^\mu l_t)$	Q_{ee}	$(\bar{e}_p \gamma_\mu e_r)(\bar{e}_s \gamma^\mu e_t)$	Q_{le}	$(\bar{l}_p \gamma_\mu l_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{qq}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{q}_s \gamma^\mu q_t)$	Q_{uu}	$(\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$	Q_{lu}	$(\bar{l}_p \gamma_\mu l_r)(\bar{u}_s \gamma^\mu u_t)$
$Q_{qq}^{(3)}$	$(\bar{q}_p \gamma_\mu \tau^I q_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	Q_{dd}	$(\bar{d}_p \gamma_\mu d_r)(\bar{d}_s \gamma^\mu d_t)$	Q_{ld}	$(\bar{l}_p \gamma_\mu l_r)(\bar{d}_s \gamma^\mu d_t)$
$Q_{lq}^{(1)}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{q}_s \gamma^\mu q_t)$	Q_{eu}	$(\bar{e}_p \gamma_\mu e_r)(\bar{u}_s \gamma^\mu u_t)$	Q_{qe}	$(\bar{q}_p \gamma_\mu q_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{lq}^{(3)}$	$(\bar{l}_p \gamma_\mu \tau^I l_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	Q_{ed}	$(\bar{e}_p \gamma_\mu e_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{u}_s \gamma^\mu u_t)$
		$Q_{ud}^{(1)}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{u}_s \gamma^\mu T^A u_t)$
		$Q_{ud}^{(8)}$	$(\bar{u}_p \gamma_\mu T^A u_r)(\bar{d}_s \gamma^\mu T^A d_t)$	$Q_{qd}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{d}_s \gamma^\mu d_t)$
				$Q_{qd}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{d}_s \gamma^\mu T^A d_t)$
$(\bar{L}R)(\bar{R}L)$ and $(\bar{L}R)(\bar{L}R)$		B -violating			
Q_{ledq}	$(\bar{l}_p^j e_r)(\bar{d}_s^j q_t^i)$	Q_{duq}	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} [(d_p^\alpha)^T C u_r^\beta] [(q_s^j)^T C l_t^k]$		
$Q_{quqd}^{(1)}$	$(\bar{q}_p^j u_r) \varepsilon_{jk} (\bar{q}_s^k d_t)$	Q_{qqqu}	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(u_s^\gamma)^T C e_t]$		
$Q_{quqd}^{(8)}$	$(\bar{q}_p^j T^A u_r) \varepsilon_{jk} (\bar{q}_s^k T^A d_t)$	Q_{qqqq}	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} \varepsilon_{mnp} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(q_s^m)^T C l_t^n]$		
$Q_{lequ}^{(1)}$	$(\bar{l}_p^j e_r) \varepsilon_{jk} (\bar{q}_s^k u_t)$	Q_{duu}	$\varepsilon^{\alpha\beta\gamma} [(d_p^\alpha)^T C u_r^\beta] [(u_s^\gamma)^T C e_t]$		
$Q_{lequ}^{(3)}$	$(\bar{l}_p^j \sigma_{\mu\nu} e_r) \varepsilon_{jk} (\bar{q}_s^k \sigma^{\mu\nu} u_t)$				

SMEFT is the new Fermi theory

Standard **M**odel **E**ffective **F**ield **T**heory:
The EFT constructed with **Standard Model** field & symmetries

→ expansion in **canonical dimensions** d :

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda} \mathcal{L}_5 + \frac{1}{\Lambda^2} \mathcal{L}_6 + \frac{1}{\Lambda^3} \mathcal{L}_7 + \frac{1}{\Lambda^4} \mathcal{L}_8 + \dots$$

$$\mathcal{L}_d = \sum_i C_i \mathcal{O}_i^{(d)}$$

SMEFT describes \sim **any beyond-SM physics living at $\Lambda \gg v$**
(nearly decoupled)



measure C_i to extract information about BSM!

A worthy challenge!

new physics seems indeed
nearly decoupled

ATLAS Exotics Searches* - 95% CL Upper Exclusion Limits

Status: May 2020

ATLAS Preliminary

$$\int \mathcal{L} dt = (32 - 139) \text{ fb}^{-1}$$

$$\sqrt{s} = 8, 13 \text{ TeV}$$

	Model	ℓ, γ	Jets †	E^{miss}	$[\mathcal{L}_{\text{obs}}(\text{fb}^{-1})]$	Limit	Reference
Extra dimensions	ADD $G_{XX} + g/g'$	$0, \mu$	1-4	Yes	36.1	3.7 TeV	$n=2$ 17110301
	ADD nonrenormalized $\gamma\gamma$	$2, \gamma$	-	-	36.7	8.8 TeV	$n=3$ 17030147
	ADD DBH	-	2j	-	37.0	8.8 TeV	$n=6$ 17030127
	ADD BH High Σp_T	$2, \mu, \tau$	2-3j	-	3.2	3.6 TeV	$n=6, M_0 = 3.7 \text{ TeV}$ (a BH) 16050285
	ADD BH multiple	$2, \mu, \tau$	2-3j	-	3.3	3.6 TeV	$n=6, M_0 = 3.7 \text{ TeV}$ (a BH) 15120286
	RS1 $G_{XX} \rightarrow \gamma\gamma$	$2, \gamma$	-	-	36.7	4.1 TeV	$M_{\text{Pl}} = 0.1$ 17070447
	Bulk RS $G_{XX} \rightarrow WW/JZ$	multi-channel	$2j/1, J$	Yes	36.1	2.3 TeV	$M_{\text{Pl}} = 10$ 16030280
	Bulk RS $G_{XX} \rightarrow t\bar{t}$	$1, \mu, \tau$	2 b, 2 J/Z	Yes	36.1	2.0 TeV	$M_{\text{Pl}} = 10$ 20041406
	Bulk RS $G_{XX} \rightarrow t\bar{t}$	$1, \mu, \tau$	2 b, 2 J/Z	Yes	36.1	3.8 TeV	$f/m = 0\%$ 16041683
	ZUFED/APP	$1, \mu, \tau$	2 b, 2 J/Z	Yes	36.1	1.6 TeV	$\text{Tr}[(\gamma, \mu, \tau) \otimes g^{(1)}] - \text{Tr}[\gamma]$ 16030287
Gauge bosons	SSM $Z' \rightarrow \ell\ell$	$2, \mu, \tau$	-	-	1.39	2.1 TeV	$f/m = 12\%$ 19030344
	SSM $Z' \rightarrow \tau\tau$	$2, \mu, \tau$	-	-	1.39	2.4 TeV	17081742
	Leptoophobic $Z' \rightarrow b\bar{b}$	-	2 b	-	36.1	2.1 TeV	16050409
	Leptoophobic $Z' \rightarrow t\bar{t}$	$0, \mu, \tau$	2 b, 2 J/Z	Yes	1.39	2.1 TeV	200505109
	SSM $W' \rightarrow \ell\ell$	$1, \mu, \tau$	-	-	1.39	4.1 TeV	16050503
	SSM $W' \rightarrow \tau\tau$	$1, \tau$	-	-	36.1	3.7 TeV	16010682
	HVT $W' \rightarrow WZ$ (4-ary model B)	$1, \mu, \tau$	2 J/Z, 1 J	Yes	1.39	4.2 TeV	20041406
	HVT $W' \rightarrow WW$ (4-ary model B)	$0, \mu, \tau$	2 J	-	1.39	3.8 TeV	16060898
	HVT $W' \rightarrow WW/ZH$ model B	multi-channel	-	-	36.1	2.93 TeV	17130618
	HVT $W' \rightarrow WH$ model B	$0, \mu, \tau$	2 b, 2 J/Z	Yes	1.39	3.2 TeV	CGP/ATLAS-2019-078
LRSM $W'_\mu \rightarrow \ell\ell$	multi-channel	-	-	36.1	3.5 TeV	160710473	
LRSM $W'_\mu \rightarrow \mu\nu$	$2, \mu$	1 J	-	80	3.8 TeV	190412679	
CI	CI $\eta\eta\eta$	-	2j	-	37.0	2.18 TeV	17030127
	CI $\ell\ell\eta\eta$	$2, \mu, \tau$	-	-	1.39	2.4 TeV	CGP/ATLAS-2019-078
	CI $\ell\ell\tau\tau$	$2, \mu, \tau$	2 b, 2 J/Z	Yes	36.1	2.97 TeV	161102305
DM	Axial-vector mediator (Dirac DM)	$0, \mu, \tau$	1-4j	Yes	36.1	1.55 TeV	$\mu_{\text{eff}} = 0.25, g_{\ell, 0}, \mu_{\text{eff}} = 1$ 171103301
	Vector-scalar mediator (Dirac DM)	$0, \mu, \tau$	1-4j	Yes	36.1	1.67 TeV	$g_{\ell, 0}, \mu_{\text{eff}} = 1$ GeV 171103301
	UV $\chi\chi$ (Dirac DM)	$0, \mu, \tau$	1, 2, 3 j	Yes	3.2	700 GeV	$\mu_{\text{eff}} = 1$ GeV 16030272
	Scalar mediator, $\chi \rightarrow \tau$ (Dirac DM)	$0, 1, \mu, \tau$	1 b, 0, 1, 1 j	Yes	36.1	3.4 TeV	$\gamma \rightarrow b, \ell \rightarrow 0.5, \mu_{\text{eff}} = 10$ GeV 16030283
LO	Scalar LO 1 st gen	$1, 2, \mu, \tau$	2 j	Yes	36.1	1.0 TeV	$\beta = 1$ 16020677
	Scalar LO 2 nd gen	$1, 2, \mu, \tau$	2 j	Yes	36.1	1.36 TeV	$\beta = 1$ 16020677
	Scalar LO 3 rd gen	$2, \mu, \tau$	2 b	Yes	36.1	1.03 TeV	$\beta(\text{SU}(2)_C - \beta) = 1$ 160206103
	Scalar LO 4 th gen	$0, 1, \mu, \tau$	2 b	Yes	36.1	9.73 GeV	$\beta(\text{SU}(2)_C - \beta) = 0$ 160206103
Heavy quarks	$WLO \text{ TT} \rightarrow H/Z/Wb + X$	multi-channel	-	-	36.1	1.37 TeV	BSU2 doublet 160602343
	$WLO \text{ BB} \rightarrow WJ/Zb + X$	multi-channel	-	-	36.1	1.24 TeV	BSU2 doublet 160602343
	$WLO \text{ T} \rightarrow \tau\tau/\tau\tau + W + X$	$2/5/5/2, \mu, \tau$	2 b, 2 J/Z	Yes	36.1	1.64 TeV	DY production, $\beta(\text{SU}(2)_C - \beta) = 1$ 160711883
	$WLO \text{ Y} \rightarrow Wb + X$	$1, \mu, \tau$	1 b, 1, 2 j	Yes	36.1	1.65 TeV	$\beta(\text{Y} \rightarrow Wb) = 1, \alpha(\text{SU}(2)_C) = 1$ 160711883
	$WLO \text{ Q} \rightarrow Hb + X$	$0, \mu, \tau$	2 b, 2 J/Z	Yes	79.8	1.21 TeV	$\alpha_{\text{em}} = 0.5$ ATLAS CONF-2018-024
$WLO \text{ QQ} \rightarrow WqWq$	$1, \mu, \tau$	2 j	Yes	20.3	600 GeV	15080481	
Excited fermions	Excited quark $q^* \rightarrow q\gamma$	-	2j	-	1.39	6.7 TeV	only u^* and d^* , $A = m(q^*)$ 191004467
	Excited quark $q^* \rightarrow q\tau$	$1, \tau$	1 j	-	26.7	6.3 TeV	$\beta(\text{SU}(2)_C - \beta) = 1, A = m(q^*)$ 170811463
	Excited quark $q^* \rightarrow b\gamma$	-	1 b, 1 j	-	36.1	2.6 TeV	16050409
	Excited lepton $\ell^* \rightarrow \ell\gamma$	$1, \mu, \tau$	-	-	20.3	3.7 TeV	$A = 3.0 \text{ TeV}$ 1411281
	Excited lepton $\ell^* \rightarrow \ell\tau$	$1, \mu, \tau$	-	-	20.3	3.6 TeV	$A = 1.6 \text{ TeV}$ 1411281
Other	Type III Seesaw	$1, \mu, \tau$	2 j	Yes	79.8	360 GeV	ATLAS CONF-2018-020
	LRSM Majorana ν	$2, \mu$	2 j	-	36.1	3.2 TeV	$m(W_2) = 4.1 \text{ TeV}, g_{\ell} = g_{\nu}$ 16091105
	Higgs triplet $H^{\pm\pm} \rightarrow \ell\ell$	$2, 3, 4, \mu, \tau$ (BS)	-	-	36.1	870 GeV	DY production 171003743
	Higgs triplet $H^{\pm\pm} \rightarrow \tau\tau$	$2, 3, \mu, \tau$	-	-	20.3	400 GeV	DY production, $\beta(\text{SU}(2)_C - \beta) = 1$ 1411281
	Multi-charged particles	-	-	-	36.1	1.22 TeV	DY production, $g_{\ell} = g_{\nu}$ 161203673
Magnetic monopoles	-	-	-	34.4	2.37 TeV	DY production, $g_{\ell} = 1 \text{ g.p.}, \alpha_{\text{em}} = 1/3$ 190510190	

