

*Christoph Englert*

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# Effective Field Theory for Higgs and Top Physics

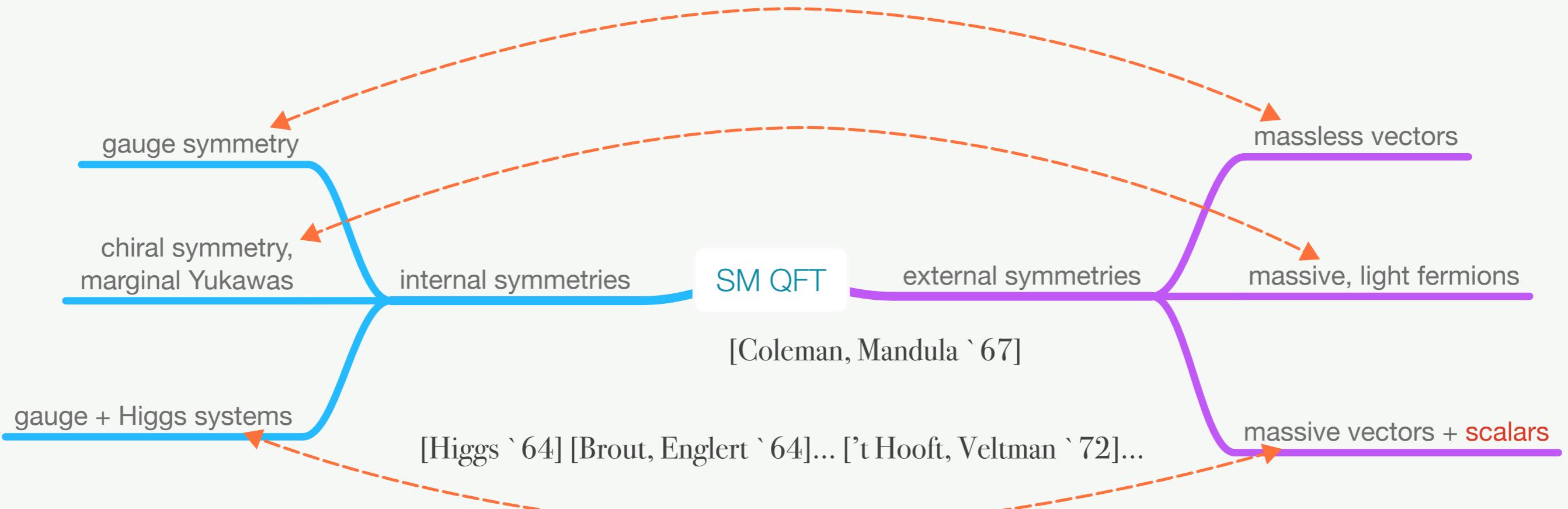
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*Oklahoma State University*

30/09/21

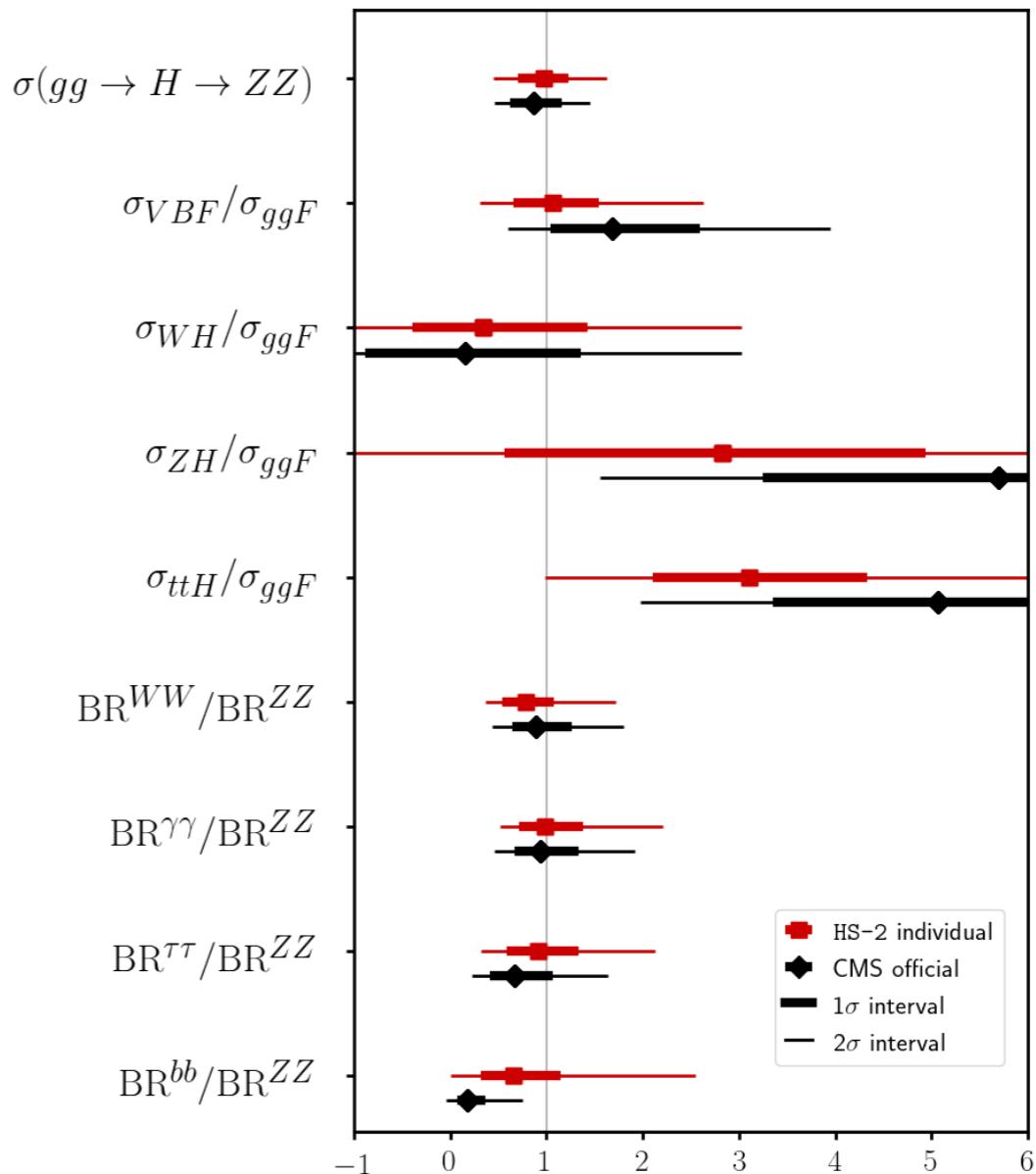
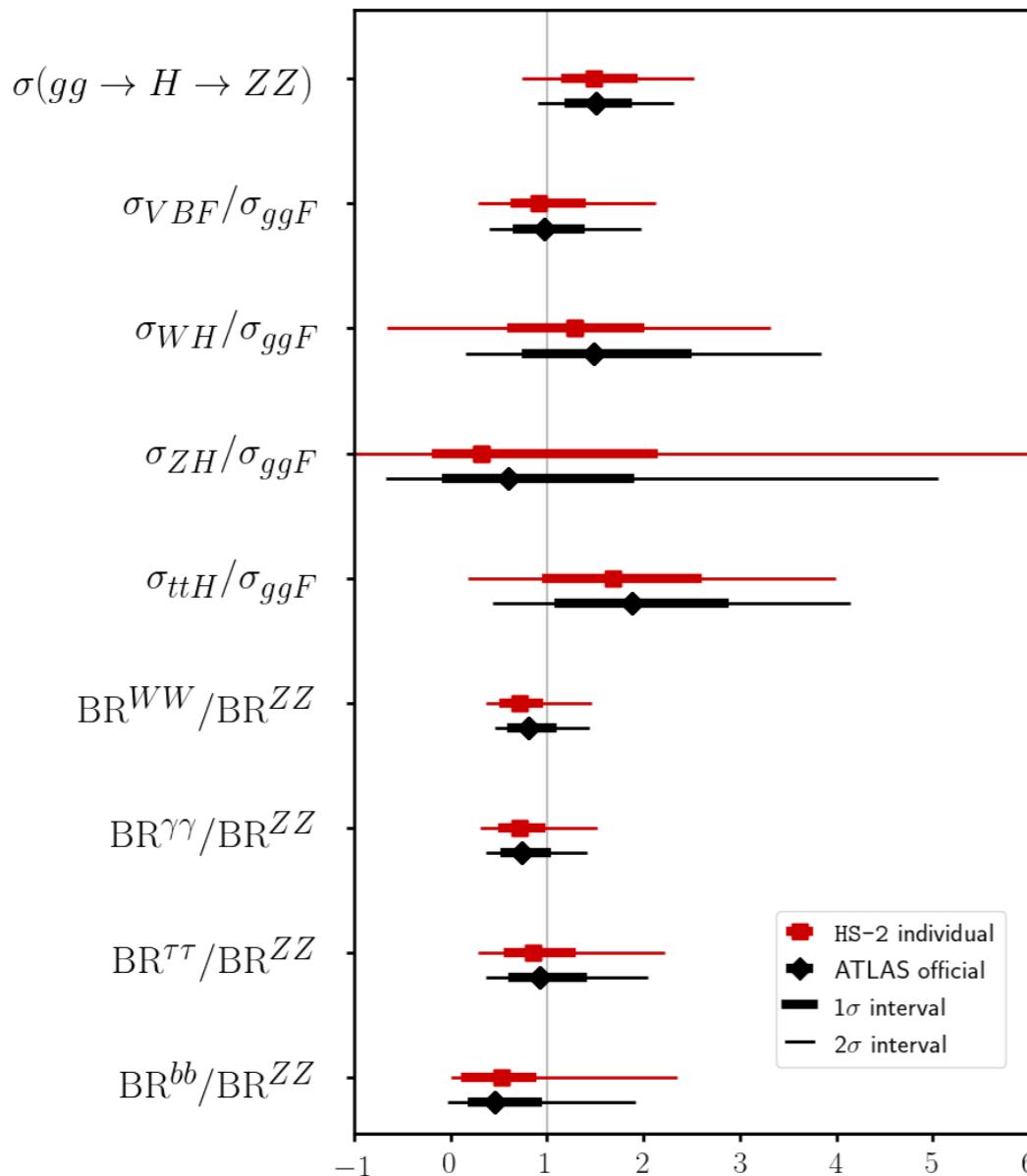
# The Standard Model: taking stock

...[Yang, Mills '54] [Mandelstam]...



“What’s next ?”

# Status of LHC measurements



[Bechtle et al. '20]

- everything is consistent with the SM Higgs hypothesis (so far)  
but what are the implications for new physics?

# Fingerprinting the lack of new physics

the SM is flawed

no evidence for exotics

## coupling/scale separated BSM physics

### Effective Field Theory

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i + \dots$$

[Buchmüller, Wyler '87]

[Hagiwara, Peccei, Zeppenfeld, Hikasa '87]

[Giudice, Grojean, Pomarol, Rattazzi '07]

[Grzadkowski, Iskrzynski, Misiak, Rosiek '10]

[Brivio, Jiang, Trott '17]....

59 B-conserving operators  $\otimes$  flavor  $\otimes$  h.c., d=6  
2499 parameters (reduces to 76 with  $N_f=1$ )

### concrete models

- ▶ extended SMEFT
- ▶ (C) Higgs portals
- ▶ 2HDMs
- ▶ simplified models
- ▶ compositeness....

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*Are EFTs collider tools to improve on the expected and perhaps even observe the unexpected?*

- ▶ *CP violating Higgs interactions ?*
- ▶ *improving our understanding Higgs propagation ?*
- ▶ *BSM interplay of top/Higgs sectors ?*

# CP violation

- Higgs sector is a primary candidate for CP violation - how is this captured in a dimension 6 approach?

matching

MC  
perturbativity

unitarity...

$$\begin{aligned} d\sigma &\sim |\mathcal{M}_{\text{SM}}|^2 + 2\text{Re}\{\mathcal{M}_{\text{SM}}\mathcal{M}_{\text{d6}}^*\} + |\mathcal{M}_{\text{d6}}|^2 \\ &\sim \Lambda^0 \quad \sim \Lambda^{-2} \quad \sim \Lambda^{-4} \end{aligned}$$

dim 8

# CP violation

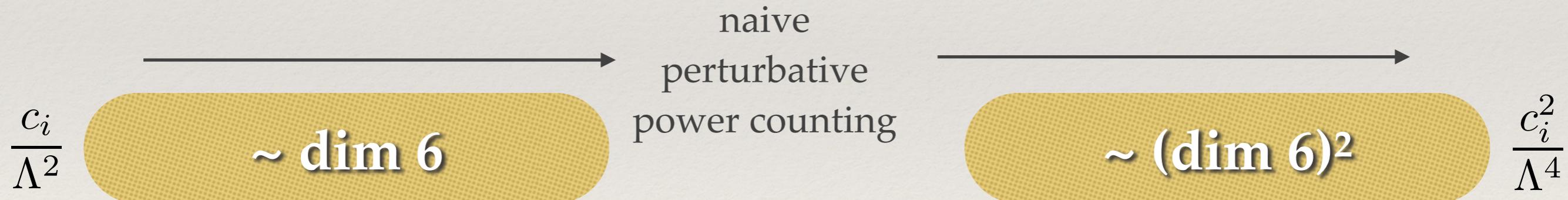
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  - matching
  - MC perturbativity
  - unitarity...
- in practice this is (often) not a huge problem for large data samples

# CP violation

- Higgs sector is a primary candidate for CP violation - how is this captured in a dimension 6 approach?