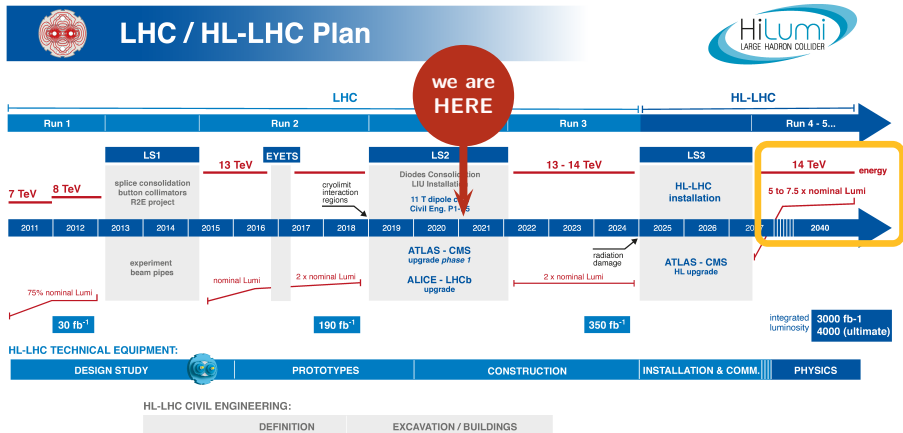
*Only a selection of the available mass limits on new states or phenomena is shown.

† Small-radius (large-radius) jets are denoted by the letter j (J).

A worthy challenge!

new physics seems indeed nearly decoupled

collider physics is entering a precision era



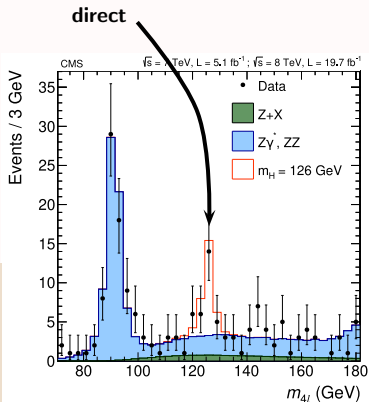
A worthy challenge!

new physics seems indeed nearly decoupled

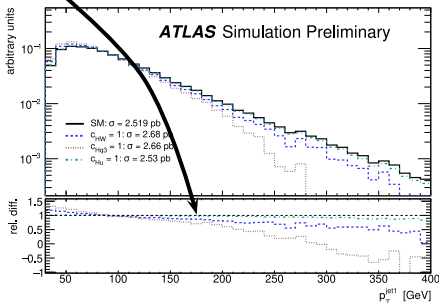
collider physics is entering a precision era



indirect searches
more and more competitive
with direct ones



indirect



CMS-HIG-13-002

ATL-PHYS-PUB-2019-042

A worthy challenge!

new physics seems indeed nearly decoupled



SMEFT-based searches at the LHC are crucial

collider physics is entering a precision era



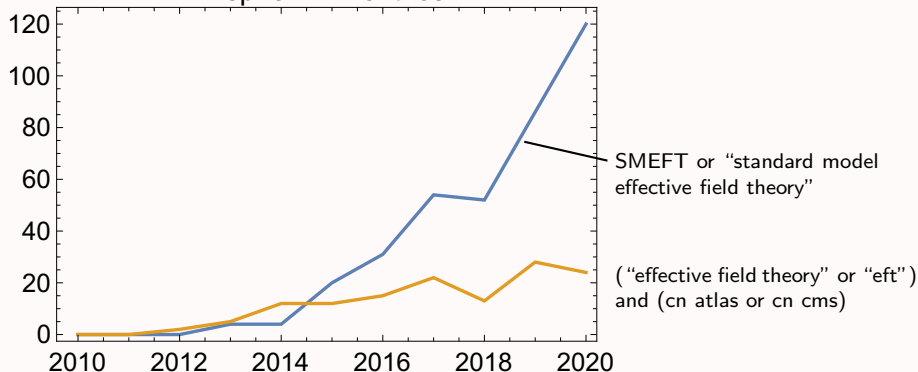
indirect searches
more and more competitive
with direct ones



- + a proper **QFT** :
renormalizable order by order, well-defined radiative corrections and RGE
- + minimal commitment to a specific UV
- + systematically includes **all** BSM effects, compatible with assumptions
- + **universal language** for data interpretation: can connect to other experiments

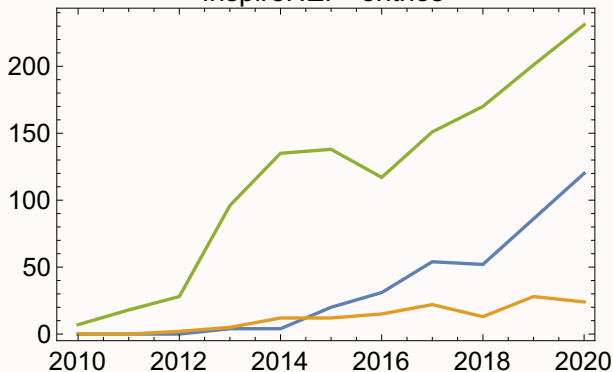
A booming field

InspireHEP entries



A booming field

InspireHEP entries



refersto:recid:866649

Grzadkowski,Iskrzynski,Misiak,Rosiek
1008.4884

SMEFT or "standard model
effective field theory"

("effective field theory" or "eft")
and (cn atlas or cn cms)

+ dedicated seminar series, conferences, working groups. . .





✓ # parameters known for all orders

Lehman(Martin) 1410.4193,1510.00372
Henning,Lu,Melia,Murayama 1512.03433

calculation



measurement



global analysis



interpretation

✓ # parameters known for all orders

✓ complete bases up to $d = 9$

5: Weinberg PRL43(1979)1566

6: Buchmller,Wyler Nucl.Phys.B268(1986)621, Grzadkowski et al 1008.4884

7: Lehman 1410.4193, Henning,Lu,Melia,Murayama 1512.0343

8: Li, Ren, Shu, Xiao, Yu, Zheng 2005.00008, Murphy 2005.00059

9: Li, Ren, Xiao, Yu, Zheng 2007.07899, Liao, Ma 2007.08125

calculation

measurement

global analysis

interpretation

- ✓ # parameters known for all orders
- ✓ complete bases up to $d = 9$
- ✓ interplay with helicity amplitudes

Shadmi, Weiss 1809.09644
Henning, Melia 1901.06747, 1902.06754, 1902.06747
Ma, Shu, Xiao 1902.06752
Aoude, Machado 1905.11433
Durieux, Kitahara, (Machado), Shadmi, Weiss
1909.10551, 2008.09652
Durieux, Machado 1912.08827
Craig, Jiang, Li, Sutherland 2001.00017

calculation

measurement

global analysis

interpretation

- ✓ # parameters known for all orders
- ✓ complete bases up to $d = 9$
- ✓ interplay with helicity amplitudes
- ✓ geometric formulation

Alonso, Jenkins, Manohar 1511.00724, 1605.03602
Dedes, Materkowska, Paraskevas, Rosiek, Suxho 1704.03888
Helset, Paraskevas, Trott 1803.08001
Corbett, Helset, Trott 1909.08470
(Hays), Helset, Martin, Trott 2001.01453, 2007.00565

calculation

measurement

global analysis

interpretation

✓ # parameters known for all orders

✓ complete bases up to $d = 9$

✓ interplay with helicity amplitudes

✓ geometric formulation

✓ RGE up to $d = 6$

Alonso, Jenkins, Manohar, Trott 1308.2627, 1310.4838, 1312.2014

Grojean, Jenkins, Manohar, Trott 1301.2588

Alonso, Chang, Jenkins, Manohar, Shotwell 1405.0486

Ghezzi, Gomez-Ambrosio, Passarino, Uccirati 1505.03706

Miro, Ingoldby, Riemann 2005.06983

Baratella, Fernandez, Pomarol 2005.07129

calculation

measurement

global analysis

interpretation

✓ # parameters known for all orders

✓ complete bases up to $d = 9$

✓ interplay with helicity amplitudes

✓ geometric formulation

✓ RGE up to $d = 6$

✓ flavor structure

Brivio, Jiang, Trott 1709.06492

Bordone, Catà, Feldmann 1910.02641

Faroughy, Isidori, Wilsch, Yamamoto 2005.05366

Brivio 2012.11343

calculation

measurement

global analysis

interpretation

- ✓ # parameters known for all orders
- ✓ complete bases up to $d = 9$
- ✓ interplay with helicity amplitudes
- ✓ geometric formulation
- ✓ RGE up to $d = 6$
- ✓ flavor structure
- ✓ NLO SMEFT

Pruna,Signer 1408.3565
Hartmann,(Shepherd),Trott 1505.02646,1507.03568,1611.09879
Ghezzi,Gomez-Ambrosio,Passarino,Uccirati 1505.03706
(Cullen,Gauld),Pecjak,Scott 1512.02508,1904.06358,2007.15238
Dawson,Giardino 1801.01136,1807.11504,1808.05948,1909.02000
Deutschmann,Duhr,Maltoni,Vryonidou 1708.00460
Grazzini,Ilnicka,Spira 1806.08832
Boughezal,Chen,Petriello,Wiegand 1907.00997
Dedes,Paraskevas,Rosiek,Suxho,Trifyllis 1805.00302
+ several NLO QCD

calculation

measurement

global analysis

interpretation

- ✓ # parameters known for all orders
- ✓ complete bases up to $d = 9$
- ✓ interplay with helicity amplitudes
- ✓ geometric formulation
- ✓ RGE up to $d = 6$
- ✓ flavor structure
- ✓ NLO SMEFT
- ✓ connection to WET/LEFT

Jenkins, Manohar, Stoffer 1709.04486
Dekens, Stoffer 1908.05295

calculation

measurement

global analysis

interpretation

- ☑ automation
 - ▶ basis handling
 - Falkowski et al 1508.05895 Aebischer et al 1712.05298 Criado 1901.03501
 - ▶ matching & running
 - (Celis),Fuentes-Martin,(Ruiz-Femenia),Vicente,Virto 1704.04504,2010.16341 Aebischer,Kumar,Straub 1804.05033
 - ▶ Feynman rules
 - Dedes,(Materkowska),Paraskevas,Rosiek,Suxho,(Trifyllis) 1704.03888,1904.03204, Brivio,Jiang,Trott 1907.04692 Corbett 2010.15852
 - ▶ LO predictions
 - Alloul,Fuks,Sanz 1310.5150 Durieux,Mattelaer 1802.07237 Brivio,(Jiang,Trott) 1907.04692,2012.11343
 - ▶ NLO QCD predictions
 - Degrande,Durieux,Maltoni,Mimasu,Vryonidou 2008.11743

SMEFT predictions in practice: SMEFTsim



a set of models that enable
LO Monte Carlo simulations
with the complete Warsaw basis

Brivio, Jiang, Trott 1709.06492
Brivio 2012.11343
🔗 [SMEFTsim.github.io](https://github.com/SMEFTsim)

Automates all the steps required:

1. bosonic: make kinetic terms canonical, define vev. properly
fermionic: fix flavor structure
2. choose input parameters
3. Feynman diagram calculation (MC). subtlety: propagator corrections

Example: $h \rightarrow e^+ e^- \mu^+ \mu^-$ at tree-level, $\mathcal{O}(\Lambda^{-2}) = \text{SM-}\mathcal{L}_6$ interference

analytic result derived in [Brivio, Corbett, Trott 1906.06949](#)

w/o flavor assumptions \mathcal{L}_6 has **2499** free parameters

$$\left\| \begin{array}{ll} O_{He,pr} = (H\overleftrightarrow{D}_\mu H)(\bar{e}_p\gamma^\mu e_r) & \text{has } \mathbf{9} \text{ independent par.} \\ O_{ledq,prst} = (\bar{l}_p^i e_r)(\bar{d}_s q_t^i) & \text{has } \mathbf{162} \end{array} \right.$$

freedom can be reduced imposing a **symmetry**. Maximal:

$$U(3)^5 \equiv U(3)_q \times U(3)_u \times U(3)_d \times U(3)_l \times U(3)_e$$

→ only invariant contractions allowed

→ Yukawa couplings typically promoted to **spurions**:

$$Y_u \mapsto \Omega_u Y_u \Omega_q^\dagger, \quad Y_d \mapsto \Omega_d Y_d \Omega_q^\dagger, \quad Y_e \mapsto \Omega_e Y_e \Omega_l^\dagger$$

$$\left\| \begin{array}{ll} O_{He,pr} \delta_{pr} & \text{has } \mathbf{1} \text{ independent par.} \\ O_{ledq,prst} (Y_e^\dagger)_{pr} (Y_d)_{st} & \text{has } \mathbf{2} \end{array} \right.$$

$\mathcal{L}_6 + U(3)^5$ has **85**
free parameters

several options implemented:

general	free indices	2499
U35	$U(3)_q \times U(3)_u \times U(3)_d \times U(3)_l \times U(3)_e$	85
MFV	$U(3)_q \times U(3)_u \times U(3)_d \times U(3)_l \times U(3)_e$ +leading spurion insertions $\propto Y_u Y_u^\dagger, Y_d Y_d^\dagger, \dots$	120
topU31	$U(2)_q \times U(2)_u \times U(2)_d \times U(3)_l \times U(3)_e$	182
top	$U(2)_q \times U(2)_u \times U(2)_d \times U(1)_{l+e}^3$	275

SM

input obs.

\mathcal{O}_n

$\alpha_s, d_N,$
 $m_Z, \alpha_{\text{em}},$
 $G_F, m_h,$
 $m_f,$
meson decay/osc

\mathcal{L} param.

g_i

$g_s, \theta_{\text{QCD}},$
 $g_1, g_w,$
 $v, \lambda,$
 $y_f,$
 $\theta_{1,2,3}, \delta$ CKM

predicted obs.

\mathcal{P}

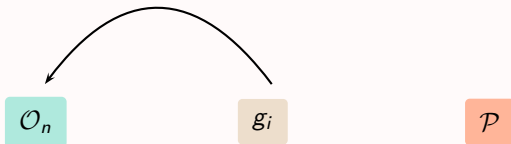
$m_W, \Gamma_h, \Gamma_Z,$
 $\sigma(pp \rightarrow X)$
...

Input parameters

SM

$$G_F = \frac{1}{\sqrt{2}v^2}$$

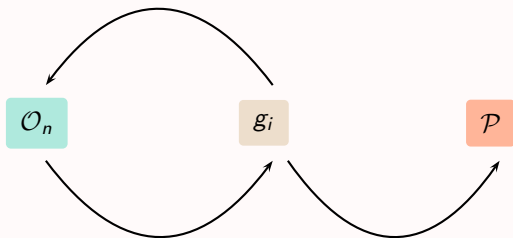
$$\mathcal{O}_n = F_n^{SM}(g)$$



SM

$$G_F = \frac{1}{\sqrt{2}v^2}$$

$$\mathcal{O}_n = F_n^{SM}(g)$$



$$g_i = K_i^{SM}(\mathcal{O})$$

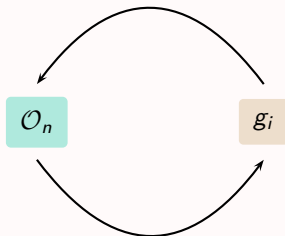
$$\mathcal{P} = P^{SM}(g) = P^{SM}(K^{SM}(\mathcal{O}))$$

$$v = \frac{1}{2^{1/4}\sqrt{G_F}}$$

SMEFT

$$G_F = \frac{1}{\sqrt{2}v^2} + \frac{1}{\sqrt{2}\Lambda^2} (2C_{HI}^{(3)} - C'_{II}) + \dots$$

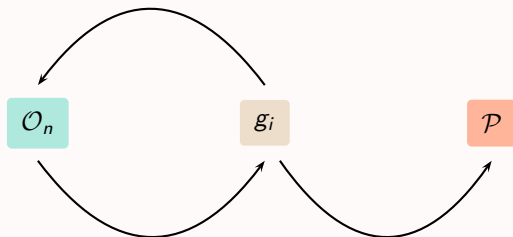
$$\mathcal{O}_n = F_n^{SM}(\mathbf{g}) + \frac{1}{\Lambda^2} F_n^{(2)}(\mathbf{g}, C) + \dots$$



$$g_i = K_i^{SM}(\mathcal{O}) + \frac{1}{\Lambda^2} K_i^{(2)}(\mathcal{O}, C) + \dots, \quad K_i^{(2)} = - \left[\frac{\partial F^{SM}}{\partial \mathbf{g}} \right]_{in}^{-1} F_n^{(2)}$$

$$v = \frac{1}{2^{1/4} \sqrt{G_F}} + \frac{1}{2^{3/4} G_F^{3/2}} \frac{2C_{HI}^{(3)} - C'_{II}}{2\Lambda^2} + \dots$$

SMEFT



$$\mathcal{P} = P^{SM}(g) + \frac{1}{\Lambda^2} \left[\underbrace{P^{(2)}(g)}_{\text{direct corr.}} - \frac{\partial P^{SM}}{\partial g_i} \left[\frac{\partial F_n^{SM}}{\partial g_i} \right]^{-1} F_n^{(2)} \right] + \dots$$

direct corr.

input shift
contribution

EW input schemes

Two most used options for the EW sector:

$$\begin{aligned} \{\alpha_{\text{em}}, m_Z, G_F\} &\rightarrow \Delta\alpha_{\text{em}} \equiv 0 \\ &\rightarrow \Delta m_W^2 \neq 0 \quad (\sim C_{HD}, C_{HWB}, C_{HI}^{(3)}, C'_{II}) \end{aligned}$$

$$\begin{aligned} \{m_W, m_Z, G_F\} &\rightarrow \Delta\alpha_{\text{em}} \neq 0 \quad (\sim C_{HD}, C_{HWB}, C_{HI}^{(3)}, C'_{II}) \\ &\rightarrow \Delta m_W^2 \equiv 0 \end{aligned}$$