- in practice this is (often) not a huge problem for large data samples
- **but qualitatively different for CP-violation:**



- only genuinely CP-sensitive observables carry information
  - signed  $\Delta\phi_{jj}$ , asymmetries, ....
  - ...[Plehn et al. '01]... [Figy et al. '06]...
- every CP-even observable carries information
  - cross sections, widths, pt spectra...

# CP violation

[Bernlochner, CE, Hays, Lohwasser, Mildner, Pilkington, Price, Spannowsky '18]

- the linearised upshot

$$O_{H\tilde{G}} = H^\dagger H G^{a\mu\nu} \tilde{G}_{\mu\nu}^a,$$

$$O_{H\tilde{W}} = H^\dagger H W^{a\mu\nu} \tilde{W}_{\mu\nu}^a,$$

$$O_{H\tilde{B}} = H^\dagger H B^{\mu\nu} \tilde{B}_{\mu\nu},$$

$$O_{H\tilde{W}B} = H^\dagger \tau^a H B_{\mu\nu} \tilde{W}^{a\mu\nu}$$

+

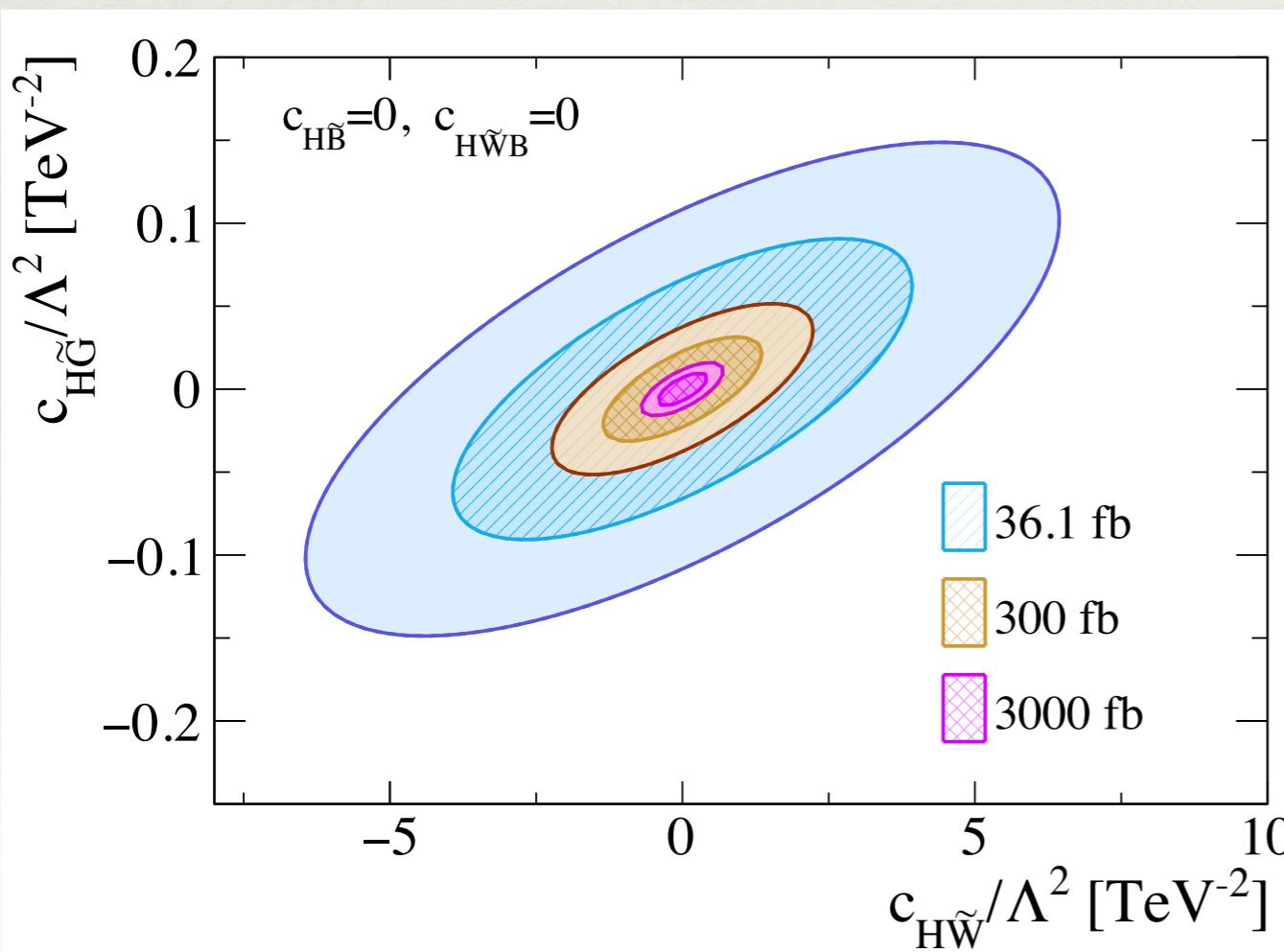
top quark

$$\sim \frac{\alpha_s}{8\pi v} G_{\mu\nu}^a \tilde{G}^{a\mu\nu} h = \tilde{O}_G$$

Yukawa phases

ignore here: Can be tackled in GF

[CE, Galler, Pilkington, Spannowsky '19]



LHC and HL-LHC  
extrapolations

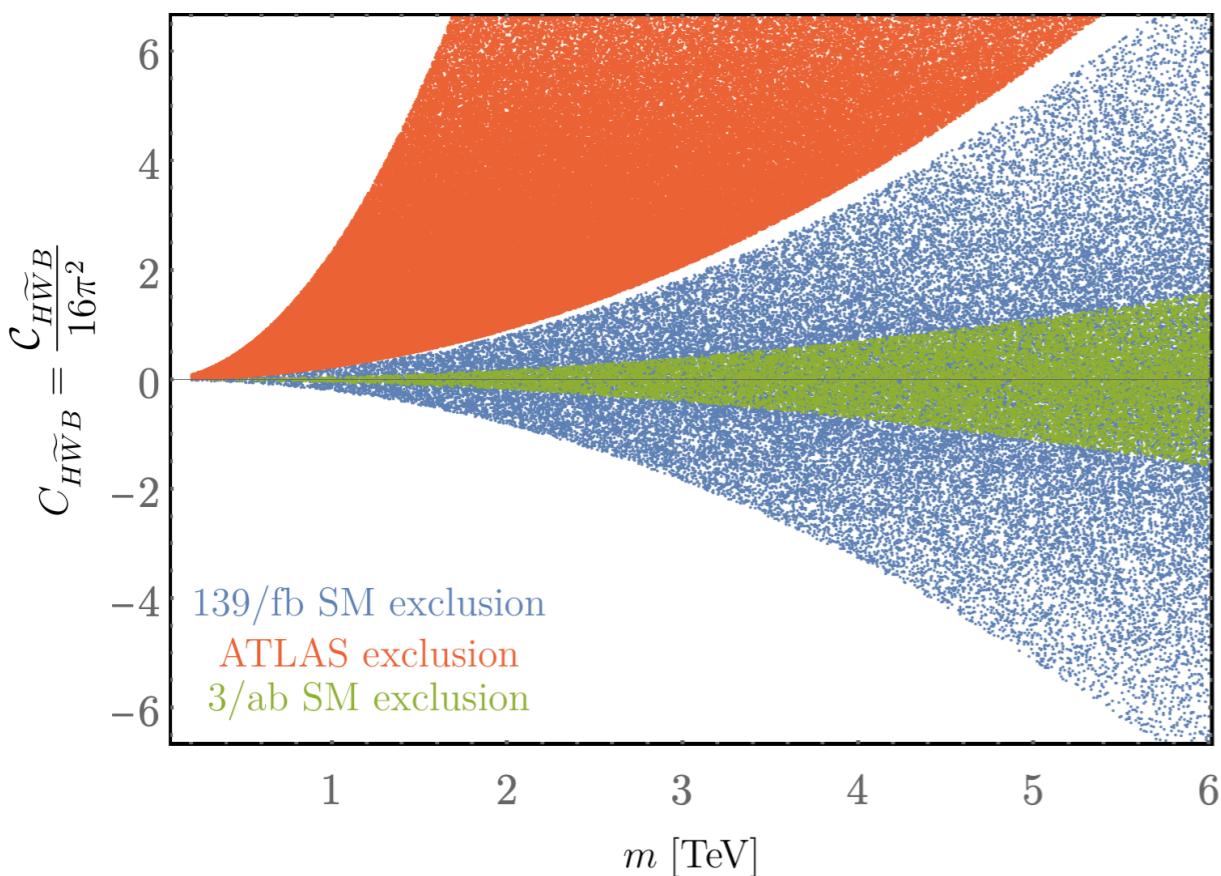
WBF+GF production and optimised  
selection, 4l final states

Coefficient [TeV <sup>-2</sup> ]	36.1 fb <sup>-1</sup>	300 fb <sup>-1</sup>	3000 fb <sup>-1</sup>
$c_{H\tilde{G}}/\Lambda^2$	[-0.19, 0.19]	[-0.067, 0.067]	[-0.021, 0.021]
$c_{H\tilde{W}}/\Lambda^2$	[-11, 11]	[-3.8, 3.8]	[-1.2, 1.2]
$c_{H\tilde{B}}/\Lambda^2$	[-5.9, 5.9]	[-2.1, 2.1]	[-0.65, 0.65]
$c_{H\tilde{W}B}/\Lambda^2$	[-14, 14]	[-4.9, 4.9]	[-1.5, 1.5]

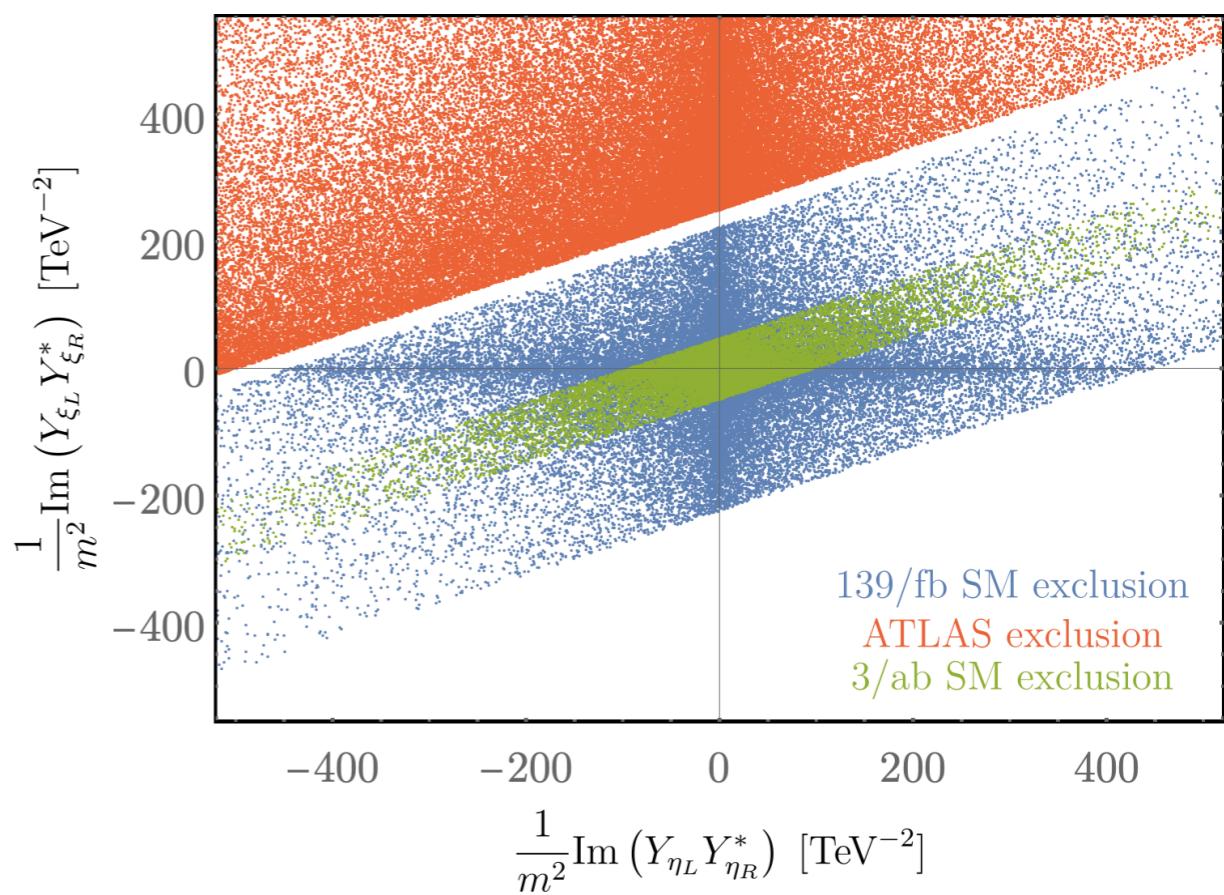
# CP violation

[ATLAS, 2006.15458]

Wilson coefficient	Includes $ \mathcal{M}_{d6} ^2$	95% confidence interval [TeV $^{-2}$ ] Expected	95% confidence interval [TeV $^{-2}$ ] Observed	$p$ -value (SM)
$c_W/\Lambda^2$	no	[-0.30, 0.30]	[-0.19, 0.41]	45.9%
	yes	[-0.31, 0.29]	[-0.19, 0.41]	43.2%
$\tilde{c}_W/\Lambda^2$	no	[-0.12, 0.12]	[-0.11, 0.14]	82.0%
	yes	[-0.12, 0.12]	[-0.11, 0.14]	81.8%
$c_{HWB}/\Lambda^2$	no	[-2.45, 2.45]	[-5.73, 1.13]	29.0%
	yes	[-3.11, 2.10]	[-6.31, 1.01]	25.0%
$\tilde{c}_{HWB}/\Lambda^2$	no	[-1.06, 1.06]	[0.23, 2.34]	1.7%
	yes	[-1.06, 1.06]	[0.23, 2.35]	1.6%



[Das Bakshi, Chakrabortty, CE, Spannowsky, Stylianou `20]



- ▶ ATLAS see a tension related to CP violation in WBF  $Z$  production
- ▶ sign for hierarchical new physics beyond the SM ?

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- ▶ ATLAS see a tension related to CP violation in WBF  $Z$  production
- ▶ sign for hierarchical new physics beyond the SM ?

- ▶ what can be learned from this?

[Das Bakshi, Chakrabortty, CE, Spannowsky, Stylianou `20]

- ▶ Assumptions of two-parameter CP fits theoretically consistent in a wide class of vector-like leptons
- ▶ Hierarchy  $|C_{H\widetilde{W}B}|/\Lambda^2 > |C_{\widetilde{W}}|/\Lambda^2$  predicted in these scenarios
- ▶ broad UV assumptions reduce complexity of fit whilst facilitating matching more straightforwardly

# Overview

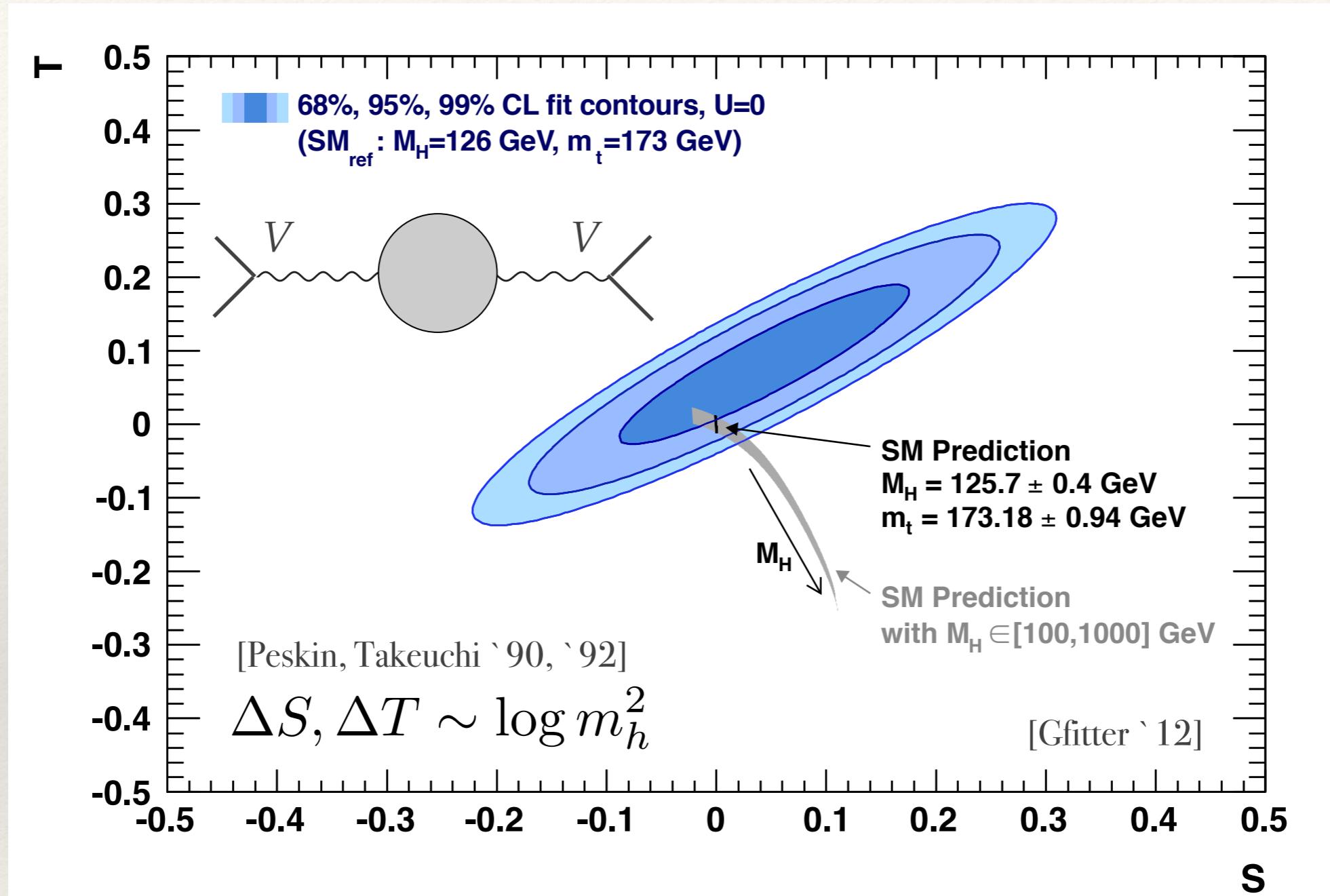
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- ▶ *CP violating Higgs interactions ?*
- ▶ *improving our understanding Higgs propagation ?*
- ▶ *BSM interplay of top/Higgs sectors ?*

# Higgs propagation

[CE, Giudice, Greljo, McCullough '19]



- specific dim 6 operators much better constrained than naively expected! Can we use similar tricks for the Higgs?

# Higgs propagation

[CE, Giudice, Greljo, McCullough '19]

- access oblique Higgs propagator corrections

$$\Delta_h(p^2) = \frac{1}{p^2 - m_h^2} - \frac{\hat{H}}{m_h^2} \quad \hat{H} = -\frac{m_h^2}{2} \Sigma''_h(m_h^2)$$

similar to

$$\dots \mathcal{L}_{\hat{W}} = -\frac{\hat{W}}{4m_W^2}(D_\rho W_{\mu\nu}^a)^2, \quad \mathcal{L}_{\hat{Y}} = -\frac{\hat{Y}}{4m_W^2}(\partial_\rho B_{\mu\nu})^2, \quad \mathcal{L}_{\hat{H}} = \frac{\hat{H}}{m_h^2} \parallel D^\mu D_\mu$$

[Barbieri et al. '04]

- excellent prospects to surpass LEP(2) sensitivity at high energy colliders due to scaling

$$\hat{T} = \mathcal{O}(q^0)$$

$$\hat{S} = \mathcal{O}(q^2)$$

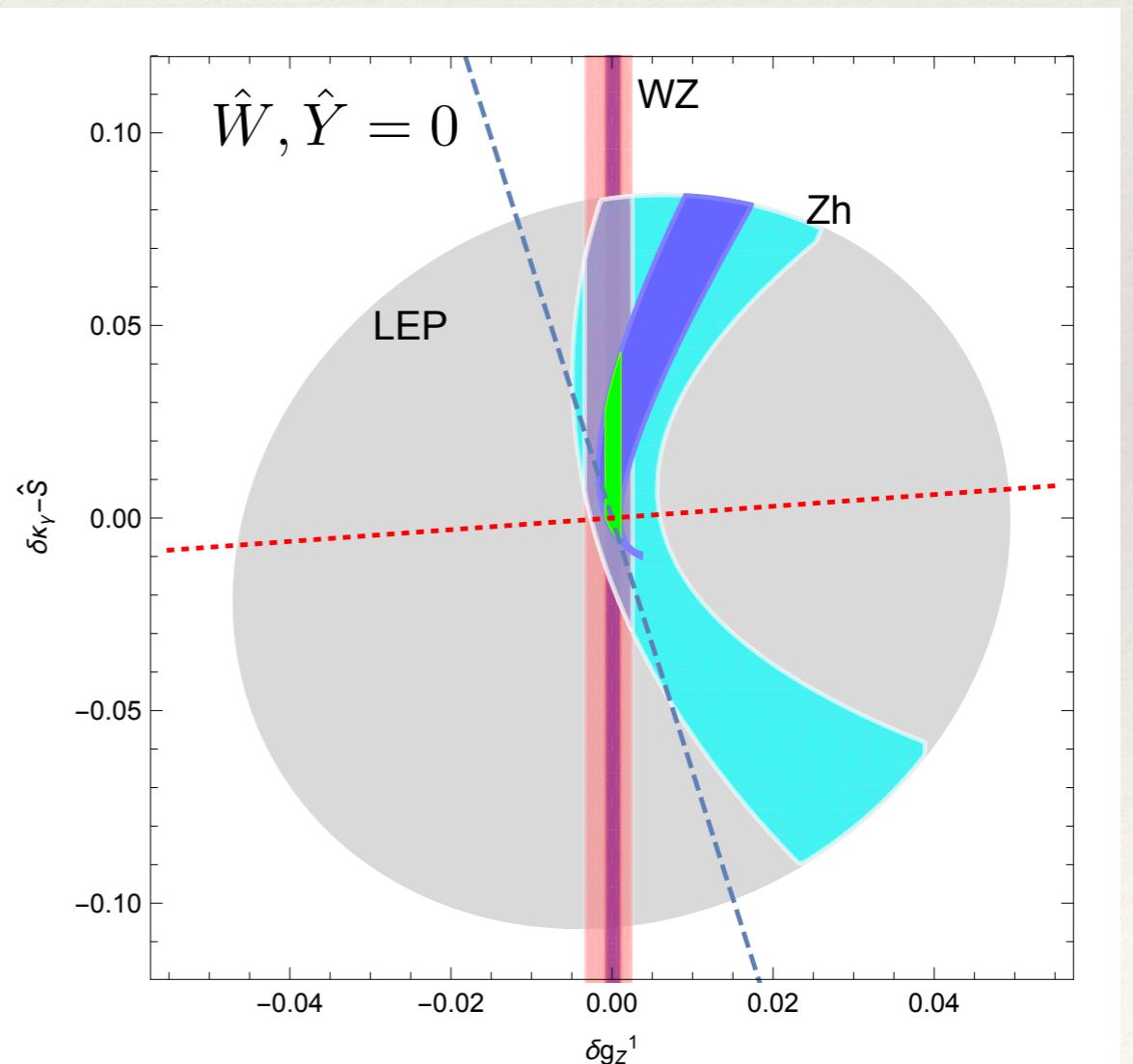
$$\hat{W}, \hat{Y} = \mathcal{O}(q^4)$$

[Farina et al. '17]

[Franceschini et al. '18]

[Banerjee, Gupta, CE, Spannowsky '18]