EW input schemes

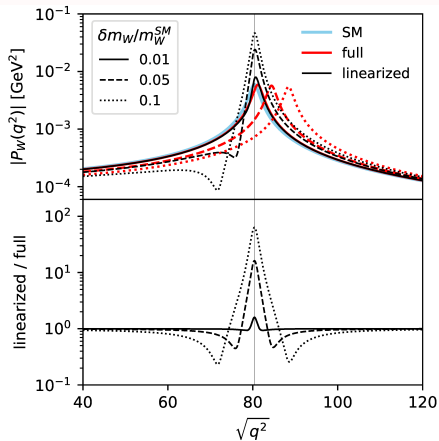
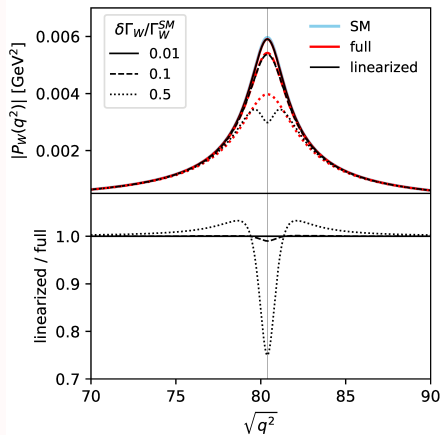
Two most used options for the EW sector:

$$\begin{aligned} \{\alpha_{\text{em}}, m_Z, G_F\} &\rightarrow \Delta\alpha_{\text{em}} \equiv 0 \\ &\rightarrow \Delta m_W^2 \neq 0 \quad (\sim C_{HD}, C_{HWB}, C_{HI}^{(3)}, C'_{II}) \end{aligned}$$

$$\begin{aligned} \{m_W, m_Z, G_F\} &\rightarrow \Delta\alpha_{\text{em}} \neq 0 \quad (\sim C_{HD}, C_{HWB}, C_{HI}^{(3)}, C'_{II}) \\ &\rightarrow \Delta m_W^2 \equiv 0 \end{aligned}$$

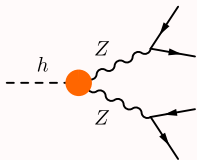
Propagator corrections

$$\frac{i(-\eta^{\mu\nu} + q^\mu q^\nu/m_W^2)}{p^2 - m_W^2 + i\Gamma_W m_W} \left[1 + \frac{im_W \Delta\Gamma_W}{p^2 - m_W^2 + i\Gamma_W m_W} - \frac{(2m_W - i\Gamma_W) \Delta m_W}{p^2 - m_W^2 + i\Gamma_W m_W} \right] + \mathcal{O}(\Lambda^{-4})$$



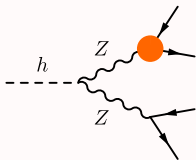
Example: $H \rightarrow 4f$ in the SMEFT

① corrections to SM diagrams

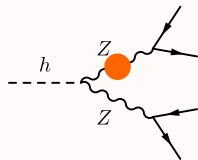


$$\propto g_{\mu\nu} \text{ (SM-like)}$$

$$\propto g_{\mu\nu} p \cdot q - p_\nu q_\mu \text{ (} Z_{\mu\nu} Z^{\mu\nu} h \text{)}$$

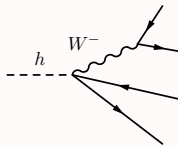
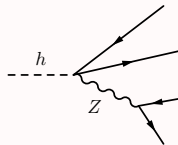
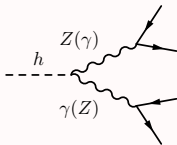
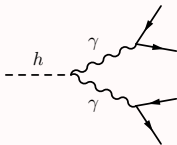


$$\delta g_L, \delta g_R$$



$$\frac{-im_Z \delta \Gamma_Z + (2m_Z - i\Gamma_Z) \delta m_Z}{p^2 - m_Z^2 + i\Gamma_Z m_Z}$$

② genuine SMEFT diagrams



Example: $h \rightarrow 4f$ in the SMEFT

Results with $U(3)^5$ symmetry, $\{m_W, m_Z, G_F\}$ inputs.

$$\Gamma = \Gamma_{SM} \left[1 + \sum_{\alpha} \bar{C}_{\alpha} N_{\alpha} \right], \quad \bar{C}_{\alpha} = C_{\alpha} \frac{v^2}{\Lambda^2}$$



	$\bar{q}q \rightarrow h\bar{q}q$ VBF-like		$h \rightarrow e^+e^-\mu^+\mu^-$	
	direct	propagators	direct	propagators
\bar{C}_{He}		$5.32 \cdot 10^{-5}$	-1.724	0.153
$\bar{C}_{Hl}^{(1)}$		$5.32 \cdot 10^{-5}$	2.144	0.153
$\bar{C}_{Hl}^{(3)}$	-6	$1.351 \cdot 10^{-3}$	-3.856	1.147
$\bar{C}_{Hq}^{(1)}$	0.109	$-1.363 \cdot 10^{-4}$		-0.39
$\bar{C}_{Hq}^{(3)}$	-5.345	$-1.423 \cdot 10^{-3}$		-1.353
\bar{C}_{Hu}	-0.323	$-7.092 \cdot 10^{-5}$		-0.203
\bar{C}_{Hd}	0.103	$5.24 \cdot 10^{-5}$		0.150
\bar{C}'_{ll}	3	$-1. \cdot 10^{-3}$	3	-0.839

calculation

measurement

global analysis

interpretation

✓ Higgs + EW

Alves et al 1211.4580,1805.11108, Butter et al 1604.03105
Corbett et al 1509.01585, de Blas et al 1608.01509,1710.05402,1910.14012
Ellis,Murphy,Sanz,You 1803.03252 da Silva Almeida et al 1812.01009
Biekötter,Corbett,Plehn 1812.07587, Dawson,Homiller,Lane 2007.01296

✓ top

Englert et al 1506.08845,1512.05560,1901.03164
Cirigliano,Dekens,deVries,Mereghetti 1605.04311 Hartland et al 1901.05965
Brivio et al 1910.0306, Durieux,Irles,Miralles,Peñuelas,Pöschl 1907.10619

✓ Higgs + EW + top

Ellis,Madigan,Mimasu,Sanz,You 2012.02779

✓ top + B physics

Bißmann,(Erdmann),Grunwald,Hiller,Kroöninger 1909.13632, 2012.10456
Bruggisser,Schäfer,Westhoff, VanDyk 2101.07273

✓ diboson (+ VBS)

Baglio,Dawson,Homiller,(Lane,Lewis) 1812.00214, 1909.11576, 2003.07862
Ethier,Gomez-Ambrosio,Magni,Rojo 2101.03180

✓ Higgs combination within ATLAS

2004.03447, ATLAS-CONF-2020-053

✓ impact analysis of NLO corrections and quadratic SMEFT terms

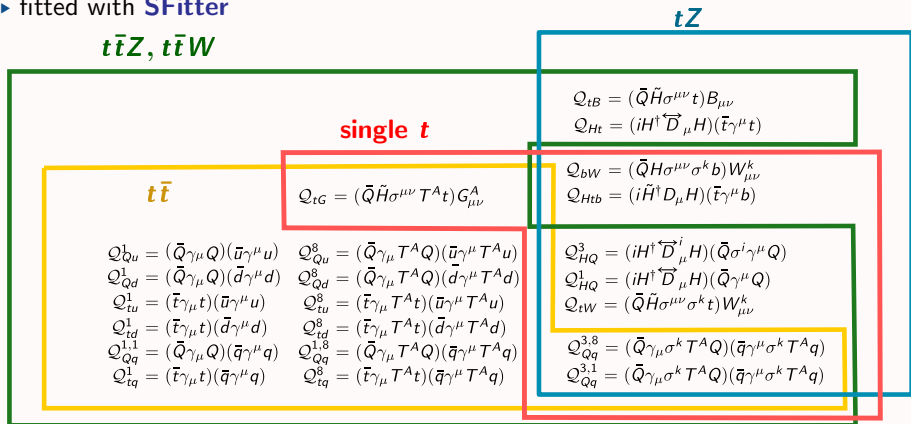
✓ several statistical methods: information geometry, bayesian reweighting, replica model, Partial Component Analysis. . .

Global analysis of top quark observables

- ▶ $U(2)_q \times U(2)_u \times U(2)_d$
- ▶ top interactions only → 24 param.
- ▶ up to NLO QCD, quadratic SMEFT (predictions with SMEFT@NLO)
- ▶ fitted with **SFitter**

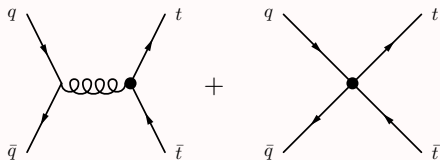
Brivio, Bruggisser, Maltoni, Moutafis, Plehn, Vryonidou, Westhoff, Zhang 1910.03606

also: SMEFIT. Hartland, Maltoni, Nocera, Rojo, Slade, Vryonidou, Zhang 1901.05965
TopFitter 1506.08845, 1512.03360



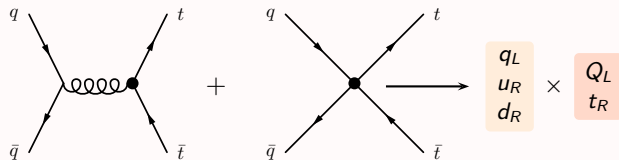
A typical issue: flat directions

e.g. $q\bar{q} \rightarrow t\bar{t}$ at tree-level, linear order in C_i :



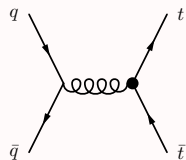
A typical issue: flat directions

e.g. $q\bar{q} \rightarrow t\bar{t}$ at tree-level, linear order in C_i :

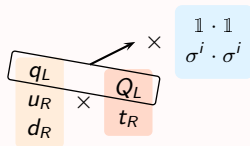
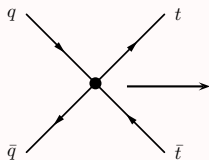


A typical issue: flat directions

e.g. $q\bar{q} \rightarrow t\bar{t}$ at tree-level, linear order in C_i :

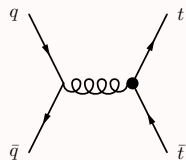


+

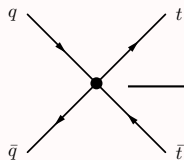


A typical issue: flat directions

e.g. $q\bar{q} \rightarrow t\bar{t}$ at tree-level, linear order in C_i :



+



$$\begin{aligned}
 & \begin{array}{|c|} \hline q_L \\ \hline u_R \\ \hline d_R \\ \hline \end{array} \times \begin{array}{|c|} \hline Q_L \\ \hline t_R \\ \hline \end{array} \times \begin{array}{|c|} \hline 1 \cdot 1 \\ \hline \sigma^i \cdot \sigma^i \\ \hline \end{array} \times \begin{array}{|c|} \hline 1 \cdot 1 \\ \hline T^A \cdot T^A \\ \hline \end{array} \\
 & = 14 \text{ operators}
 \end{aligned}$$

interfering with SM:

7 4-fermion operators

$$C_{LL}^8 \sim C_{Qq}^{(1,8)}, C_{Qq}^{(3,8)}$$

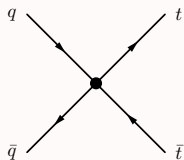
$$C_{LR}^8 \sim C_{tq}^8$$

$$C_{RL}^8 \sim C_{Qu}^8, C_{Qd}^8$$

$$C_{RR}^8 \sim C_{tu}^8, C_{td}^8$$

A typical issue: flat directions

e.g. $q\bar{q} \rightarrow t\bar{t}$ at tree-level, linear order in C_i :



$$\beta_t^2 = 1 - 4m_t^2/s$$
$$c_t = \cos\theta(\vec{p}_t, \vec{p}_q) \text{ in c.m. frame}$$

$$\Delta\sigma_{t\bar{t}}^{int} \propto \left[C_{LL}^8 + C_{RR}^8 + C_{LR}^8 + C_{RL}^8 \right] \left(1 + \beta_t^2 c_t^2 + \frac{2m_t^2}{s} \right) + \left[C_{LL}^8 + C_{RR}^8 - C_{LR}^8 - C_{RL}^8 \right] 2\beta_t c_t$$

breaking down each pairs requires
→ NLO QCD
→ $(C_i C_j)$ terms
→ other processes (eg. single top)

the two pairs can be distinguished
already at this order, with
kinematics

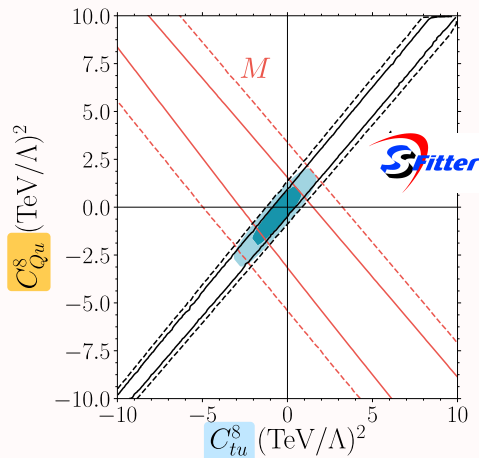
Same vs. different chiralities in $t\bar{t}$

$$\Delta\sigma_{t\bar{t}}^{int} \propto \left[C_{LL}^8 + C_{RR}^8 + C_{LR}^8 + C_{RL}^8 \right] \left(1 + \beta_t^2 c_t^2 + \frac{2m_t^2}{s} \right) + \left[C_{LL}^8 + C_{RR}^8 - C_{LR}^8 - C_{RL}^8 \right] 2\beta_t c_t$$

likelihood contours:

$$\ln L_{\max} - \ln L = \begin{array}{ll} 1/2 & \text{——} \\ 2 & \text{----} \end{array}$$

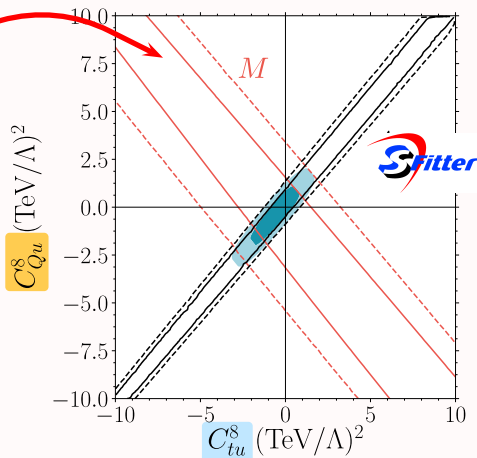
($\sim \Delta\chi^2 = 1, 4$ in Gaussian limit)



Same vs. different chiralities in $t\bar{t}$

$$\Delta\sigma_{t\bar{t}}^{int} \propto \left[C_{LL}^8 + C_{RR}^8 + C_{LR}^8 + C_{RL}^8 \right] \left[1 + \beta_t^2 c_t^2 + \frac{2m_t^2}{s} \right] + \left[C_{LL}^8 + C_{RR}^8 - C_{LR}^8 - C_{RL}^8 \right] 2\beta_t c_t$$

$\sigma_{t\bar{t}} + m_{t\bar{t}}$ dist



likelihood contours:

$$\ln L_{\max} - \ln L = \begin{array}{ll} 1/2 & \text{——} \\ 1 & \text{-----} \\ 2 & \text{-----} \end{array}$$

($\sim \Delta\chi^2 = 1, 4$ in Gaussian limit)

Same vs. different chiralities in $t\bar{t}$

$$\Delta\sigma_{t\bar{t}}^{int} \propto \left[C_{LL}^8 + C_{RR}^8 + C_{LR}^8 + C_{RL}^8 \right] \left[1 + \beta_t^2 c_t^2 + \frac{2m_t^2}{s} \right] + \left[C_{LL}^8 + C_{RR}^8 - C_{LR}^8 - C_{RL}^8 \right] 2\beta_t c_t$$

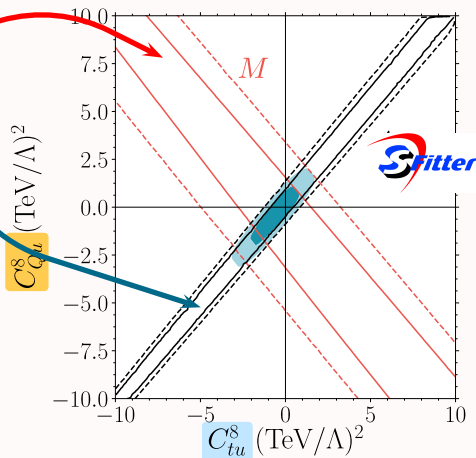
$\sigma_{t\bar{t}} + m_{t\bar{t}}$ dist

charge asymmetry
 A_C

likelihood contours:

$$\ln L_{\max} - \ln L = \begin{array}{ll} 1/2 & \text{—} \\ 2 & \text{- - - -} \end{array}$$

($\sim \Delta\chi^2 = 1, 4$ in Gaussian limit)



Impact of quadratic SMEFT contributions

$$\Delta\sigma_{t\bar{t}}^{quad} \propto \left[(C_{LL}^8)^2 + (C_{RR}^8)^2 + (C_{LR}^8)^2 + (C_{RL}^8)^2 + \frac{9}{2} \left((C_{LL}^1)^2 + (C_{RR}^1)^2 + (C_{LR}^1)^2 + (C_{RL}^1)^2 \right) \right] (1 + \beta_t^2 c_t^2)$$

$$+ \left[(C_{LL}^8)^2 + (C_{RR}^8)^2 - (C_{LR}^8)^2 - (C_{RL}^8)^2 + \frac{9}{2} \left((C_{LL}^1)^2 + (C_{RR}^1)^2 - (C_{LR}^1)^2 - (C_{RL}^1)^2 \right) \right] 2\beta_t c_t$$

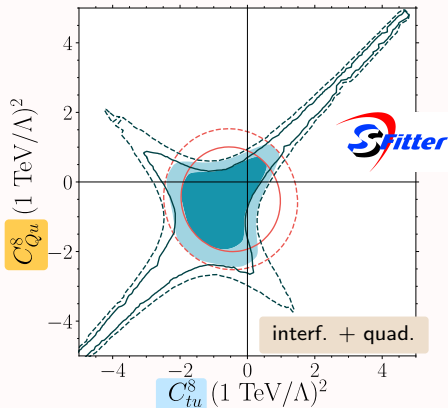
$$+ \left[C_{LL}^8 C_{LR}^8 + C_{RR}^8 C_{RL}^8 + \frac{9}{2} \left(C_{LL}^1 C_{LR}^1 + C_{RR}^1 C_{RL}^1 \right) \right] 4m_t^2/s$$

Quadratic terms modify the **geometry** of the fit

likelihood contours:

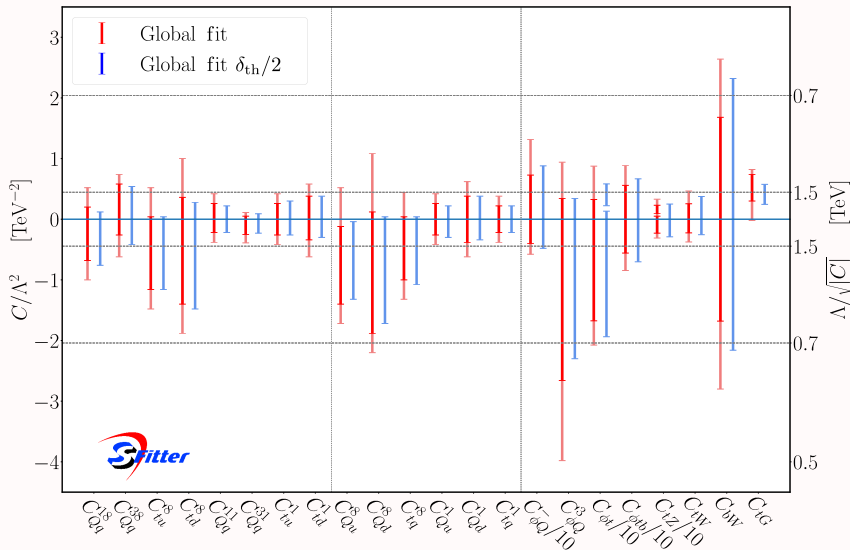
$$\ln L_{\max} - \ln L = \frac{1/2}{2}$$

($\sim \Delta\chi^2 = 1, 4$ in Gaussian limit)



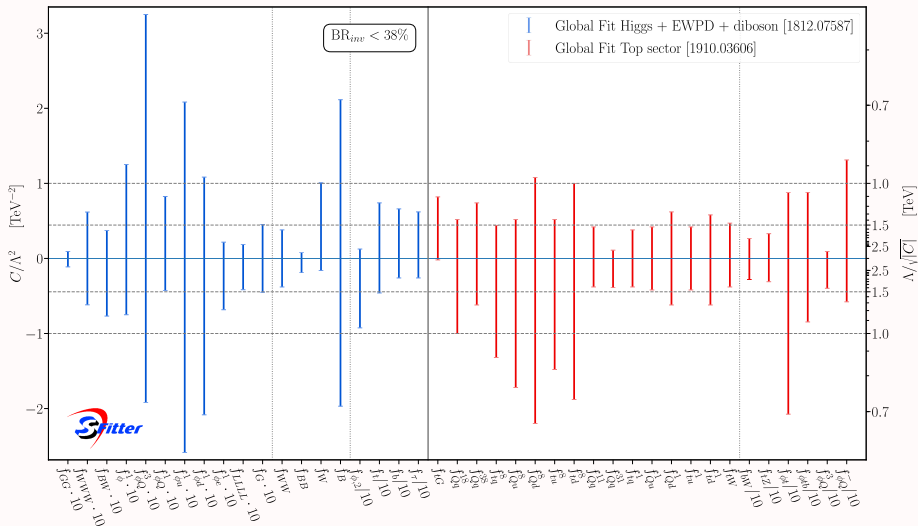
Global analysis of top quark observables

Run II, ATLAS+CMS, 68% and 95% C.L.