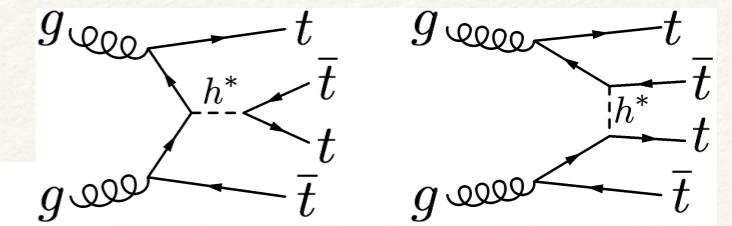
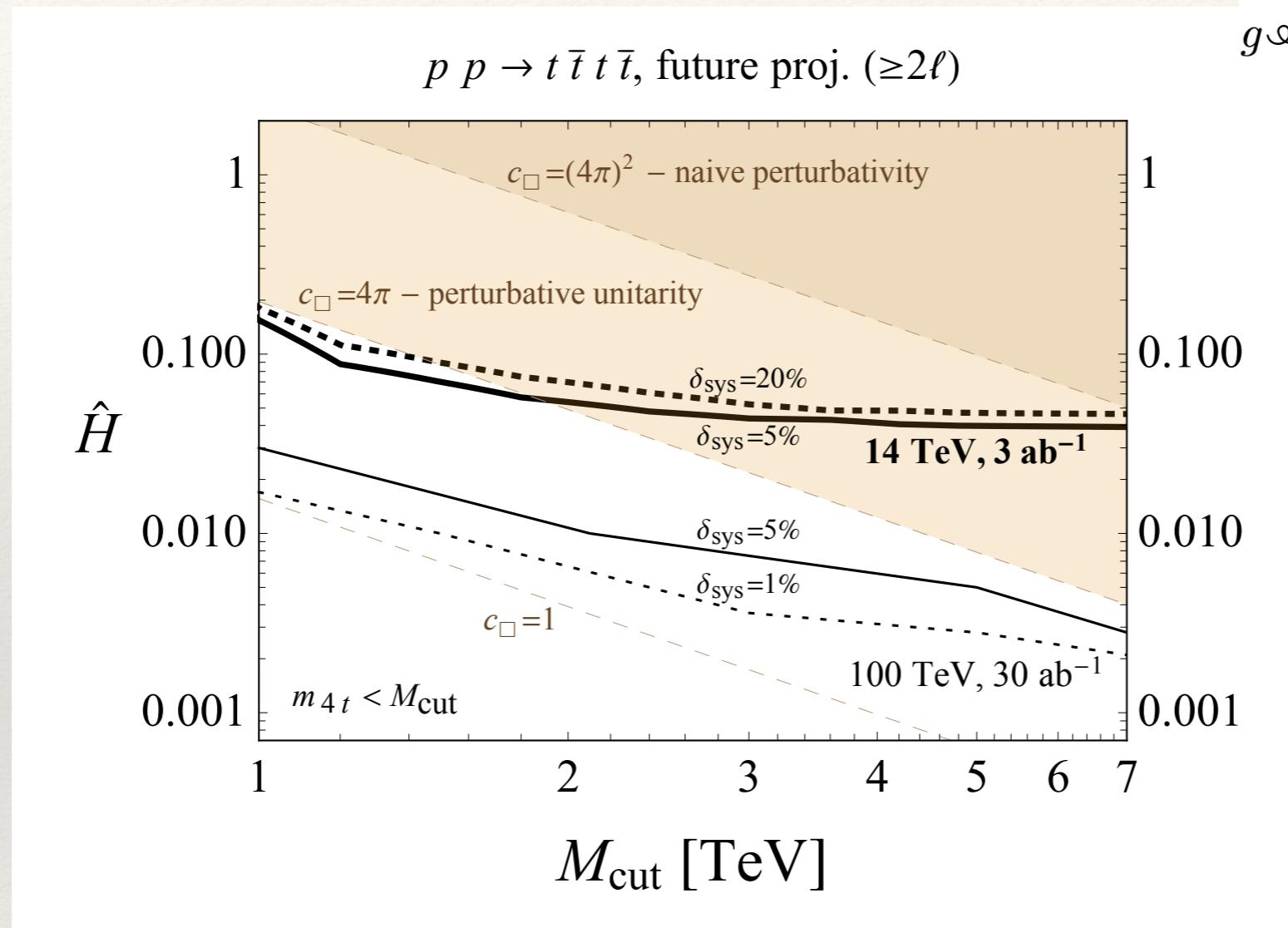
...



# Higgs propagation

[CE, Giudice, Greljo, McCullough '19]

- phenomenological details captured via



cf. [ATL-PHYS-PUB-2018-047]  
 [CMS-PAS-FTR-18-031]

- high energy frontier is an efficient probe at large cutoff

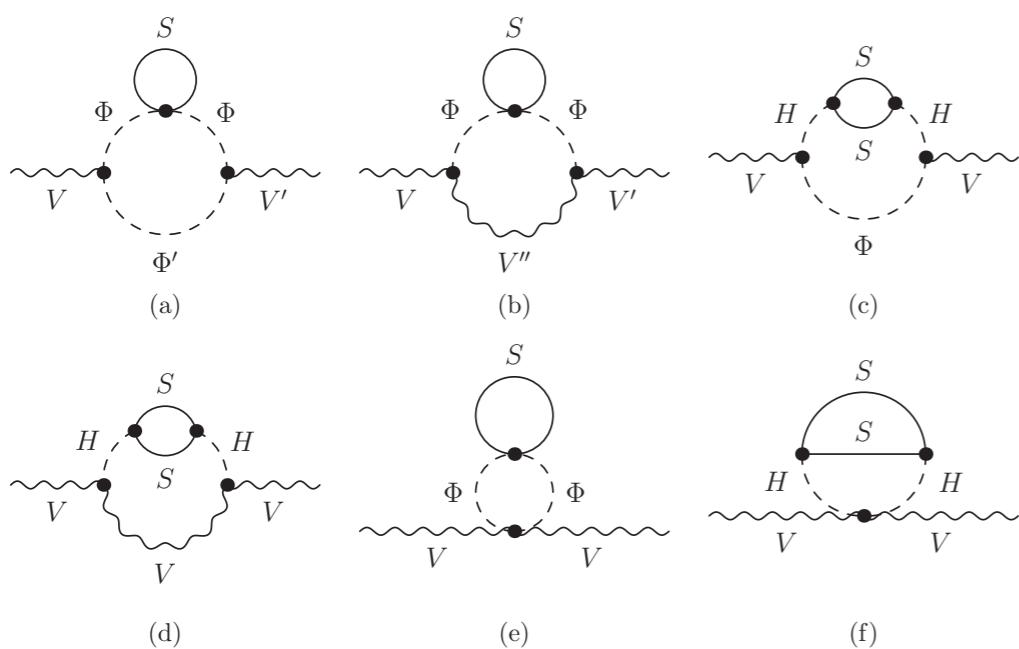
FCC-ee  $|\hat{H}| \lesssim 0.5\%$

[FCC Collaboration '19]

# Higgs propagation

...in loops...

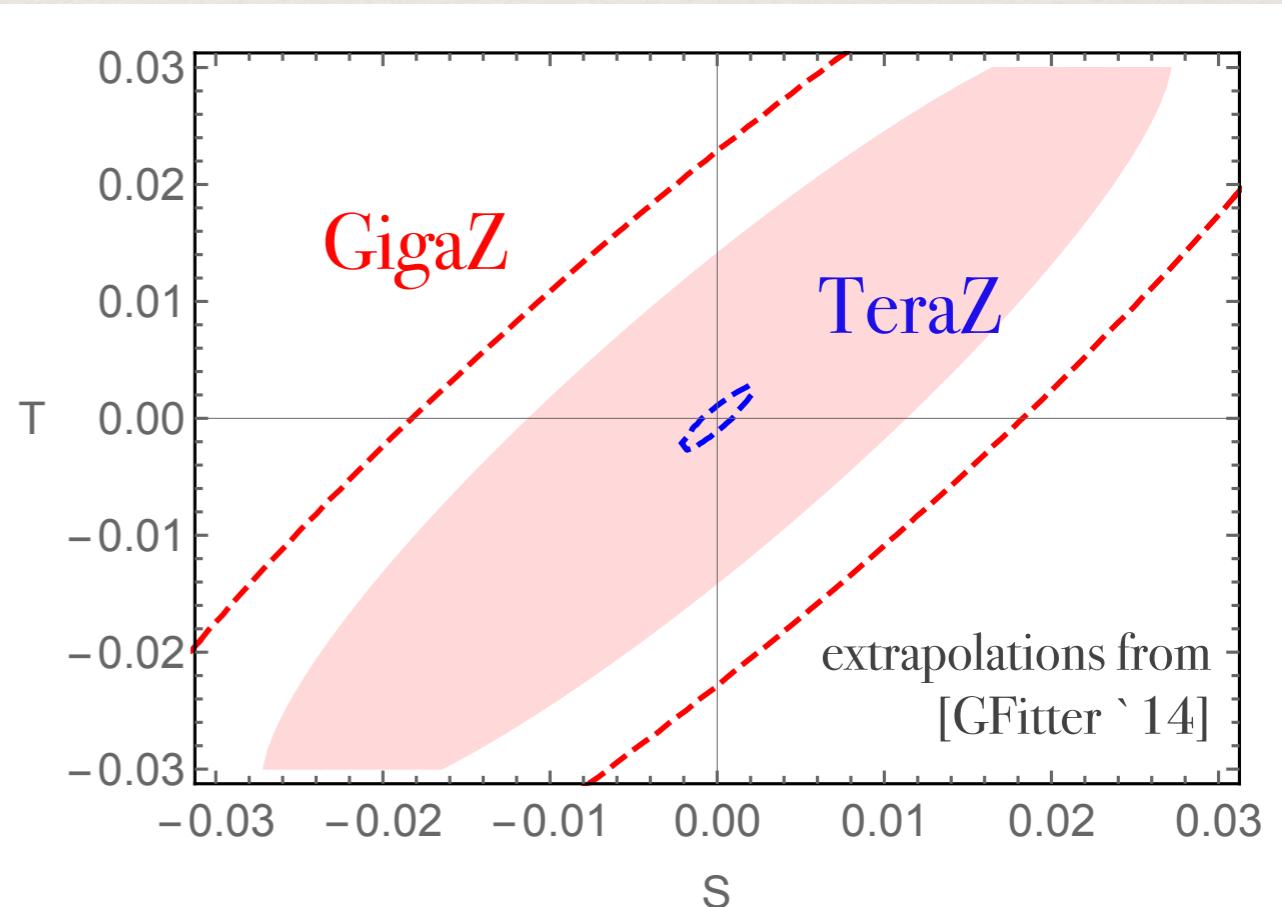
- precision analysis of  $Z$ -pole measurements ( $e^+ e^- \rightarrow ff'$ ) sensitive to Higgs corrections  $\mathcal{L} \supset -\lambda S^2(\Phi^\dagger \Phi - v^2/2)$  [CE, Jaeckel, Spannowsky, Stylianou '20]



GigaZ gives non-trivial constraint

massive improvement for TeraZ (if attainable)

- Oblique corrections suppressed, but large statistics and clean measurement at Higgs factories!



# Higgs propagation

singlets above threshold

- $\mathbb{Z}_2$ -symmetric Higgs portal

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{1}{2}(\partial_\mu S)^2 - \frac{m_S^2}{2}S^2 - \lambda S^2(\Phi^\dagger\Phi - v^2/2)$$

[Craig, Lou. et al. '14]  
[Curtin, Meade, Yu '14]

...

- for  $m_S > m_H/2$  no direct SM Higgs decays

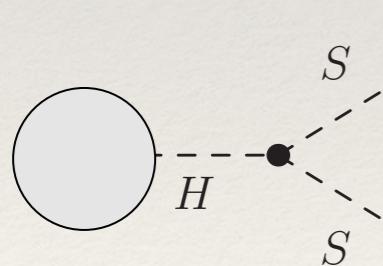
- BSM Higgs physics via momentum- or loop-suppressed effects

off-shell  
production

di-Higgs  
physics

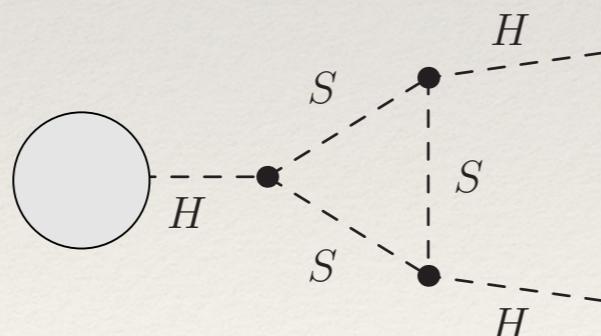
Higgs  
couplings

Electroweak  
precision



[Craig, Lou. et al. '14]

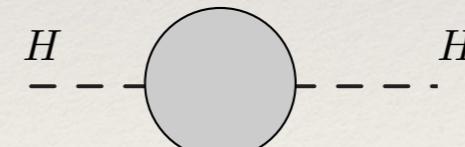
[Ruhdorfer, Salvioni, Weiler '19]



[Curtin, Meade, Yu '14]

[He, Zhu '16]

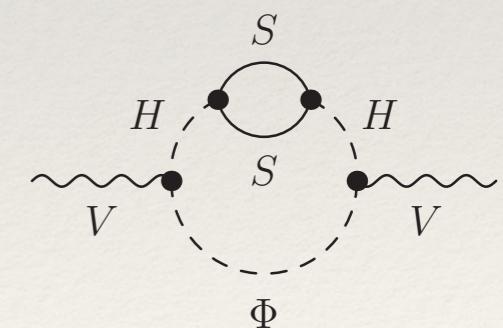
[Voigt, Westhoff '17]



[CE, McCullough '13]

[Craig, CE, McCullough '13]

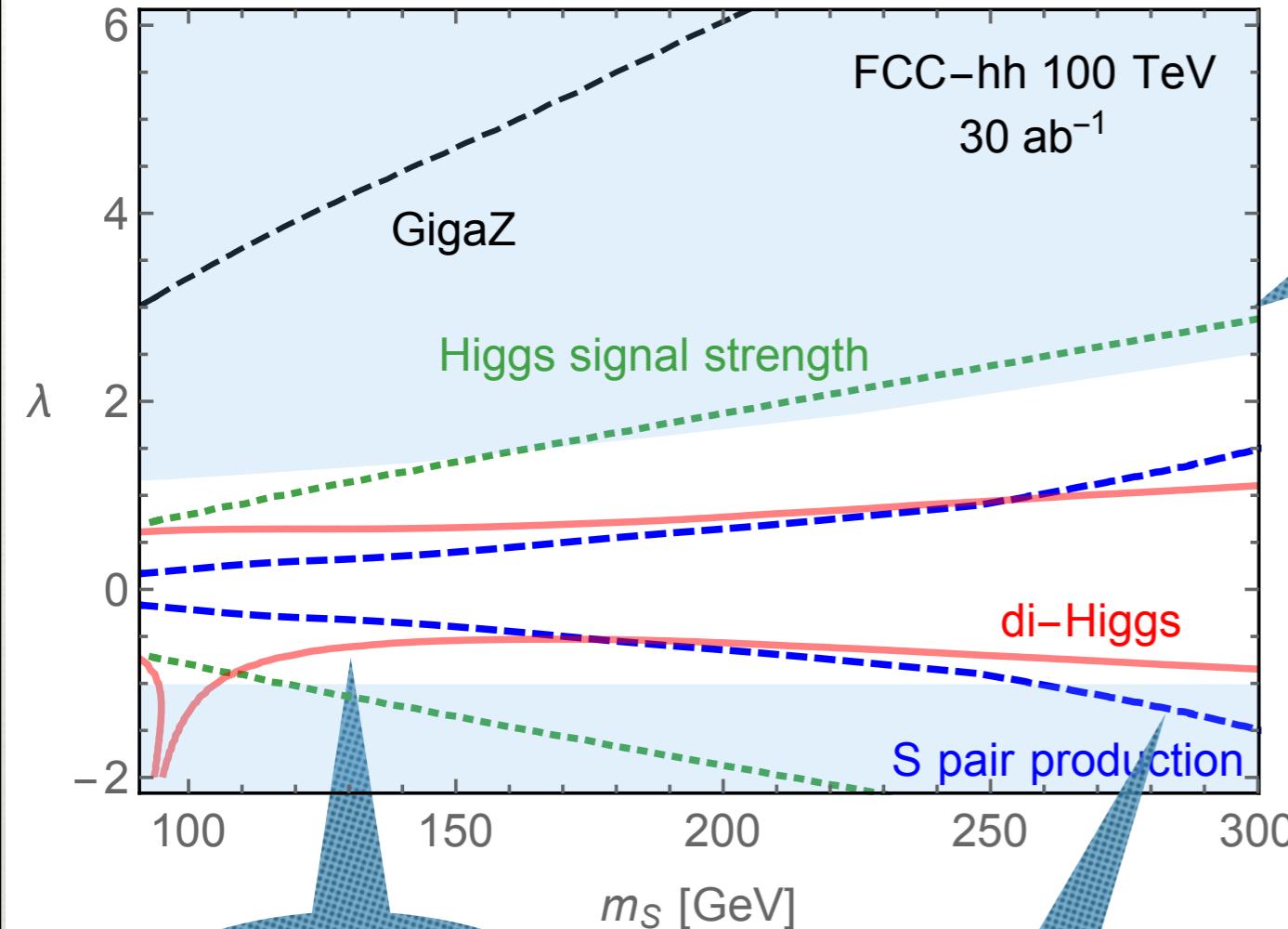
[Goncalves, Han, Mukhopadhyay '18]



[CE, Jaeckel, Spannowsky, Stylianou '20]

# Higgs propagation

[CE, Jaeckel, Spannowsky, Stylianou '20]



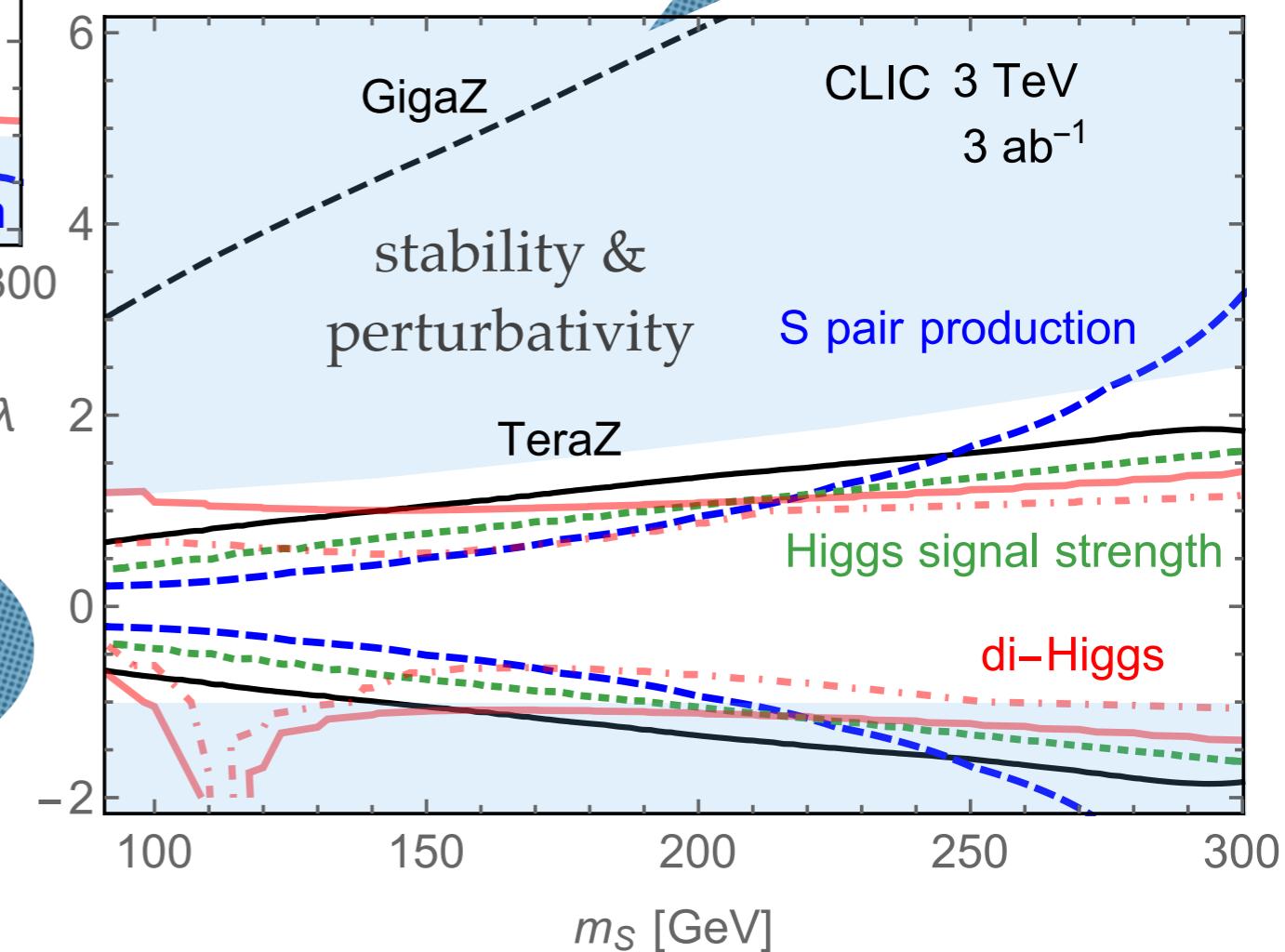
di-Higgs physics

off-shell production

# singlets above threshold

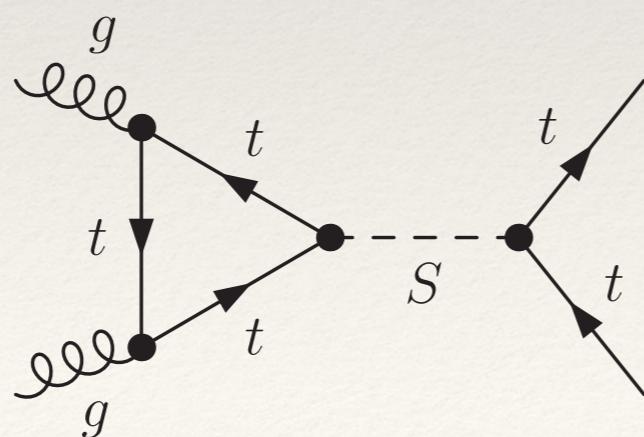
Higgs  
couplings

Oblique  
corrections



*Are EFTs collider tools to improve on the expected and perhaps even observe the unexpected?*

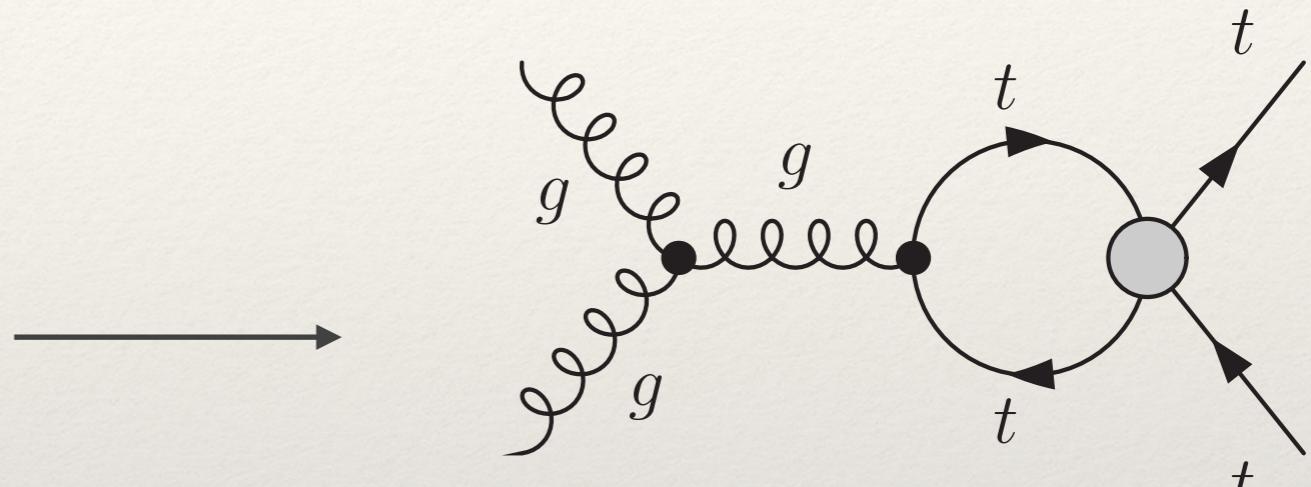
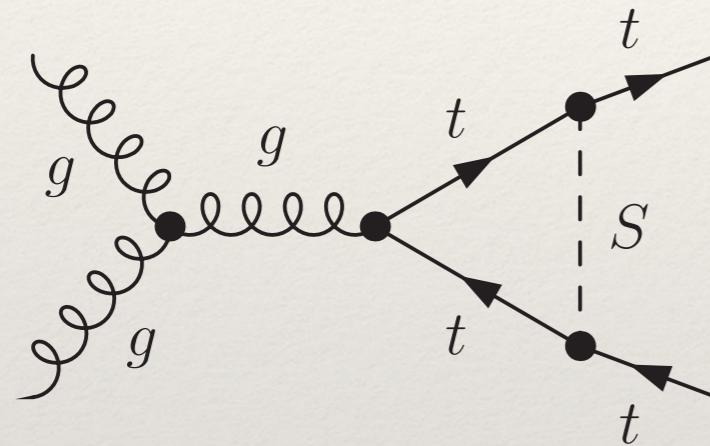
- ▶ *CP violating Higgs interactions ?*
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What do tops have  
to say about the presence  
of new scalar states?