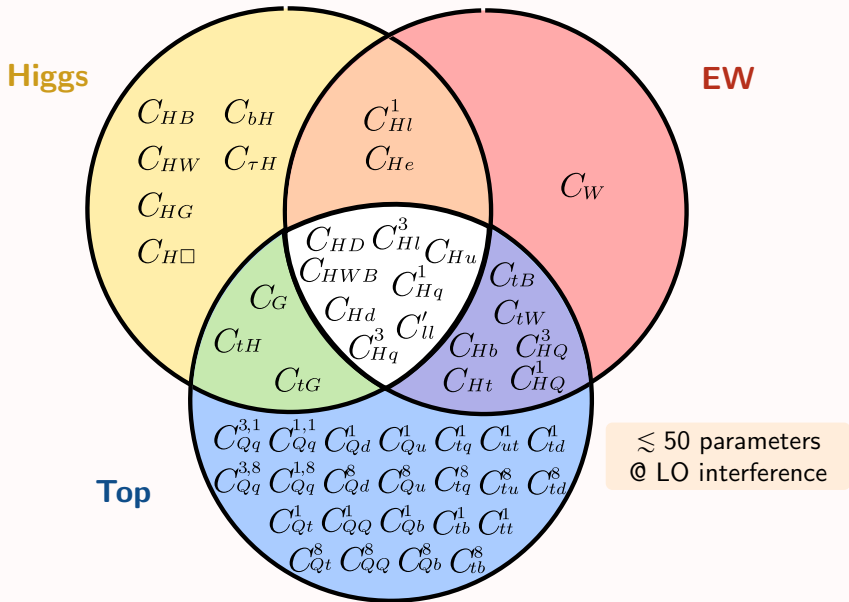


Global analysis of top quark observables

EWPD + LHC Run I + II, 95% C.L.

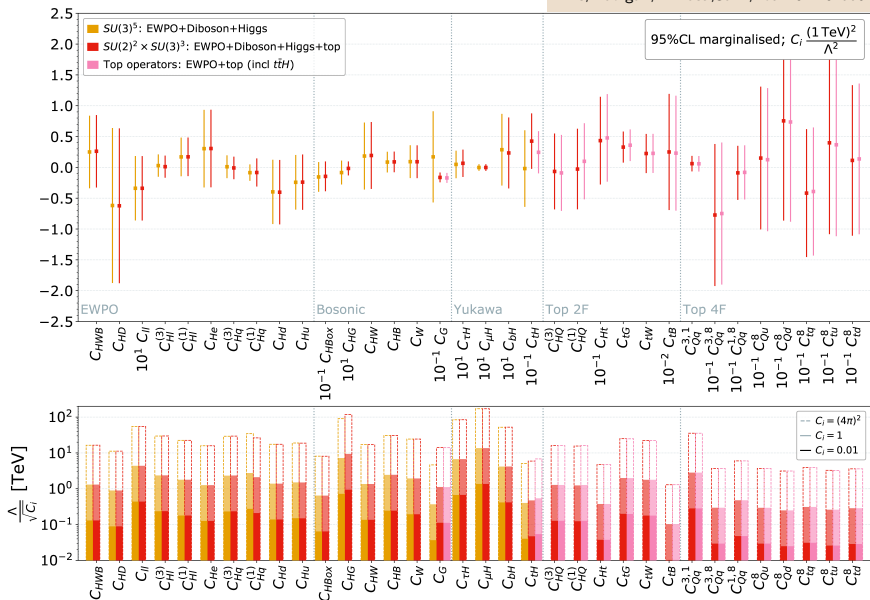


Next step: top + EW + Higgs



Next step: top + EW + Higgs

Ellis, Madigan, Mimasu, Sanz, You 2012.02779





calculation

measurement

global analysis

interpretation

▶ CDE / UOLEA: up to 1-loop

Henning, Lu, Murayama 1412.1837, 1604.01019
del Aguila, Kunszt, Santiago 1602.00126
Drozd, Ellis, Quevillon, You 1512.03003
Boggia, Gomez-Ambrosio, Passarino 1603.03660
Ellis, Quevillon, (Vuong), You, Zhang
1604.02445, 1706.07765, 2006.16260
Fuentes-Martin, Portoles, Ruiz-Femenia 1607.02142
Zhang 1610.00710, Cohen, Lu, Zhang 2011.02484
(Krämer), Summ, Voigt 1806.05171, 1908.04798

▶ matching via amplitudes up to 1-loop

Craig, Jiang, Li, Sutherland 2001.00017

▶ complete tree-level dictionary

de Blas, Criado, Pérez-Victoria, Santiago 1711.10391

▶ “v-improved” matching

(Brehmer), Freitas, López-Val, Plehn 1510.03443, 1607.08251

▶ automated matching

Criado 1710.06445, Bashki, Chakraborty, Kumar Patra 1808.04403
Cohen, Lu, Zhang 2012.07851
Fuentes-Martin, König, Pagès, Thomsen, Wilsch 2012.08506

▶ reduced fits

Gorbahn, No, Sanz 1502.07352, Drozd, Ellis, Quevillon, You 1504.02409
Ellis, (Madigan, Mimasu, Murphy), Sanz, You 1803.03252, 2012.02779
Dawson, Homiller, Lane 2007.01296, 2102.02823
Bakshi, Chakraborty, Englert, Spannowsky, Stylianou 2009.13394
Anisha, Bakshi, Chakraborty, Kumar Patra 2010.04088

Example: SM + vector triplet

Brivio, Bruggisser, Geoffray, Krämer, Plehn, Summ in preparation

$$\begin{aligned}\mathcal{L}_V = & -\frac{1}{4} V_{\mu\nu}^i V^{i\mu\nu} - \frac{g_M}{2} V_{\mu\nu}^i W^{i\mu\nu} + \frac{m_V^2}{2} V_\mu^i V^{i\mu} + \frac{g_H}{2} V_\mu^i (H^\dagger i \overleftrightarrow{D}^{i\mu} H) \\ & + \frac{g_l}{2} V_\mu^- \bar{\ell} \gamma^\mu \sigma^i \ell + \frac{g_q}{2} V_\mu^- \bar{q} \gamma^\mu \sigma^i q + \frac{g_{VH}}{2} (H^\dagger H) V_\mu^i V^{i\mu}\end{aligned}$$

Example: SM + vector triplet

Brivio, Bruggisser, Geoffroy, Krämer, Plehn, Summ in preparation

$$\mathcal{L}_V = -\frac{1}{4} V_{\mu\nu}^i V^{i\mu\nu} - \frac{g_M}{2} V_{\mu\nu}^i W^{i\mu\nu} + \frac{m_V^2}{2} V_\mu^i V^{i\mu} + \frac{g_H}{2} V_\mu^i (H^\dagger i \overleftrightarrow{D}^{i\mu} H) \\ + \frac{g_l}{2} V_\mu^- \bar{\ell} \gamma^\mu \sigma^i \ell + \frac{g_q}{2} V_\mu^- \bar{q} \gamma^\mu \sigma^i q + \frac{g_{VH}}{2} (H^\dagger H) V_\mu^i V^{i\mu}$$

field redefinition to
remove kinetic mixing

$$\left\{ \begin{array}{l} V_\mu^i \rightarrow \frac{1}{\sqrt{1-g_M^2}} V_\mu^i \\ W_\mu^i \rightarrow W_\mu^i - \frac{g_M}{\sqrt{1-g_M^2}} V_\mu^i \end{array} \right.$$

matching onto Warsaw basis, up to 1-loop

SMEFT fit \rightarrow constraints on g_i couplings

Example: SM + vector triplet

Brivio, Bruggisser, Geoffroy, Krämer, Plehn, Summ in preparation

$$\begin{aligned} \mathcal{L}_V = & -\frac{1}{4} V_{\mu\nu}^i V^{i\mu\nu} - \frac{g_M}{2} V_{\mu\nu}^i W^{i\mu\nu} + \frac{m_V^2}{2} V_\mu^i V^{i\mu} + \frac{g_H}{2} V_\mu^i (H^\dagger i \overleftrightarrow{D}^{i\mu} H) \\ & + \frac{g_l}{2} V_\mu^- \bar{\ell} \gamma^\mu \sigma^i \ell + \frac{g_q}{2} V_\mu^- \bar{q} \gamma^\mu \sigma^i q + \frac{g_{VH}}{2} (H^\dagger H) V_\mu^i V^{i\mu} \end{aligned}$$

e.g. $Q_{HI}^{(3)} = (H^\dagger i \overleftrightarrow{D}^{i\mu} H) (\bar{\ell} \gamma^\mu \sigma^i \ell)$

$$\begin{aligned} (C_{HI}^{(3)})_{ij} = & -\frac{g_l g_H}{4m_V^2} \delta_{ij} + \frac{1}{36864\pi^2 m_V^2} \frac{\delta_{ij}}{1-g_M^2} \left[g_w^4 (288 + 1531g_M^2 + 2989g_M^4) \right. \\ & + g_w^3 (2642g_H g_M + 2340g_l g_M + 7942g_H g_M^3 + 6732g_l g_M^3) \\ & + g_w^2 (g_l^2 (-102 + 3054g_M^2) + g_H^2 (49 + 5711g_M^2)) \\ & + g_w g_M (1080g_H^3 + 5400g_H^2 g_l + 2304g_H g_l^2 + 432g_l^3 + 1440h_H g_{VH} + 1440g_l g_{VH}) \\ & + g_H g_l (1080g_H^2 - 360g_H g_l + 432g_l^2 + 1440g_{VH} + (1 + g_w^2)(2160 + 12600g_M^2)) \\ & \left. + 1440g_M^2 g_{VH} \right] + \frac{3}{3032\pi^2 m_V^2} (g_l - g_H)(g_l + g_w g_M)(Y_e Y_e^\dagger)_{ij} \end{aligned}$$

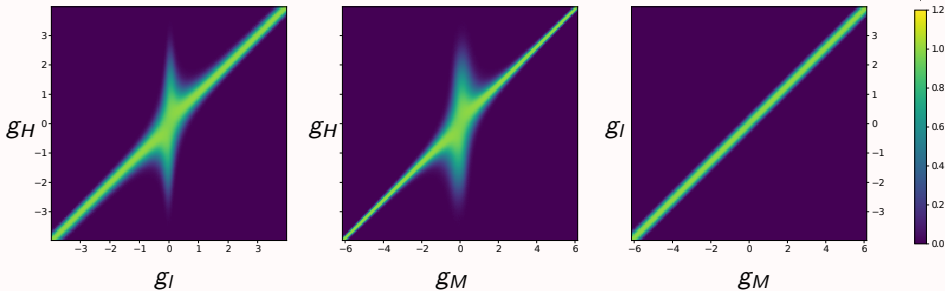
Example: SM + vector triplet

$$\mathcal{L}_V = -\frac{1}{4} V_{\mu\nu}^i V^{i\mu\nu} - \frac{g_M}{2} V_{\mu\nu}^i W^{i\mu\nu} + \frac{m_V^2}{2} V_\mu^i V^{i\mu} \\ + \frac{g_H}{2} V_\mu^i (H^\dagger i \overleftrightarrow{D}^{i\mu} H) + \frac{g_l}{2} V_\mu^- \bar{\ell} \gamma^\mu \sigma^i \ell$$

fit to EWPO+EW+Higgs preliminary!