# New physics in tops

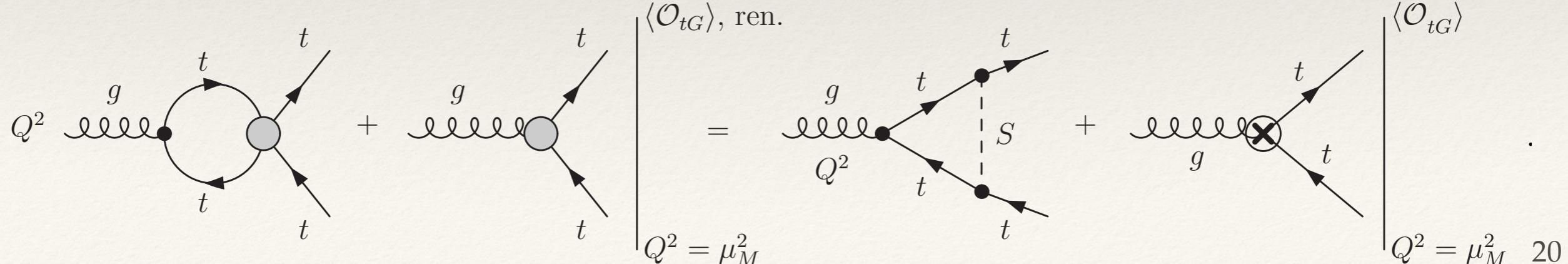
- new topophilic states arise in many BSM theories:  $- (c_S \bar{t}_L t_R S + \text{h.c.})$
  - top pair production with large cross section could fingerprint such states



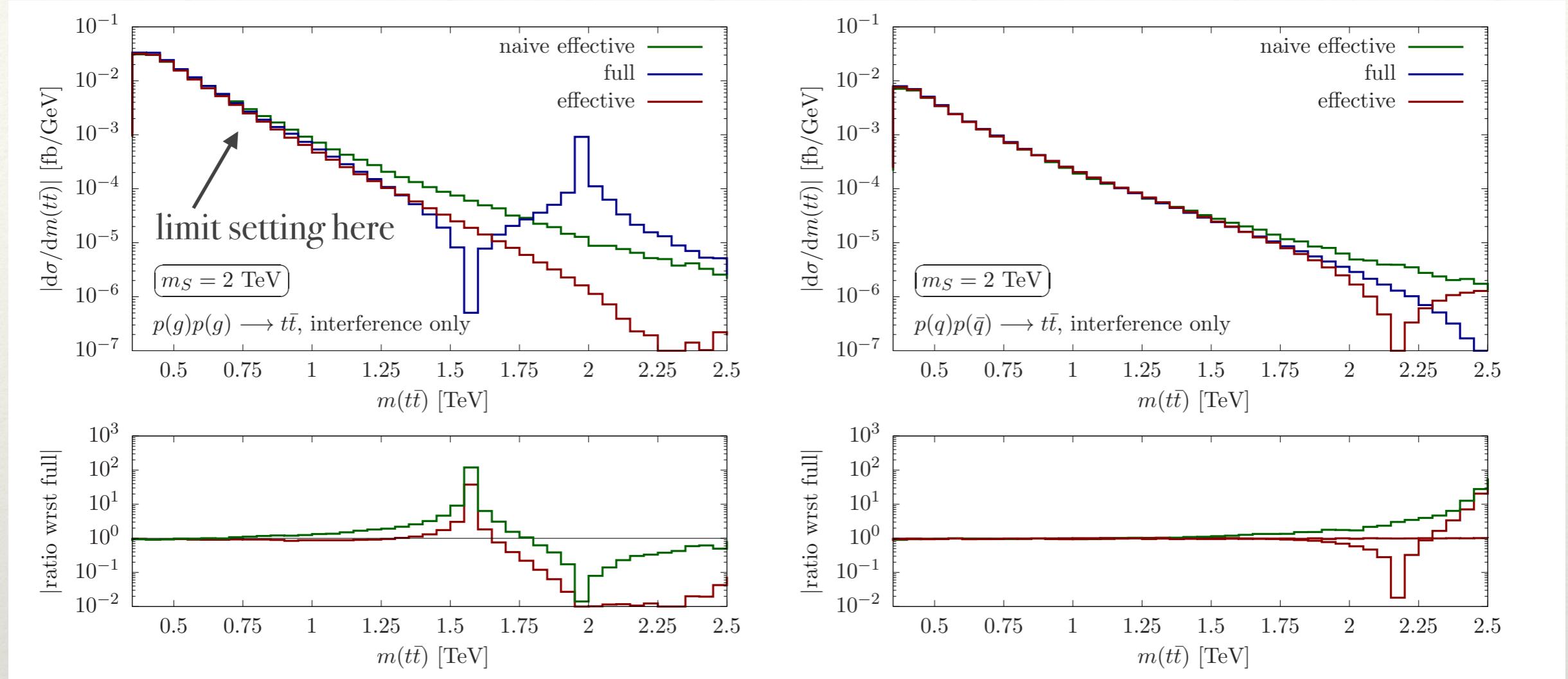
## *3 relevant operators*

- EFT is suitable tool to constrain such states model-independently,  
*however matching is crucial!* [CE, Galler, White]

[CE, Galler, White '19]

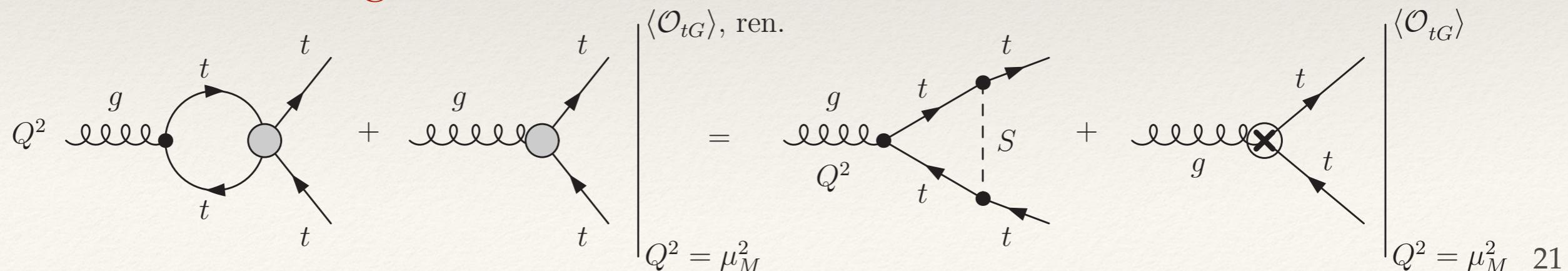


# New physics in tops



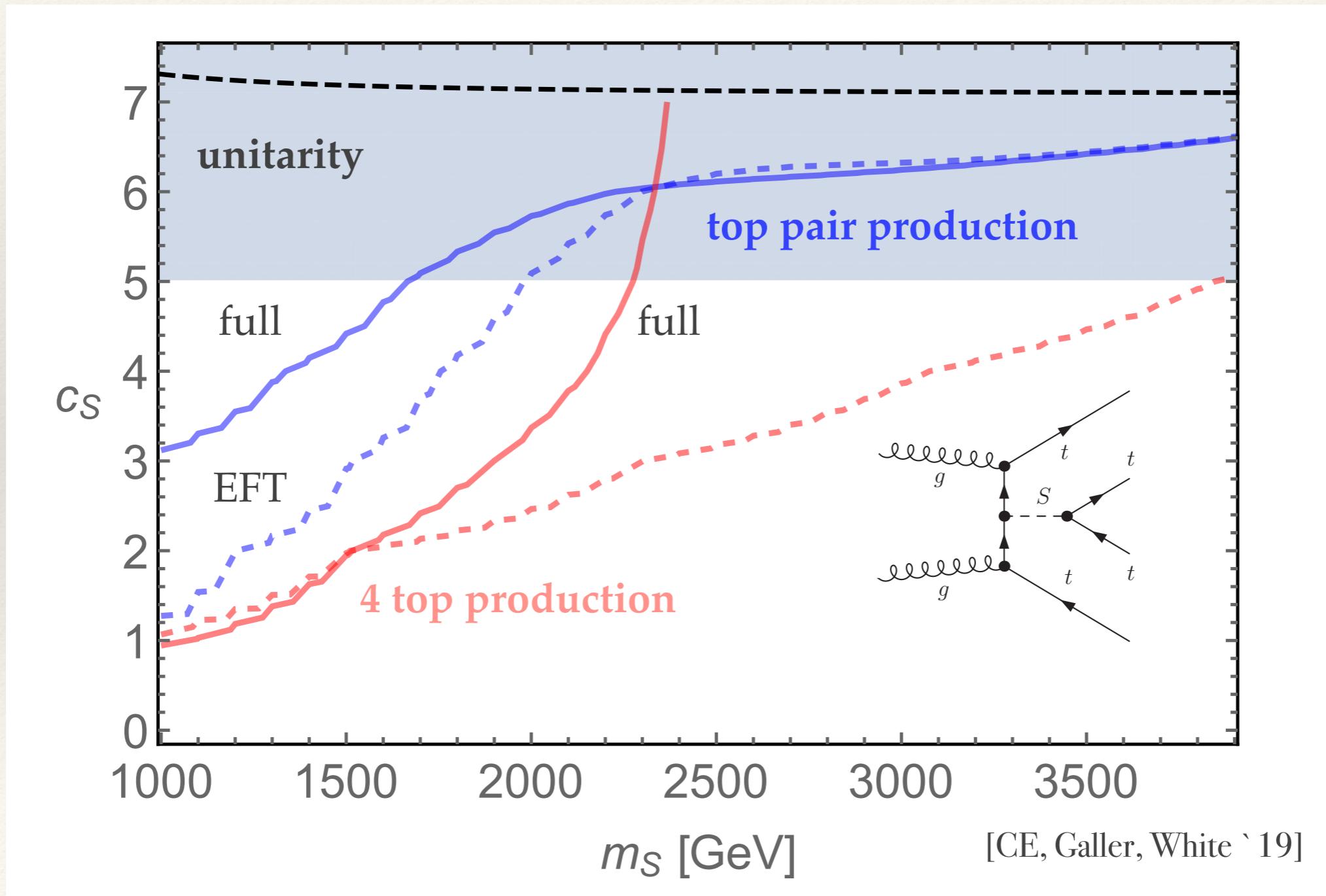
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*however matching is crucial!*

[CE, Galler, White '19]



# New physics in tops

- EFT is suitable tool to constrain such states model-independently,  
*however matching is crucial and so are expected uncertainties*



# Strong interactions? Compositeness....

---

- gauge boson masses through symmetry choices e.g. [Contino `10]
- fermion masses through mixing with baryonic matter (part. compositeness)
- minimal pheno model  $\text{SO}(5) \rightarrow \text{SO}(4) \simeq \text{SU}(2)_L \times \text{SU}(2)_R$
- fermions (and hypercolour baryons) in a 5 of  $\text{SO}(5)$

## A concrete model of compositeness

---

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  - fermions (and hypercolour baryons) in a 5 of  $\text{SO}(5)$
- so far no UV completion known for this!**
- but

$$\underbrace{\text{SU}(4) \times \text{SU}(5) \times \text{SU}(3) \times \text{SU}(3)' \times \text{U}(1)_X \times \text{U}(1)'}_{G_{\text{HC}}} \quad \underbrace{\text{SU}(3) \times \text{SU}(3)' \times \text{U}(1)'}_{G_F} \quad [\text{Ferretti '14}]$$

could work with

$$G_F/H_F = \frac{\text{SU}(5)}{\text{SO}(5)} \times \frac{\text{SU}(3) \times \text{SU}(3)'}{\text{SU}(3)} \times \text{U}(1)'$$

# A concrete model of compositeness

- model predicts a number of exotics phenomenological implications

$$G_F/H_F = \frac{SU(5)}{SO(5)} \times \frac{SU(3) \times SU(3)'}{SU(3)} \times U(1)'$$

[CE, Schichtel, Spannowsky '17]



$$\mathbf{1}_0 + \mathbf{2}_{\pm 1/2} + \mathbf{3}_0 + \mathbf{3}_{\pm 1}$$

top partners and  
top coupling  
modifications

hyperpions

[Belyaev et al. '17]

[Ferretti '14]

[Matsedonskyi, Panico, Wulzer '15]

[Brown, CE, Galler, Stylianou '20]

- Higgs coupling constraints
- compatibility with exotics searches
- cosmology
- here: focus on elw top properties

$$\begin{aligned} J_{W+}^\mu/e = & c_{XT} \bar{X} \gamma^\mu T + c_{XY} \bar{X} \gamma^\mu Y + c_{XR} \bar{X} \gamma^\mu R \\ & + c_{TB} \bar{T} \gamma^\mu B + c_{YB} \bar{Y} \gamma^\mu B + c_{RB} \bar{R} \gamma^\mu B , \end{aligned}$$

....

# A concrete model of compositeness

- model predicts a number of exotics phenomenological implications

$$G_F/H_F = \frac{SU(5)}{SO(5)} \times \frac{SU(3) \times SU(3)'}{SU(3)} \times U(1)'$$

partial compositeness  
(MCHM5 “lookalike”)

[Agashe, Contino, Pomarol '04]  
[Contino, da Rold, Pomarol '06]

$$\begin{aligned} -\mathcal{L} \supset & M \bar{\Psi} \Psi + \lambda_q f \hat{Q}_L \Sigma \Psi_R + \lambda_t f \hat{t}_R \Sigma^* \Psi_L \\ & + \sqrt{2} \mu_b \text{Tr}(\hat{Q}_L U \hat{b}_R) + \text{h.c.} \end{aligned}$$

[Ferretti '14]

gauge interactions

$$\mathcal{L} \supset \bar{\Psi} \gamma^\mu \left( \frac{2}{3} e A_\mu - \frac{2}{3} t_w e Z_\mu + v_\mu + K p_\mu \right) \Psi$$

$$\Psi = \frac{1}{\sqrt{2}} \begin{pmatrix} iB - iX \\ B + X \\ iT + iY \\ -T + Y \\ \sqrt{2}iR \end{pmatrix} \quad \hat{Q}_L = \begin{pmatrix} ib_L \\ b_L \\ it_L \\ -t_L \\ 0 \end{pmatrix}, \quad \hat{t}_R = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \\ t_R \end{pmatrix}, \quad \hat{b}_R = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \\ b_R \end{pmatrix}$$

$$(T, B) \in (\mathbf{3}, \mathbf{2})_{1/6}, \quad R \in (\mathbf{3}, \mathbf{1})_{2/3}, \quad (X, Y) \in (\mathbf{3}, \mathbf{2})_{7/6}.$$



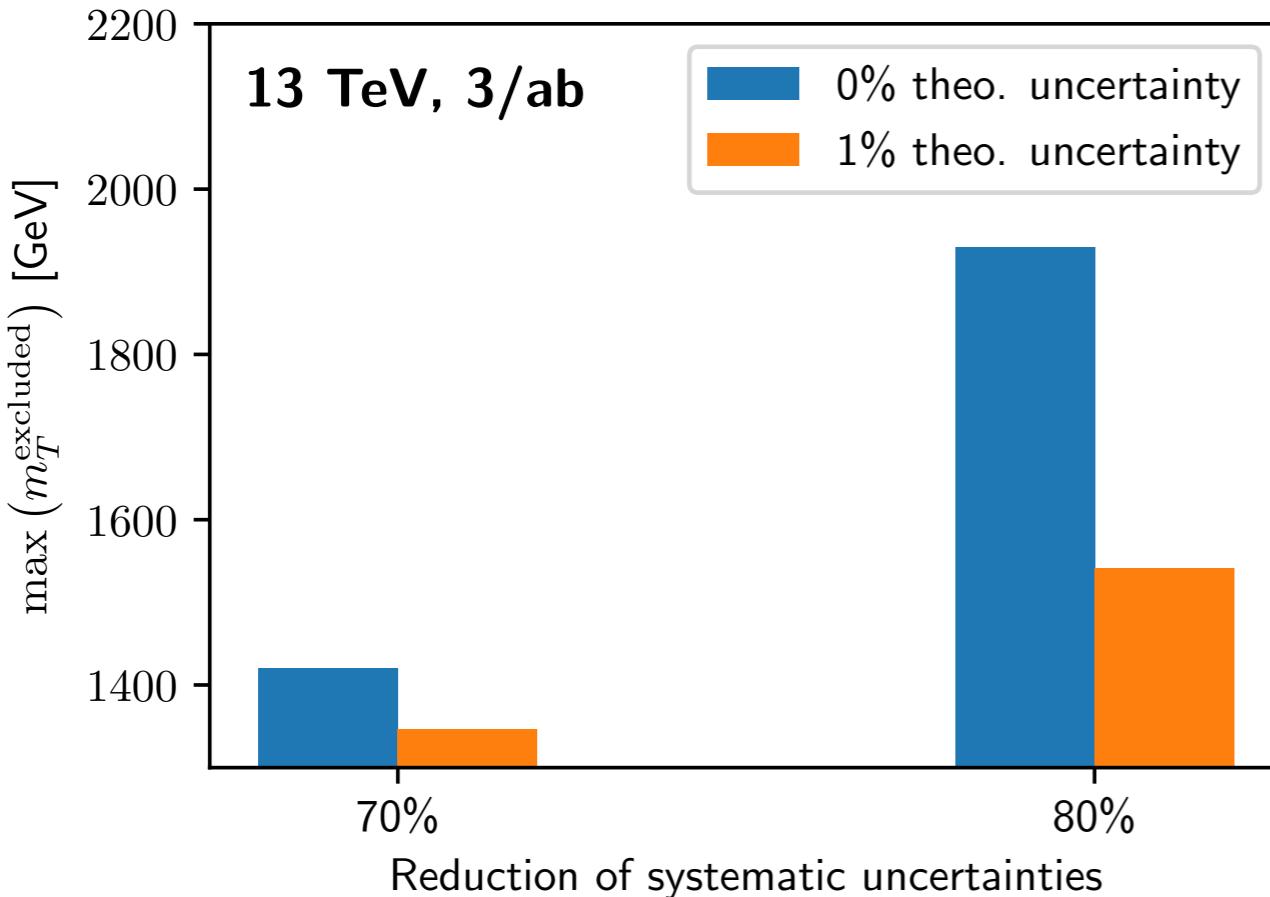
- include range of data (for extrapolation)

Analysis	Collaboration	$\sqrt{s}$ [TeV]	Observables	dof	Analysis	Collaboration	$\sqrt{s}$ [TeV]	Observables	dof					
single top $t$ -channel														
1503.05027 [45]	CDF, D0	1.96	$\sigma_{\text{tot}}$	1	1509.05276 [55]	ATLAS	8	$\sigma_{\text{tot}}$	1					
1406.7844 [46]	ATLAS	7	$\frac{\sigma_t}{\sigma_{\bar{t}}}, \frac{1}{\sigma} \frac{d\sigma}{dp_t^{\perp}}, \frac{1}{\sigma} \frac{d\sigma}{dp_{\bar{t}}^{\perp}}, \frac{1}{\sigma} \frac{d\sigma}{d y_t }, \frac{1}{\sigma} \frac{d\sigma}{d y_{\bar{t}} }$	1, 8, 6	1510.01131 [56]	CMS	8	$\sigma_{\text{tot}}$	1					
1902.07158 [47]	ATLAS, CMS	7, 8	$\sigma_{\text{tot}}$	2	1901.03584 [57]	ATLAS	13	$\sigma_{\text{tot}}$	1					
1609.03920 [48]	ATLAS	13	$\sigma_t, \frac{\sigma_t}{\sigma_{\bar{t}}}$	2	1907.11270 [58]	CMS	13	$\sigma_{\text{tot}}, \frac{1}{\sigma} \frac{d\sigma}{dp_Z^{\perp}}, \frac{1}{\sigma} \frac{d\sigma}{d \cos \theta_Z^*}$	4, 3					
1812.10514 [49]	CMS	13	$\frac{\sigma_t}{\sigma_{\bar{t}}}, \sigma_t$	2	$W$ boson helicity fractions									
single top $s$ -channel														
1402.5126 [50]	CDF, D0	1.96	$\sigma_{\text{tot}}$	1	1211.4523 [59]	CDF	1.96	$F_0, F_R$	2					
1902.07158 [47]	ATLAS, CMS	7, 8	$\sigma_{\text{tot}}$	2	1205.2484 [60]	ATLAS	7	$F_0, F_L, F_R$	3					
$tW$														
1902.07158 [47]	ATLAS, CMS	7, 8	$\sigma_{\text{tot}}$	2	1308.3879 [61]	CMS	7	$F_0, F_L, F_R$	3					
1612.07231 [51]	ATLAS	13	$\sigma_{\text{tot}}$	1	1612.02577 [62]	ATLAS	8	$F_0, F_L$	2					
top quark decay width														
1902.07158 [47]	ATLAS, CMS	7, 8	$\sigma_{\text{tot}}$	2	1201.4156 [63]	D0	1.96	$\Gamma_t$	1					
1612.07231 [51]	ATLAS	13	$\sigma_{\text{tot}}$	1	1308.4050 [64]	CDF	1.96	$\Gamma_t$	1					
1805.07399 [52]	CMS	13	$\sigma_{\text{tot}}$	1	1709.04207 [65]	ATLAS	8	$\Gamma_t$	1					
$tjZ$														
1710.03659 [53]	ATLAS	13	$\sigma_{\text{tot}}$	1	see also									
1812.05900 [54]	CMS	13	$\sigma_{\text{tot}}$	1	[TopFitter `15 `16]									

+ checks that resonance contributions are negligible away from resonance

[SMEFiT `19]  
[SFitter `19]  
[Durieux et al. `19]

# indirect top sector constraints



$$\begin{aligned} \mathcal{L} \supset & \bar{t}\gamma^\mu [g_L^t P_L + g_R^t P_R] tZ_\mu \\ & + \bar{b}\gamma^\mu [g_L^b P_L + g_R^b P_R] bZ_\mu \\ & + (\bar{b}\gamma^\mu [V_L P_L + V_R P_R] tW_\mu^+ + \text{h.c.}) \end{aligned}$$

$$V_L = -\frac{g}{\sqrt{2}} [1 + \delta_{W,L}] \quad \text{etc.}$$

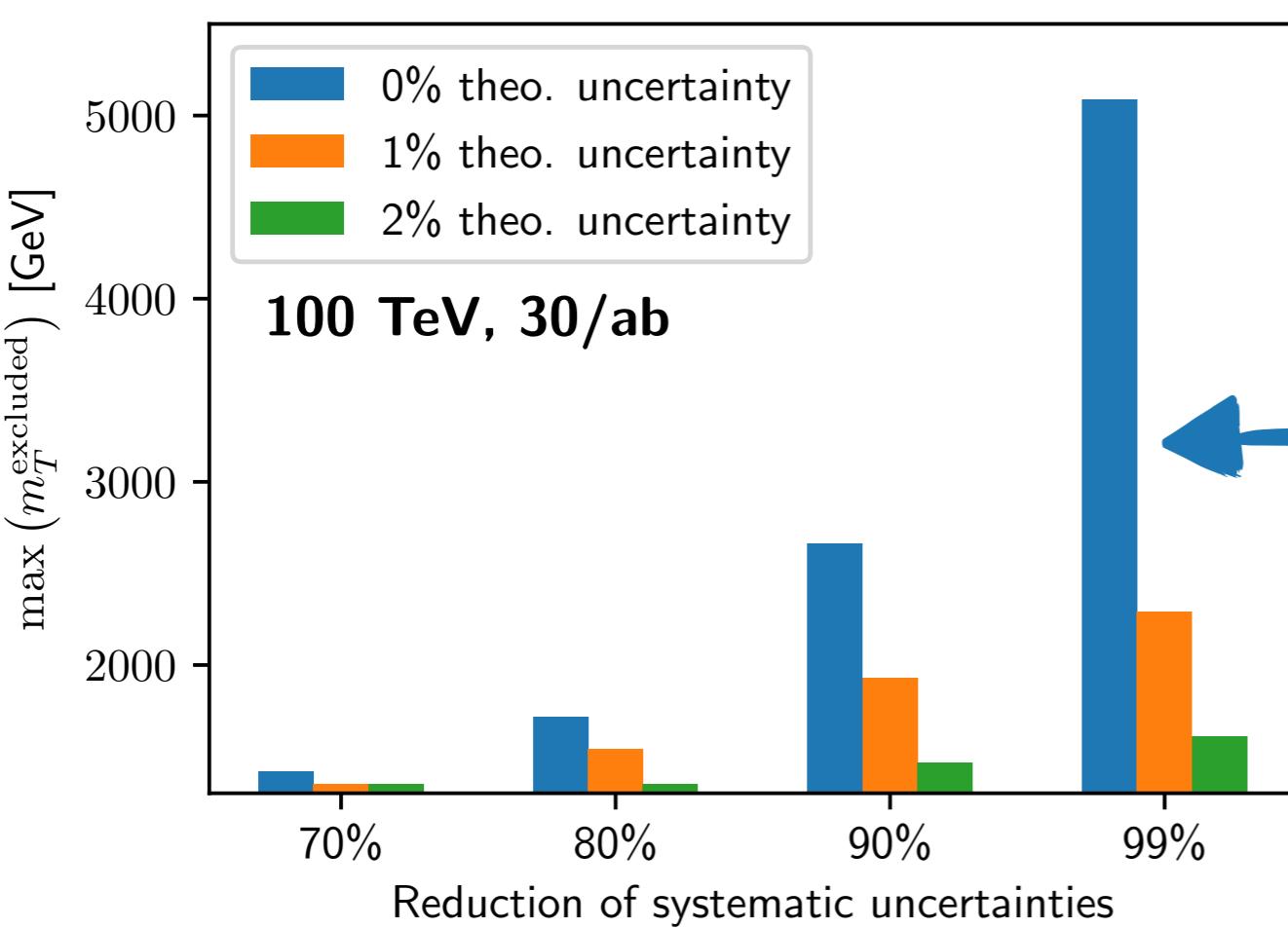
$$\begin{aligned} \delta_{W,L} &\in [-0.029, 0.019], & \delta_{W,R} &\in [-0.009, 0.009], \\ \delta_{Z,L}^t &\in [-0.639, 0.277], & \delta_{Z,R}^t &\in [-1.566, 1.350]. \end{aligned}$$

model correlations

$$\begin{aligned} \delta_{W,L} &\in [-0.025, 0.02], & \delta_{W,R} &\in [-0.0014, 0.0013], \\ \delta_{Z,L}^t &\in [-0.073, 0.06], & \delta_{Z,R}^t &\in [-0.33, 0.37] \end{aligned}$$