$m_V = 2 \text{ TeV}$, tree-level matching

scaled likelihood



all g_i unconstrained

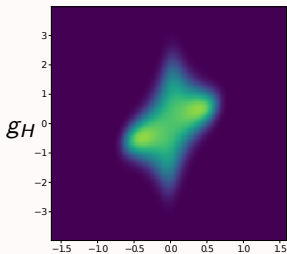
Example: SM + vector triplet

$$\mathcal{L}_V = -\frac{1}{4} V_{\mu\nu}^i V^{i\mu\nu} - \frac{g_M}{2} V_{\mu\nu}^i W^{i\mu\nu} + \frac{m_V^2}{2} V_\mu^i V^{i\mu} \\ + \frac{g_H}{2} V_\mu^i (H^\dagger i \overleftrightarrow{D}^{i\mu} H) + \frac{g_l}{2} V_\mu^- \bar{\ell} \gamma^\mu \sigma^i \ell$$

fit to EWPO+EW+Higgs preliminary!

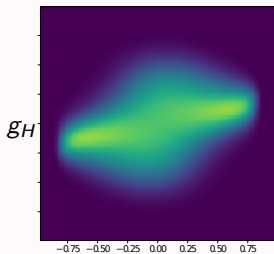
$m_V = 2 \text{ TeV}$, 1-loop matching

scaled likelihood



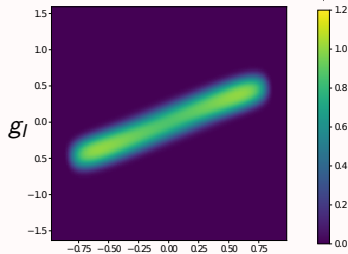
g_I

$$|g_I| \lesssim 0.7$$



g_M

$$|g_H| \lesssim 2.5$$

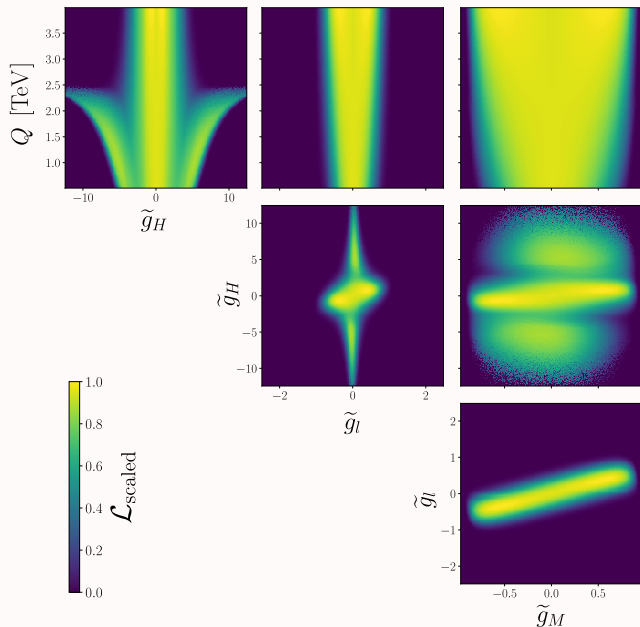


g_M

$$|g_M| \lesssim 0.8$$

SM + vector triplet: matching scale dependence

PRELIMINARY



$$m_V = 4 \text{ TeV}$$

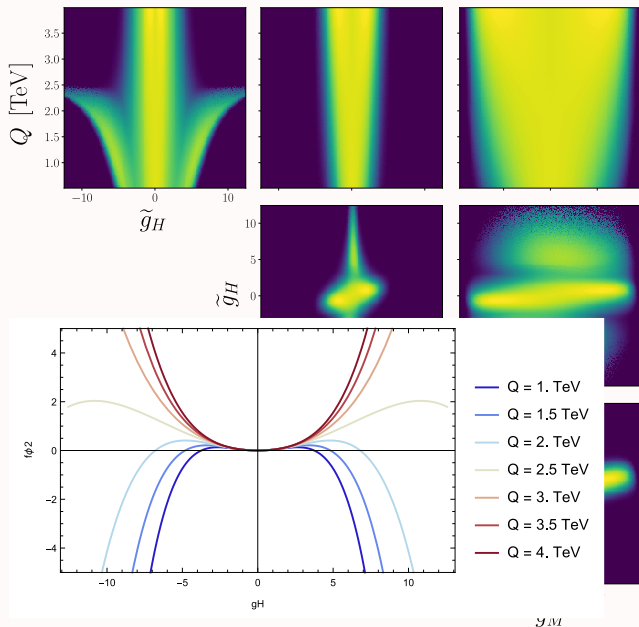
1-loop matching

SM + vector triplet: matching scale dependence

$m_V = 4 \text{ TeV}$

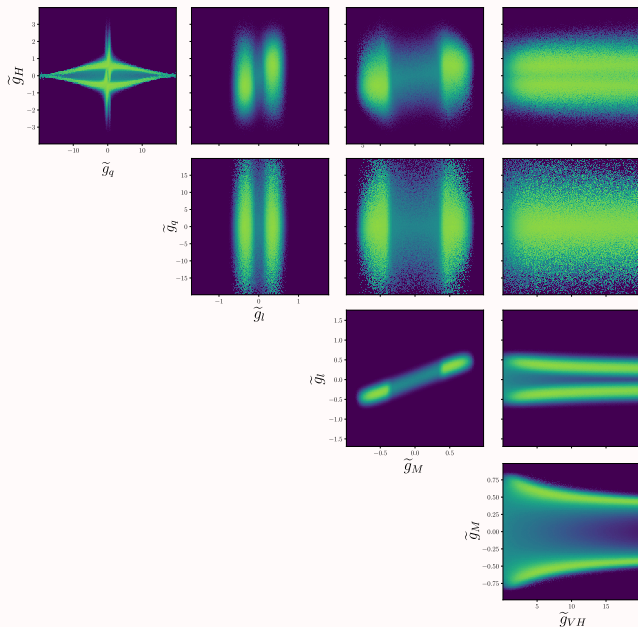
1-loop matching

PRELIMINARY



SM + vector triplet: 5-parameter fit

PRELIMINARY



$$m_V = 2 \text{ TeV}$$
$$Q = m_V$$

1-loop matching

calculation

measurement

global analysis

interpretation

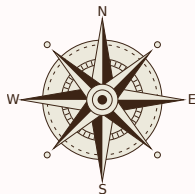
- ▶ treatment of SMEFT uncertainties, \mathcal{L}_8 terms
- ▶ anomaly constraints?
- ▶ automation of NLO SMEFT EW
- ▶ more NLO SMEFT results
- ▶ machine learning predictions?
- ▶ SMEFT effects in PDFs, hadronization, extraction of $\alpha_s \dots$

- ▶ SMEFT or HEFT?
- ▶ how to combine with direct searches?
- ▶ 1-loop matching dictionary
- ▶ 2-loop RG running

- ▶ fits with 50-60 parameters
- ▶ combination EW+Higgs+top+flavor
- ▶ combinations within the experiments
- ▶ improved treatment of correlations, uncertainties etc

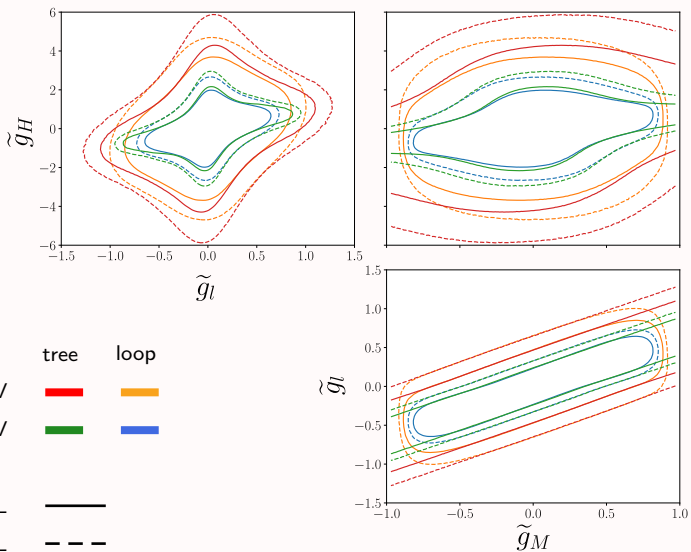
Summary

- ▶ SMEFT is the best framework for indirect searches of new physics at LHC
 - current direct bounds consistent with $(v/\Lambda_{BSM}) \ll 1$
 - collider physics entering **precision era**
- ▶ A challenging task! Developments are needed in
 - 🎯 precision calculations
 - 📊 numerical predictions & analysis strategies
 - 📈 statistical analyses
 - 💡 theory understanding
 - ...
- ▶ A community effort: combining is key
 - both theory and experiment increasingly involved
 - first exp. combination **EW + Higgs + top** sectors upcoming
 - incorporation of constraints from **flavor** experiments
 - complementarity with **direct** searches

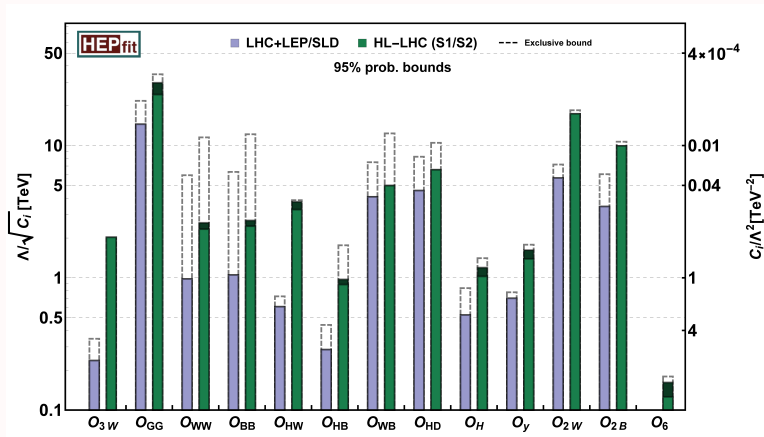


Backup slides

SM + vector triplet: contours for 3D fit

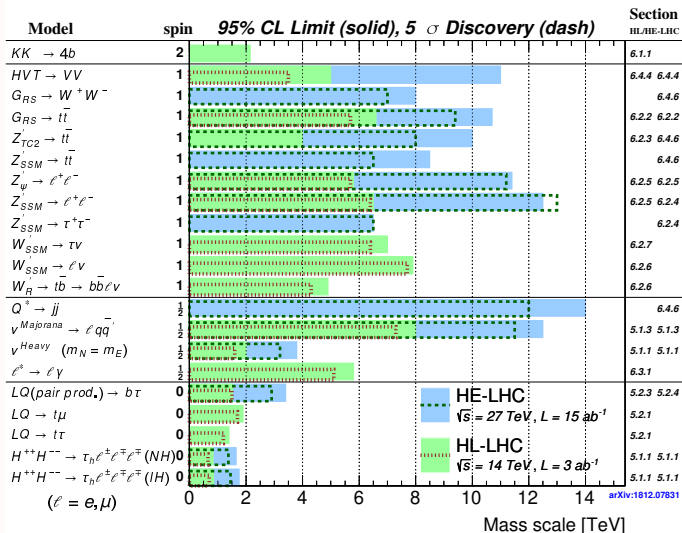


Prospects for HL-LHC



Yellow report on LH-LHC and HE-LHC physics 1902.00134

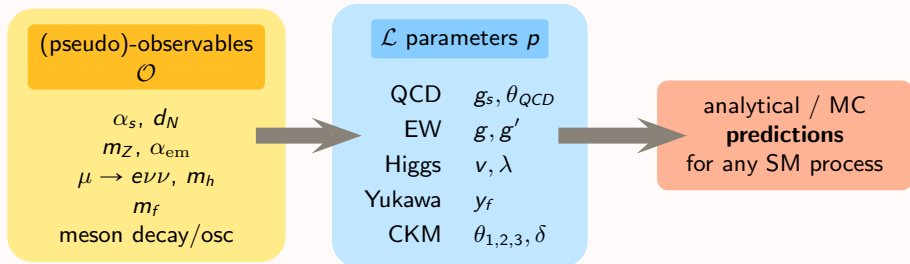
Prospects for HL-LHC



Yellow report on LH-LHC and HE-LHC physics 1812.07831

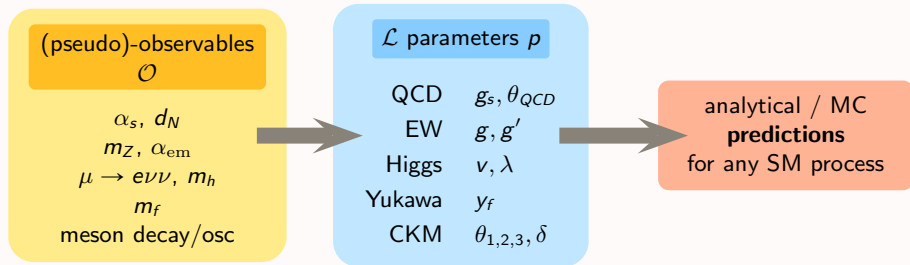


SM





SM

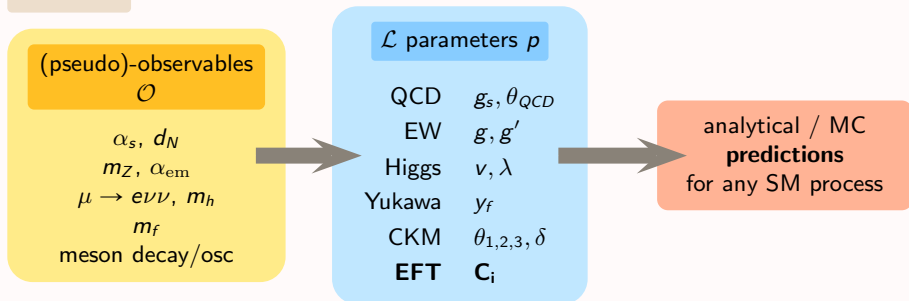


usual procedure

1. choose set \mathcal{O} to define all p unambiguously
2. calculate $\mathcal{O}(p)$ at given order
3. invert the relations $\rightarrow p(\mathcal{O})$
4. the renormalized p has a well-defined numerical value



SMEFT



SMEFT procedure

1. choose set \mathcal{O} to define all p_{SM} unambiguously
2. calculate $\mathcal{O}(p_{SM}, \mathbf{C}_i)$ at given order
3. invert the relations expanding around the SM solution
 $\rightarrow p_{SM}(\mathcal{O}) + \delta p(\mathbf{C}_i)$
4. p_{SM} has a numerical value, δp encodes EFT corrections

Top fit – observables

$pp \rightarrow t\bar{t}$

- ▶ 5 $\sigma_{t\bar{t}}$ measurements at 8 and 13 TeV
- ▶ 5 A_C measurements at 8 and 13 TeV
- ▶ 2 $d\sigma/dm_{t\bar{t}}$ dist. at 8 and 13 TeV (15 bins tot)
- ▶ 4 $d\sigma/dp_T^t(p_T^l, p_T^h)$ dist. at 8 and 13 TeV (30 bins tot)
- ▶ 1 $d^2\sigma/dm_{t\bar{t}}dt_{t\bar{t}}$ dist at 8 TeV (16 bins)
- ▶ 2 dist in high- p_T region ($p_T^t, m_{t\bar{t}}$) at 8 and 13 TeV (13 bins tot)

$pp \rightarrow t\bar{t}Z, pp \rightarrow t\bar{t}W$

- ▶ 2 $\sigma_{t\bar{t}V}$ measurements for each V at 8 and 13 TeV

Single-top

- ▶ 6 $\sigma_{tq, \bar{t}q}$ measurements in t -channel at 7, 8, 13 TeV
- ▶ 3 $\sigma_{t\bar{b}, \bar{t}b}$ measurements in s -channel at 7, 8 TeV
- ▶ 6 $\sigma_{tW, \bar{t}W}$ measurements in tW channel at 7, 8, 13 TeV
- ▶ 1 σ_{tZq} measurement in tZq at 13 TeV

Top decays

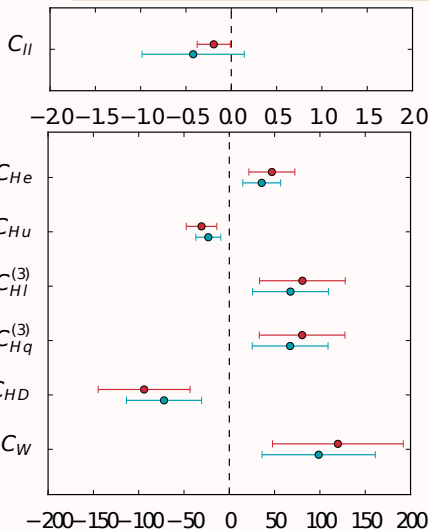
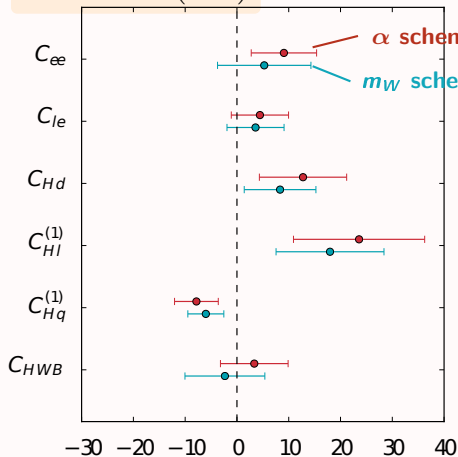
- ▶ 4 measurements of W helicity at 7, 8, 13 TeV



Z-pole + m_W + bhabha + WW (LEP2)

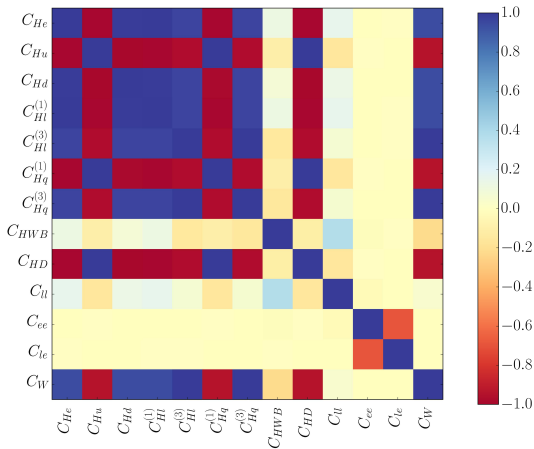
Berthier, (Bjørn), Trott 1508.05060, 1606.06693
 Brivio, Trott 1701.06424

1σ regions for $(C_i \frac{v^2}{\Lambda^2})$





Correlation matrix (α scheme)



2 blind directions
in EWPD

Grojean, Skiba, Terning 0602154



reparameterization
invariance
of 2 \rightarrow 2 scattering

Brivio, Trott 1701.06424

WW breaks (weakly) the
invariance
leaving strong correlations



Production

Brivio,Hays,Smith,Trott,Žemaitytė in preparation

- ▶ $gg \rightarrow h$ █ known to NLO SMEFT
- ▶ $qq \rightarrow qqh$ (VBF/VH) █ parton level inferred from $h \rightarrow 4l$ via crossing sym.
- ▶ $qq/gg \rightarrow hll/hl\nu$ (VH) █ in progress
- ▶ $gg \rightarrow t\bar{t}h$
- ▶ $qq \rightarrow thj$

Manohar,Wise 0601212
 Deutschmann,Duhr,Maltoni,Vryonidou 1708.00460
 Grazzini,Ilnińska,Spira 1806.08832
 $h \rightarrow 4l$ via crossing sym.

Maltoni,Vryonidou,Zhang 1607.05330

Degrande,Maltoni,Mimasu,Vryonidou, Zhang 1804.07773

Acceptance

$$A = \frac{n_{\text{kin.cuts}}}{n_{\text{tot}}} \quad \text{assumed to be SM-like in STXS extraction}$$

- ▶ SMEFT terms with **non-SM Lorentz** structure ($hV_{\mu\nu}V^{\mu\nu}$, $hV_{\mu}\bar{\psi}\gamma^{\mu}\psi \dots$) modify distributions \rightarrow ΔA
- ▶ ΔA calculable for cuts in Lorentz-invariants, requires MC for arbitrary cuts