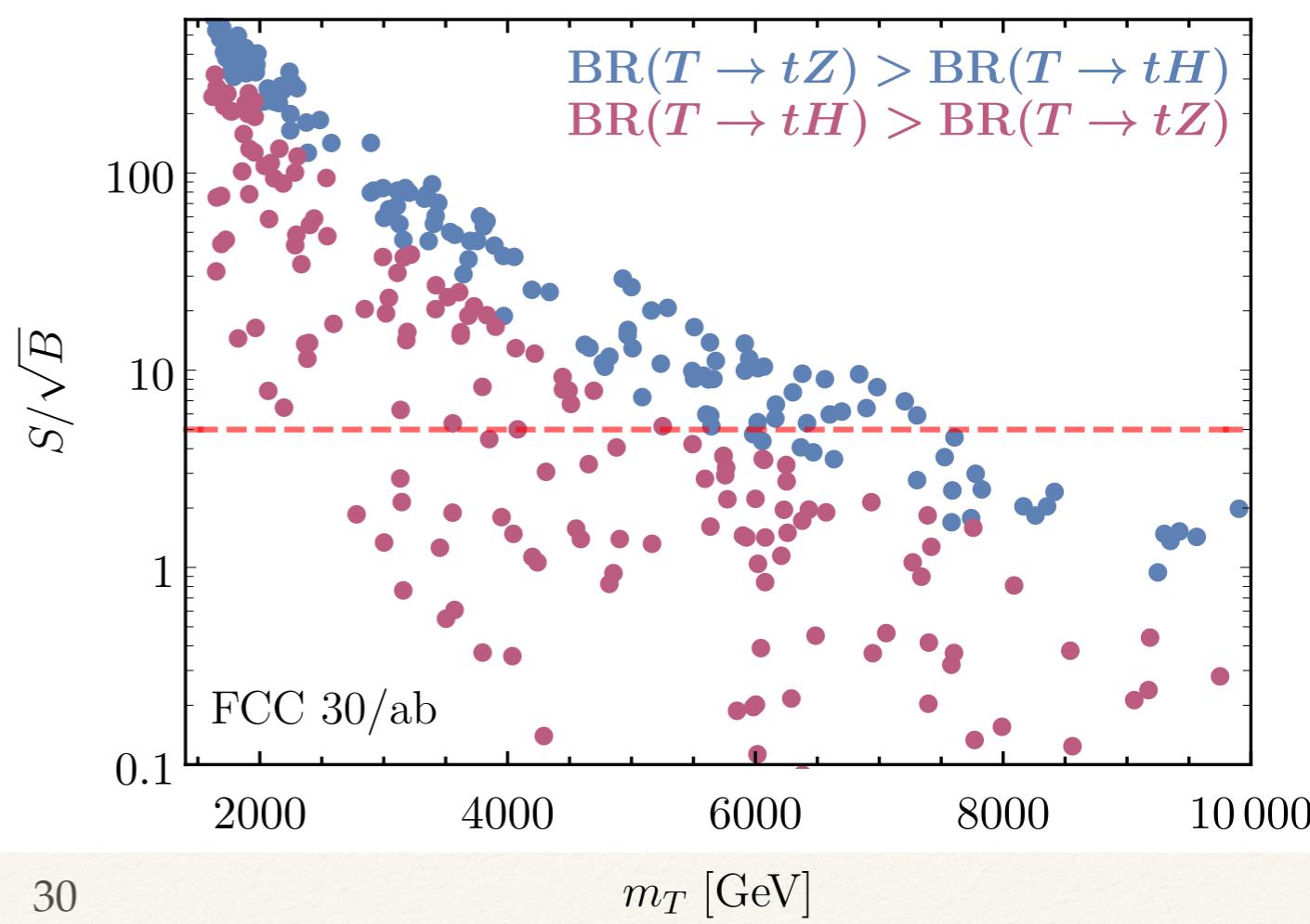
- existing direct top partner constraints in the range of  $\gtrsim 1.5$  TeV compatible [Matsedonskyi, Panico, Wulzer '15]
- theoretical uncertainties is main sensitivity limitation, adding additional channels does not change this picture dramatically

- direct top partner searches in electroweak channels providing direct sensitivity up to 8 TeV
  - [de Simone et al. '14]
  - [Azatov et al. '14]
  - [Matsedonskyi et al. '14]
  - [Golling et al. '16]
  - [Barducci et al. '17]
  - [Li et al. 19]



## indirect top sector constraints

- optimistic extrapolations provide indirect sensitivity up to about 5 TeV



$\max(m_T^{\text{excluded}})$  [GeV]

5000

4000

3000

2000

70%

80%

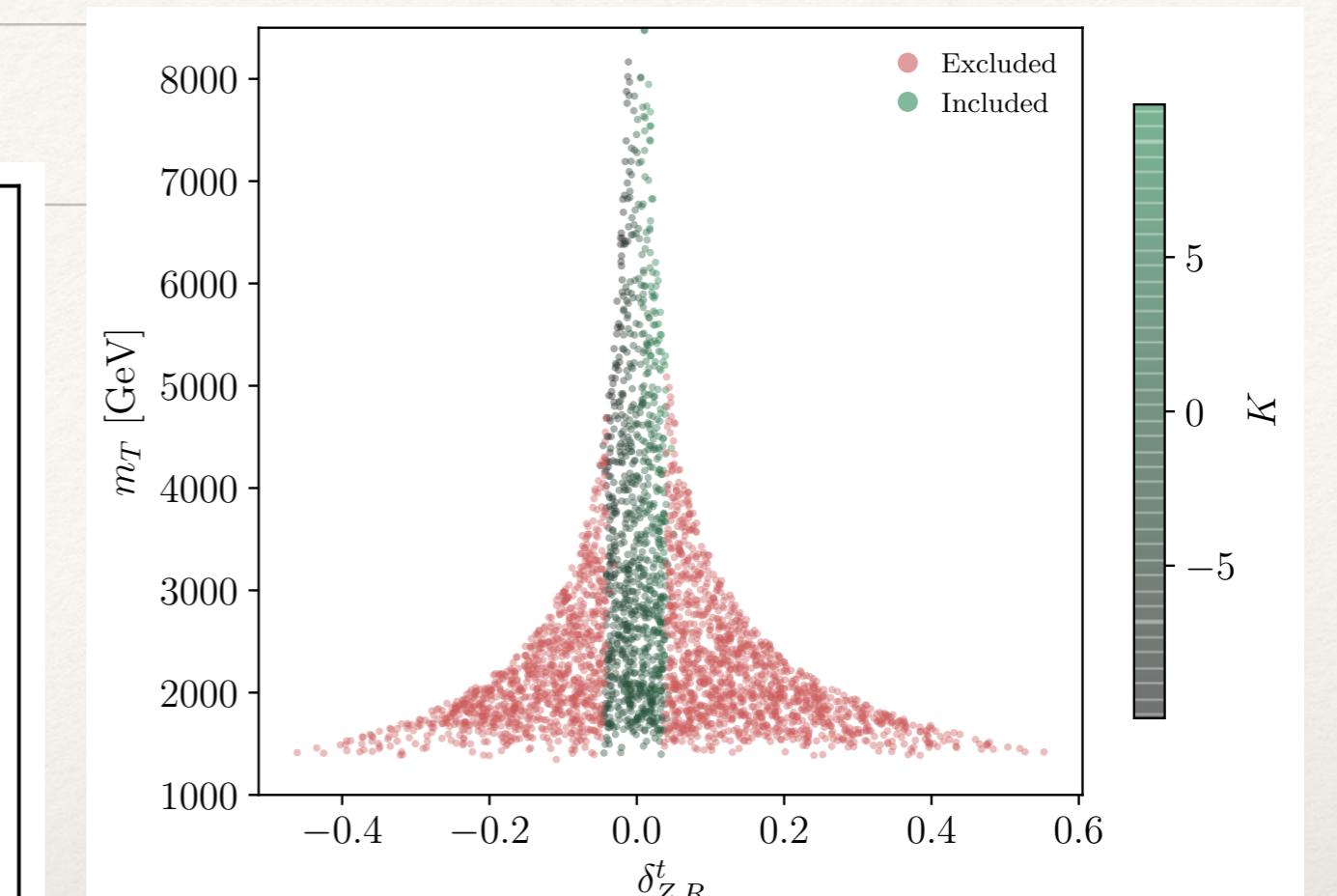
90%

99%

Reduction of systematic uncertainties

- 0% theo. uncertainty
- 1% theo. uncertainty
- 2% theo. uncertainty

100 TeV, 30/ab



- direct top partner searches in electroweak channels providing direct sensitivity up to 8 TeV

[de Simone et al. '14]

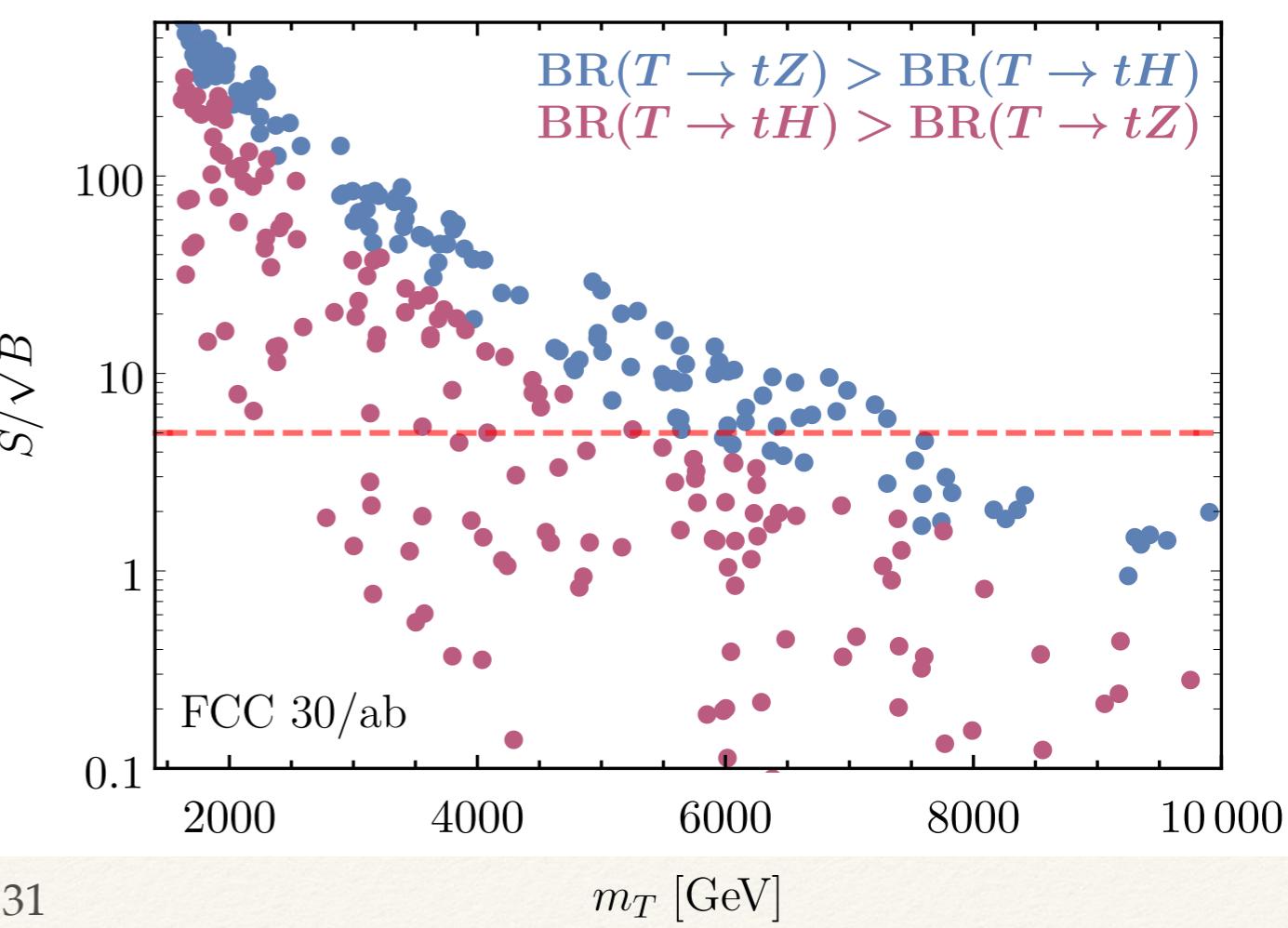
[Azatov et al. '14]

[Matsedonskyi et al. '14]

[Golling et al. '16]

[Barducci et al. '17]

[Li et al. '19]



- ▶ *EFT @ colliders progress has been rapid*
  - ▶ matching, validity re:momentum coverage at hadron machines
  - ▶ but no sensitivity when uncertainties are large
  - ▶ uncertainties/deviations crucial for continued EFT efforts to be fruitful; adopt UV inspired-restrictions as way out?
- ▶ *Opportunity to link the Higgs/top sector to new physics*
  - ▶ cure SM shortcomings (CP violation, hierarchy, DM, ...)
  - ▶ (multi-)Higgs/(multi-)top production as an avenue for BSM
  - ▶ LHC not enough to achieve this in full glory